

# Transverse Target Asymmetry in Exclusive Charged Pion Production at 11 GeV

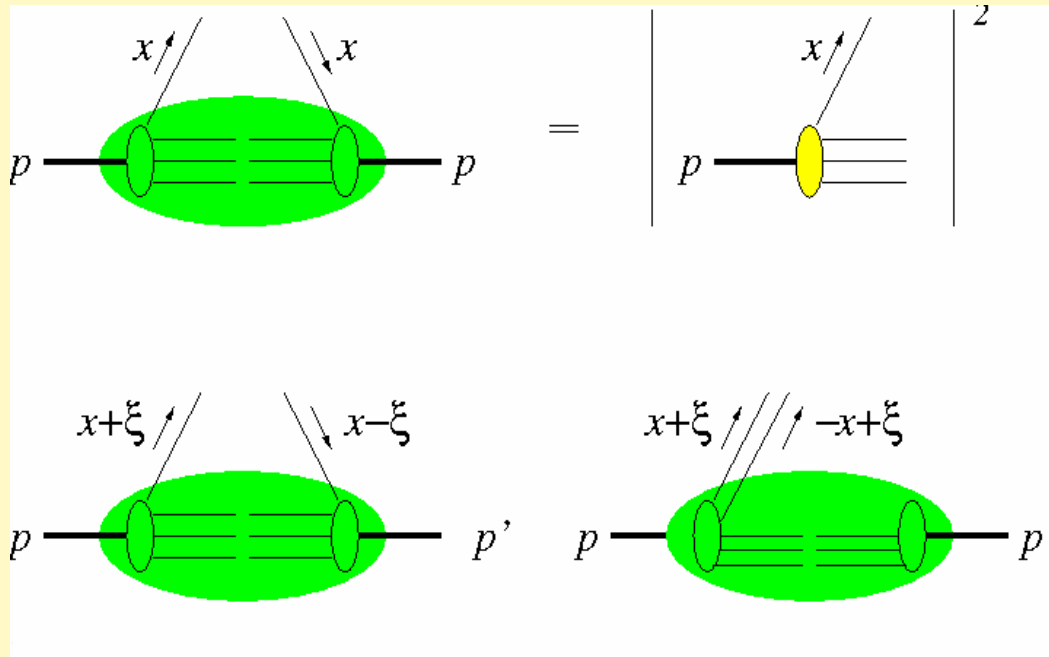
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# Generalized Parton Distributions

Over the last decade, tremendous progress has been made on the theory of generalized parton distributions (GPD).

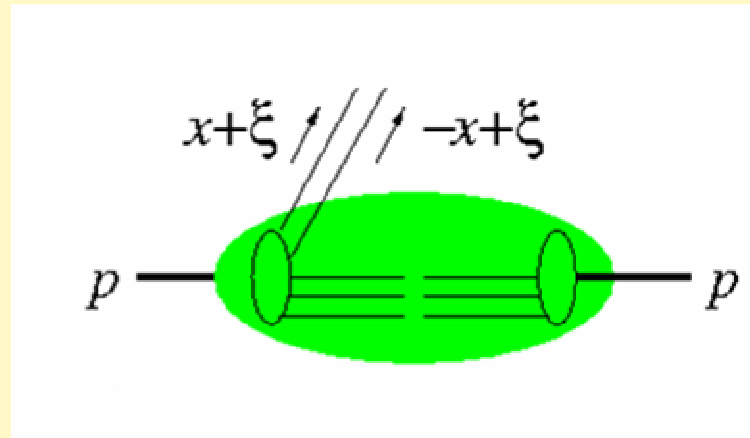


**PDFs** : Squared hadronic wavefunctions = probability of finding a parton with specified longitudinal momentum fraction and polarization in fast moving hadron.

**GPDs** : interference between wavefn's of parton with momentum fraction  $x+\xi$  and parton with momentum fraction  $x-\xi$ .

- In addition to  $x$  and  $\xi$ , GPDs depend also on  $t = -(p-p')^2$ .
  - $t$  is independent of  $x$ ,  $\xi$  since  $p$ ,  $p'$  may differ in either their longitudinal or transverse components.
- **GPDs interrelate the longitudinal and transverse momentum structure of partons within a fast moving hadron.**

# GPDs in Deep Exclusive Meson Production



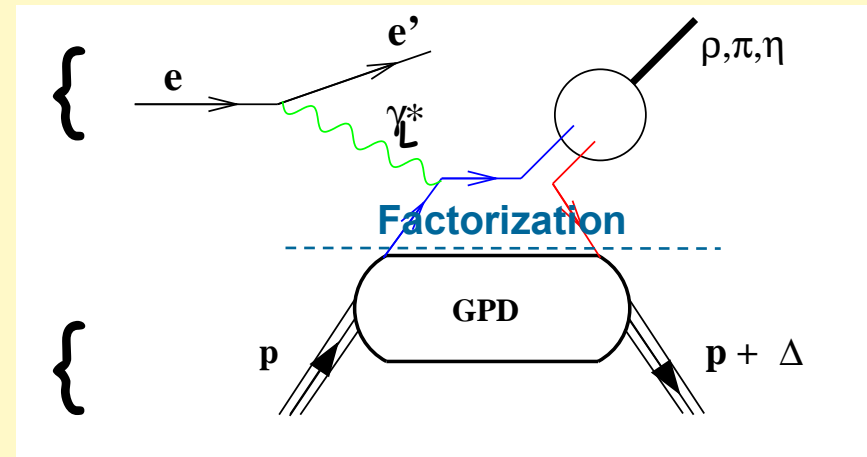
- A special kinematic regime is probed in deep exclusive meson production, where initial hadron emits a  $q\bar{q}$  or  $gg$  pair.
- This has no counterpart in usual PDFs.
- Since GPDs correlate different parton configurations in the hadron at the quantum mechanical level,
  - GPDs determined in this regime carry information about  $q\bar{q}$  and  $gg$ -components in the hadron wavefunction.

# GPDs require Hard Exclusive Reactions

- In order to access the physics contained in GPDs, one is restricted to the hard scattering regime.

- Factorization property of hard reactions:

- Hard probe creates a small size  $q\bar{q}$  and gluon configuration,
  - interactions can be described by pQCD.
- Non-perturbative part describes how hadron reacts to this configuration, or how the probe is transformed into hadrons (parameterized by GPDs).



- Hard exclusive meson electroproduction first shown to be factorizable by Collins, Frankfurt & Strikman [PRD 56(1997)2982].
- Factorization applies when the  $\gamma^*$  is longitudinally polarized.
  - corresponds to small size configuration compared to transversely polarized  $\gamma^*$ .

# Leading Twist GPD Parameterization

- **GPDs are universal quantities and reflect nucleon structure independently of the probing reaction.**

- At leading twist-2, four quark chirality conserving GPDs for each quark, gluon type.
- Because quark helicity is conserved in the hard scattering regime, the produced meson acts as a helicity filter.

$H^{q,g}(x, \xi, t)$   
spin avg  
no hel. flip

$E^{q,g}(x, \xi, t)$   
spin avg  
helicity flip

$\tilde{H}^{q,g}(x, \xi, t)$   
spin diff  
no hel. flip

$\tilde{E}^{q,g}(x, \xi, t)$   
spin diff  
helicity flip

## Leading order QCD predicts:

- Vector meson production sensitive to unpolarized GPDs,  $H$  and  $E$ .
- Pseudoscalar mesons sensitive to polarized GPDs,  $\tilde{H}$  and  $\tilde{E}$ .

# Links to other nucleon structure quantities

- First moments of GPDs are related to nucleon elastic form factors through model-independent sum rules:

$$\sum_q e_q \int_{-1}^{+1} dx H^q(x, \xi, t) = F_1(t)$$

$$\sum_q e_q \int_{-1}^{+1} dx E^q(x, \xi, t) = F_2(t)$$

$$\sum_q e_q \int_{-1}^{+1} dx \tilde{H}^q(x, \xi, t) = G_A(t)$$

$$\sum_q e_q \int_{-1}^{+1} dx \tilde{E}^q(x, \xi, t) = G_P(t)$$

Dirac and Pauli elastic nucleon form factors.  
 $t$ -dependence fairly well known.

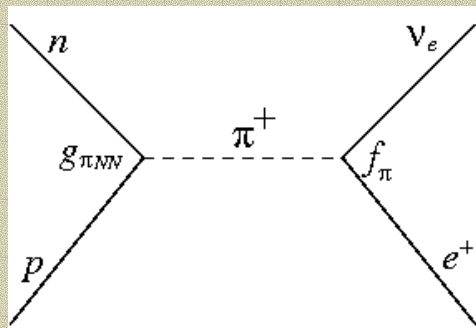
Isovector axial form factor.  
 $t$ -dep. poorly known.

Pseudoscalar form factor.  
Very poorly known.

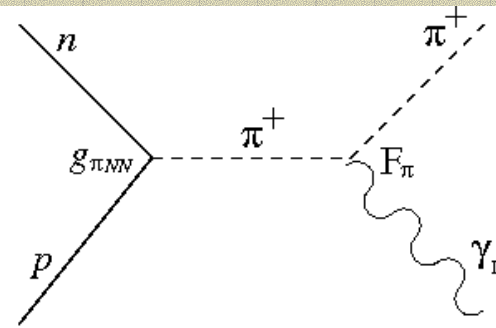
# Spin-flip GPD $\tilde{E}$

$$\sum_q e_q \int_{-1}^{+1} dx \tilde{E}^q(x, \xi, t) = G_p(t)$$

- $G_p(t)$  is highly uncertain because it is negligible at the momentum transfer of  $\beta$ -decay.
- Because of PCAC,  $G_p(t)$  alone receives contributions from  $J^{PG}=0^-$  states.
  - These are the quantum numbers of the pion, so  $\tilde{E}$  contains an important pion pole contribution.



Pion pole contribution to  $G_p(t)$



Pion pole contribution to meson electroproduction at low  $-t$ .

**For this reason, a pion pole-dominated ansatz is typically assumed:**

$$\tilde{E}^{u,d}(x, \xi, t) = F_{\pi}(t) \frac{\theta(\xi > |x|)}{2\xi} \phi_{\pi} \left( \frac{x + \xi}{2\xi} \right)$$

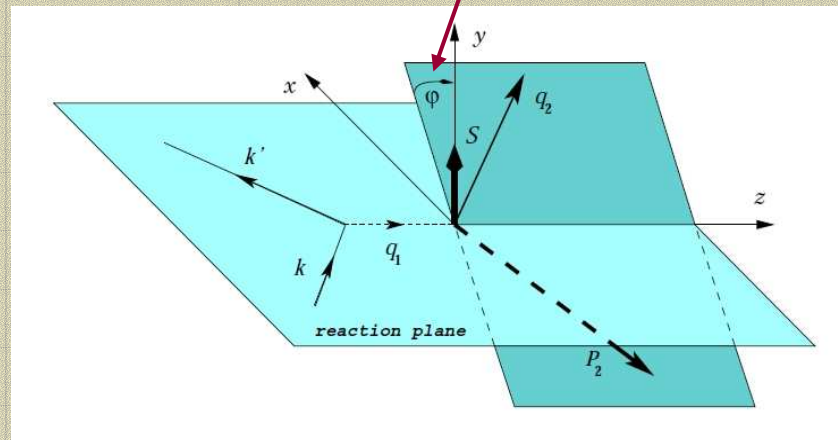
where  $F_{\pi}$  is the pion FF and  $\phi_{\pi}$  the pion PDF.

# How to determine $\tilde{E}$

- **GPD  $\tilde{E}$  not related to an already known parton distribution.**
- **Experimental information on  $\tilde{E}$  can provide new nucleon structure info unlikely to be available from any other source.**
- The most sensitive observable to probe  $\tilde{E}$  is the transverse single-spin asymmetry in exclusive  $\pi$  production:

$$A_{\perp} = \frac{\int_0^{\pi} d\beta \frac{d\sigma_{\pi}^{\pi^-}}{d\beta} - \int_{\pi}^{2\pi} d\beta \frac{d\sigma_{\pi}^{\pi^-}}{d\beta}}{\int_0^{2\pi} d\beta \frac{d\sigma_{\pi}^{\pi^-}}{d\beta}}$$

$d\sigma_{\pi}^L$  = exclusive  $\pi$  cross section for longitudinal  $\gamma^*$   
 $\beta$  = angle between transversely polarized target vector and the reaction plane.



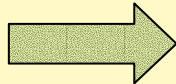


# Single Spin Asymmetry in Exclusive $\pi$ Production

- Frankfurt et al. [PRD 60(1999)014010] have shown the  $A_L^\perp$  asymmetry vanishes if  $\tilde{E}$  is zero.
  - If  $\tilde{E} \neq 0$ , the asymmetry will display a  $\sin\beta$  dependence.
- $A_L^\perp$  is also expected to display precocious factorization at moderate  $Q^2 \sim 2-4 \text{ GeV}^2$ .
  - Precocious factorization of the  $\pi$  production amplitude into three blocks is likely:
    1. overlap integral between  $\gamma, \pi$  wave functions.
    2. the hard interaction.
    3. the GPD.
  - Higher order corrections, which may be significant at low  $Q^2$ , likely cancel in the asymmetry ratio.

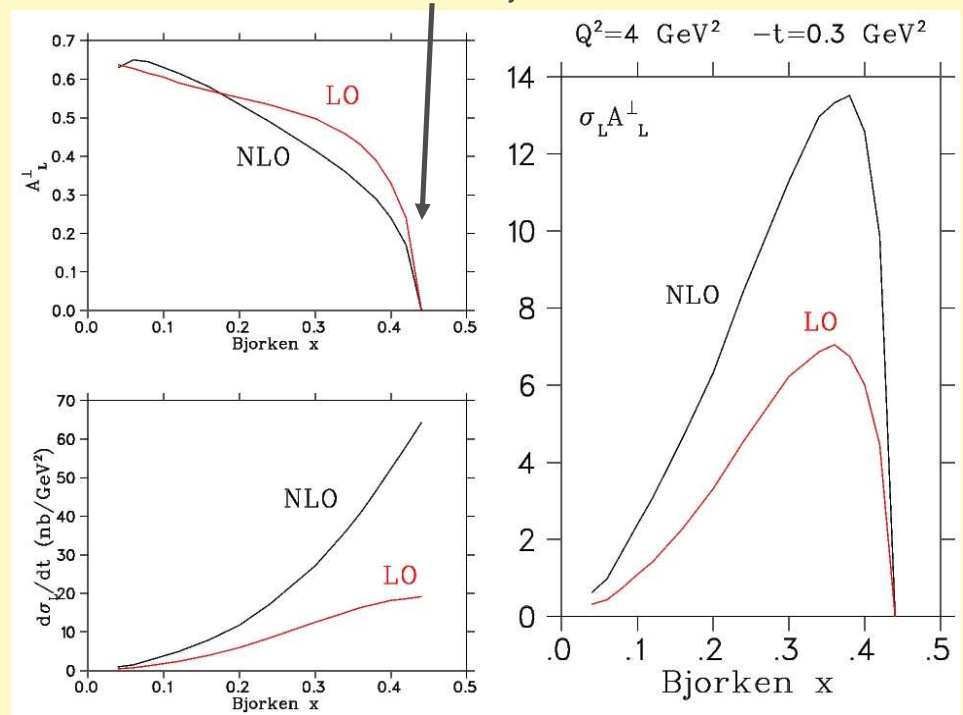
# Cancellation of Higher Twist Corrections in $A_L^\perp$

- **Belitsky and Müller GPD based calc. reinforces this expectation:**



- Even at  $Q^2=10 \text{ GeV}^2$ , NLO effects can be large, but cancel in the asymmetry,  $A_L^\perp$  (PL B513(2001)349).
- At  $Q^2=4 \text{ GeV}^2$ , higher twist effects even larger in  $\sigma_L$ , but still cancel in asymmetry (CIPANP 2003).

$A_L^\perp=0$  at parallel kinematic limit, where  $P_y$  is not well defined.



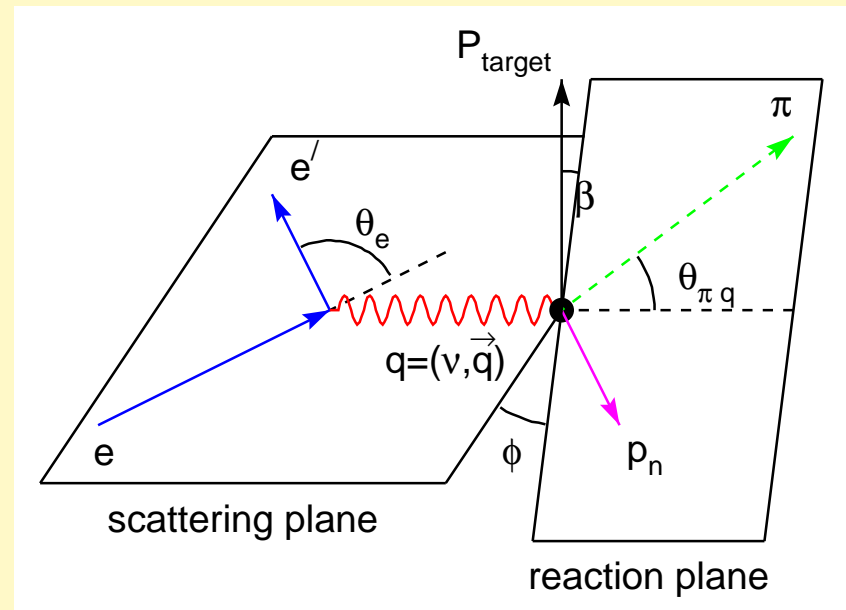
**This relatively low value of  $Q^2$  for the expected onset of precocious scaling is important, because it is experimentally accessible at JLab 12 GeV.**

# L/T Separations Essential

- In hard meson electroproduction, factorization can only be applied to **longitudinal photons**.
- Unlike other ongoing or proposed experiments, where dominance of longitudinal contribution is simply assumed, JLab's unique contribution to this field is in:
  - ability to take measurements at multiple beam energies.
  - unambiguous isolation of  $A_L^\perp$  using Rosenbluth separation.
- A JLab  $A_L^\perp$  measurement could thus establish the applicability of the GPD formalism, and precocious scaling expectations, for other  $A^\perp$  experiments.

# Require Target Polarization Parallel to $\hat{q} \times \hat{p}_\pi$

- Target polarization components ( $P_x, P_y$ ) are defined relative to reaction plane.
- $\beta$  = azimuthal angle between (transverse) target polarization and reaction plane
- $P_x = P_\perp \cos\beta$  and  $P_y = P_\perp \sin\beta$
- $P_y \parallel \hat{q} \times \hat{p}_\pi$  uniquely defined only in non-parallel kinematics.



Unpolarized  
Cross section

$$\frac{d\sigma}{d\Omega} = \sigma_T + \epsilon\sigma_L + \sqrt{\frac{1}{2}\epsilon(\epsilon+1)}\sigma_{LT} \cos\phi + \epsilon\sigma_{TT} \cos 2\phi$$

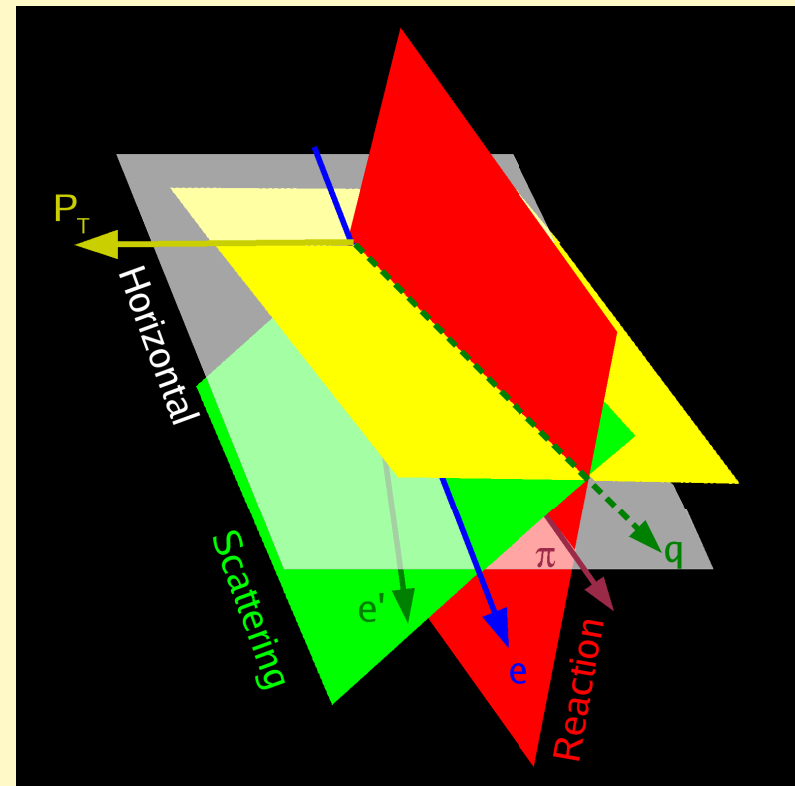
$$A_\perp = \frac{1}{P_\perp} \frac{2}{\pi} \frac{2\sigma_L^y}{\sigma_L}$$

$$\begin{aligned} \sigma_t = & P_x \left[ -\sqrt{2\epsilon(1+\epsilon)} \sin\phi \sigma_{LT}^x - \epsilon \sin 2\phi \sigma_{TT}^x \right] \\ & - P_y \left[ \sigma_{TT}^y + \epsilon \cos 2\phi \sigma_{TT}^y + 2\epsilon \sigma_L^y + \sqrt{2\epsilon(1+\epsilon)} \cos\phi \sigma_{LT}^y \right] \\ & + P_z \left[ \epsilon \sin 2\phi \sigma_{TT}^z + \sqrt{2\epsilon(1+\epsilon)} \sin\phi \sigma_{LT}^z \right] \end{aligned}$$

# SHMS+HMS Kinematic Considerations

SHMS and HMS have their largest acceptances in vertical direction.

- **Scattered  $e^-$  detected at some vertical angle in the HMS.**
  - Forces scattering plane and  $q$ -vector to be non-horizontal.
- $\pi^-$  is detected in SHMS.
  - Either above or below  $q$ -vector, depending if scattered  $e^-$  is detected above or below horizontal plane.
- Target polarization is horizontal, parallel to  $\hat{q} \times \hat{p}_\pi$ .
  - Nearly transverse to  $\vec{q}$  for all angles between the scattering and reaction planes.

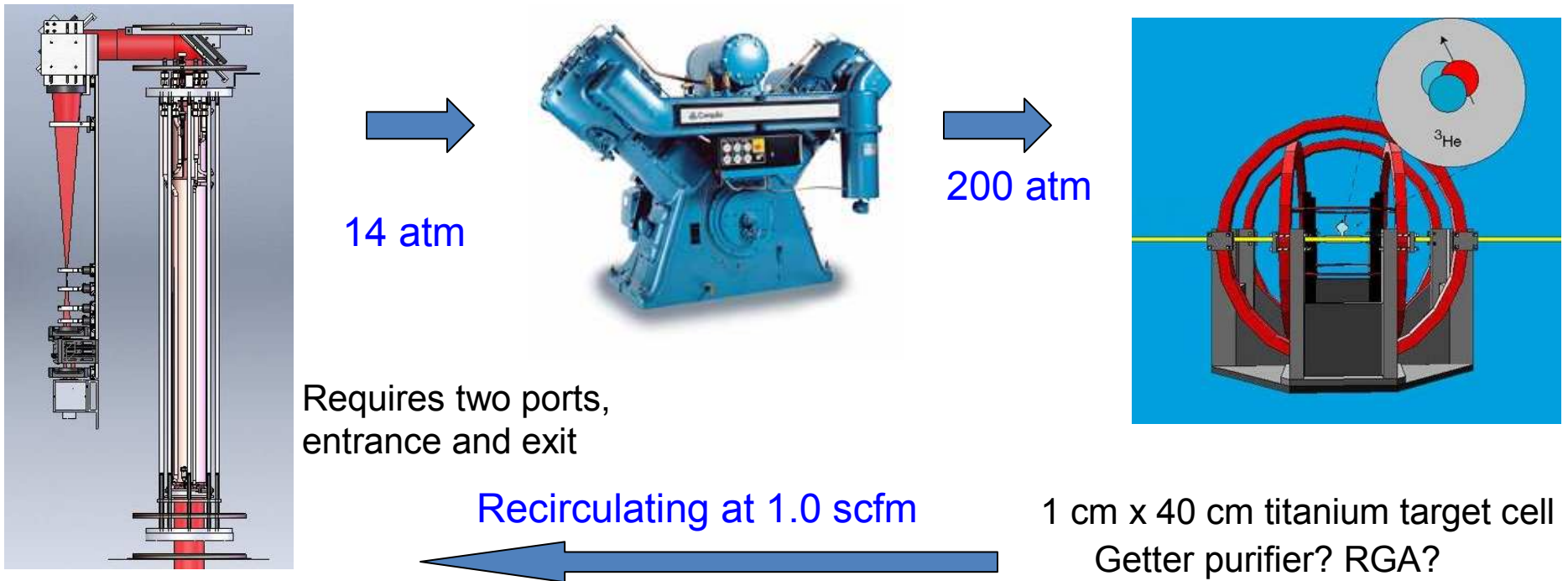


# High Luminosity Essential

- **Physics case for a measurement of  $A_L^\perp$  is compelling.**
- **High luminosity required:**
  - $\sigma_L$  is largest in parallel kinematics, where  $A_L^\perp=0$ .
  - $\sigma_L$  is small where  $A_L^\perp$  is maximal.
- We have performed numerous studies, but the measurement has not been feasible to date because of the lack of a polarized target that can handle the required high luminosity.
- Recent advancements in polarized  $^3\text{He}$  target technology may allow the measurement to proceed via the  $n(e,e'\pi^-)p$  reaction.

# XeMed/UNH Target Loop Concept

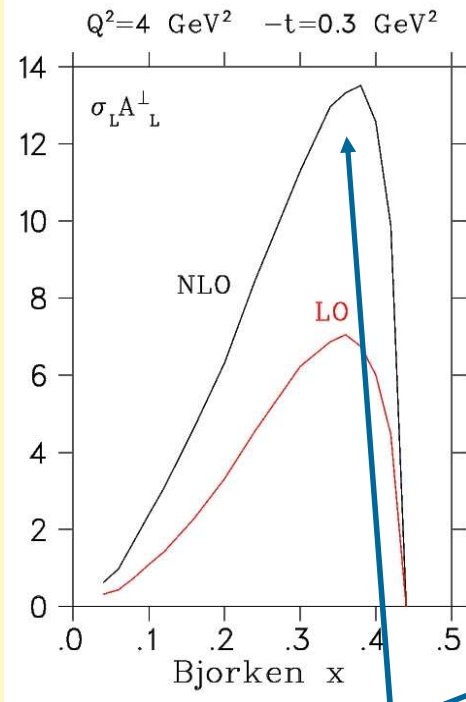
- Compress polarized  $^3\text{He}$  and deliver to 40cm long titanium target cell
- Commercial compressors achieve >3500 psi (238 bar)
- Requires compression ratio ~16, immersion in magnetic field, rubidium-free gas leaving polarizer, entrance and exit, <3% polarization loss



Property	Hall A	UNH
Polarization (%)	55	60
Beam Current ( $\mu\text{A}$ )	15	100
Pressure (atm)	10	200
Cell type	Glass/sealed	Ti/continuous flow
“Spin UP” time (h)	7	4
Beam Relaxation ( $\text{h}^{-1}$ )	41	0.1
Laser Power (W)	150	1500-2500
Thickness ( $\text{cm}^{-2}$ )	1.07E+22	1E+24
<b><i>FOM (P<sup>2</sup>L)</i></b>	0.22E+36	0.55E+38



# SHMS+HMS Kinematics



Near peak of Figure of Merit in Belitsky's calculation

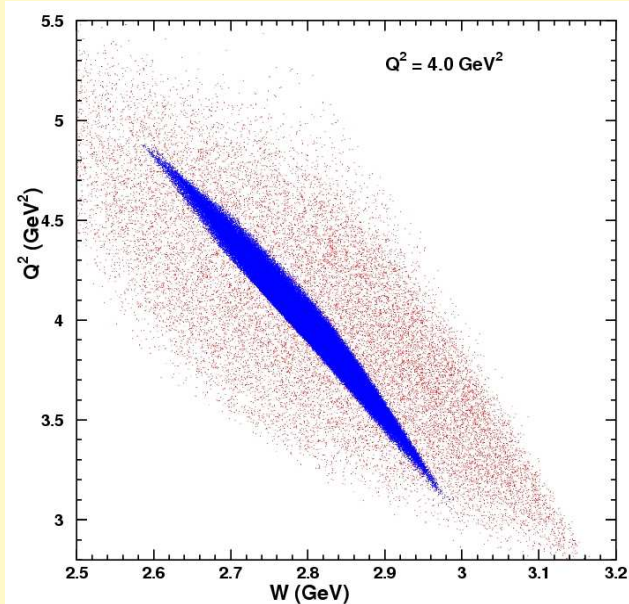
<b>n(e,e'π-)p Kinematics at Q<sup>2</sup>=4.0 GeV<sup>2</sup>, W=2.8 GeV</b>							
<b>E<sub>e</sub></b> (GeV)	<b>θ<sub>e</sub></b> (deg)	<b>p<sub>π</sub></b> (GeV/c)	<b>θ<sub>π</sub></b> (deg)	<b>Θ<sub>πq</sub></b> (deg)	<b>-t</b> (GeV <sup>2</sup> )	<b>x</b>	<b>z</b>
<b>High ε=0.745 Setting, E<sub>beam</sub>=11.00 GeV</b>							
5.160	15.25	5.744	-12.70	0	0.175	0.365	0.984
		5.666		±3	0.322	0.365	0.970
		5.531		±5	0.576	0.365	0.947
<b>Low ε=0.200 Setting, E<sub>beam</sub>=6.60 GeV</b>							
0.860	49.23	5.744	-6.06	0	0.175	0.365	0.984
		5.666		±3	0.322	0.365	0.970
		5.531		±5	0.576	0.365	0.947

- Scattered electron in HMS, π<sup>-</sup> in SHMS.
- Θ<sub>πq</sub> is π lab angle wrt  $\vec{q}$ , mostly above or below scattering plane.
- For Q<sup>2</sup>=4 GeV<sup>2</sup>, x=0.365 → -t<sub>max</sub> ≈ 1 - (M<sup>2</sup>x<sup>2</sup>/Q<sup>2</sup>) = 0.97 GeV<sup>2</sup>.

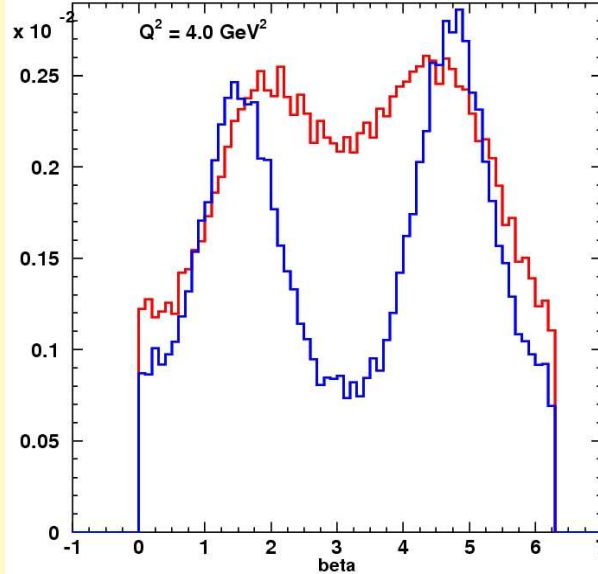
# Simulated SHMS+HMS Acceptance

$$n(e, e' \pi) p$$

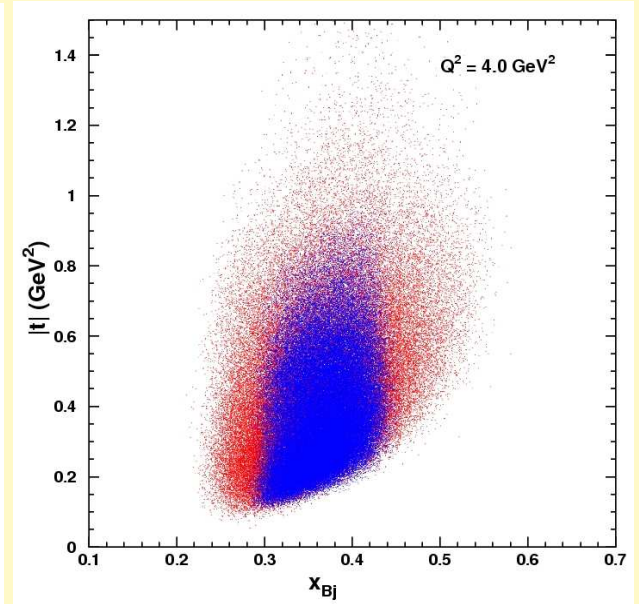
low  $\varepsilon$  high  $\varepsilon$   $\Delta\varepsilon = 0.55$



$Q^2$ - $W$  acceptance at high and low  $\varepsilon$ .

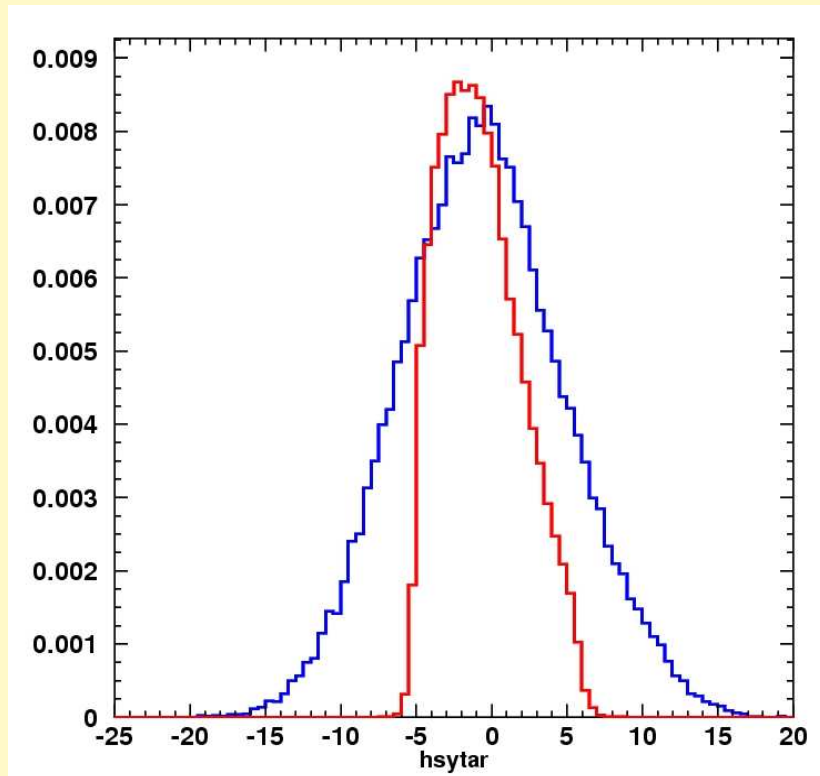


Azimuthal angle of (transversely) polarized target wrt hadron reaction plane.



SHMS+HMS acceptance covers  $0.1 < -t < 0.7$  GeV<sup>2</sup> at nearly fixed  $x_{Bj}$ .

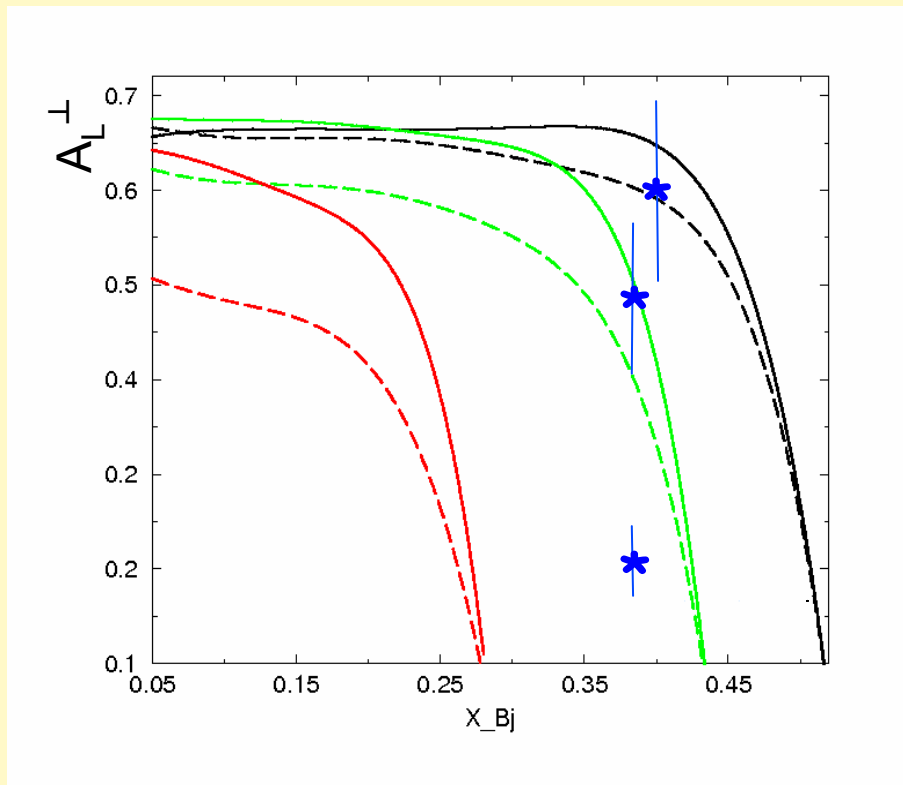
# Reliable L/T Separations require shorter $^3\text{He}$ cells



low  $\epsilon$  high  $\epsilon$   $\Delta\epsilon = 0.55$

- For reliable L/T separations, an 8cm  $^3\text{He}$  cell seems optimal.
- The UNH target is designed for nominal 40cm cells, but Bill Hersman does not believe shorter cells will cause any problem.
- In fact, by cooling the entrance and exit lines of the cell with  $\text{LN}_2$  he believes he can reduce the wall thickness by  $\sim x3$  or increase the pressure while keeping the wall thickness same.
- These issues need closer investigation to better understand target cell backgrounds for UNH target.

# Estimated Rates and Uncertainties



**Solid: asymptotic pion distribution amp.**  
**Dashed: CZ pion dist. amp.**  
**t = -0.5 GeV<sup>2</sup>**  
**t = -0.3 GeV<sup>2</sup>**  
**t = -0.1 GeV<sup>2</sup>**

- Simulated error bars after 18 days:
  - 12days @low  $\epsilon$ , 5k evts in largest  $-t$  bin.
  - 6days @high  $\epsilon$ , 30k evts in largest  $-t$  bin.
- Luminosity= $1.2 \times 10^{37}/\text{cm}^2/\text{s}$  (8cm tgt).
- $P_{\text{targ}}=65\% \rightarrow P_n=55.3\%$ .
  - No target dilution since exclusive  $\pi$  can be only from neutron.
- 2% random systematic uncertainties
  - slightly larger than assumed for Fpi-12.
- $\sigma_L/\sigma_T$  values similar to pionCT  $^1\text{H}$  data.

$Q^2=4.0, W=2.8, x=0.365$			
$-t$ (GeV <sup>2</sup> )	$R=$ $\sigma_L/\sigma_T$	$A_L^\perp$	$\delta A_L^\perp$
0.2	1.0	0.2	0.04
0.4	1.0	0.5	0.08
0.6	1.5	0.6	0.10

# Summary

## PAC24 Comments on our 2003 6 GeV LOI:

- The experiment is extremely challenging since it requires first the isolation of a Fourier component in the polarized target cross section and then, by Rosenbluth techniques, the separation of the cross section for longitudinally polarized photons.
- The measurement may allow for an extraction of further information on GPDs and is complementary to DVCS. Deep virtual electroproduction of pions is sensitive only to the GPDs  $\tilde{H}$  and  $\tilde{E}$ ;  $H$  and  $E$  do contribute.
- Moreover, since the asymmetry requires proton helicity flip, the experiment may allow the extraction of  $\tilde{E}$ , one of the two GPDs not constrained by knowledge of ordinary parton distributions. The measurement is therefore very important.
- **The lack of a transversely polarized cryotarget that can handle the required high luminosity has precluded our development of this experiment since 2003.**
- **A transversely polarized  $^3\text{He}$  target based on the UNH design offers the best hope of measuring this observable.**