

Proton Spin Polarizabilities with Polarized Compton Scattering at MAMI



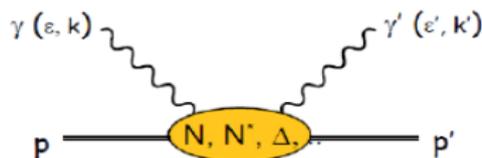
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Nuclear Compton Scattering and Polarizabilities

- Polarizabilities are very low energy fundamental structure constants and Nuclear Compton scattering off a single proton is used to access these Internal structure constants of a nucleon.



- $\gamma(k) + P(p) \rightarrow \gamma(k') + P(p')$
- Low energy outgoing photon plays a role of an applied EM dipole field

$$H_{eff}^{(0)} = \frac{(\vec{p} - e\vec{A})^2}{2m} + e\phi \quad (1)$$

$$H_{eff}^{(1)} = \frac{e(1 + \kappa)}{2m} \vec{\sigma} \cdot \vec{H} - \frac{e(1 + 2\kappa)}{8m^2} \vec{\sigma} \cdot [\vec{E} \times \vec{p} - \vec{p} \times \vec{E}] \quad (2)$$

What are Spin Polarizabilities

- Effective Hamiltonian in second order contains scalar polarizabilities (α_{E1} and β_{M1}) which are the evidence of proton's internal structure

$$H_{eff}^{(2)} = -4\pi \left[\frac{1}{2} \alpha_{E1} \vec{E}^2 + \frac{1}{2} \beta_{M1} \vec{H}^2 \right] \quad (3)$$

- The third order effective Hamiltonian term in the expansion:

$$H_{eff}^{(3)} = -4\pi \left[\frac{1}{2} \gamma_{E1E1} \vec{\sigma} \cdot (\vec{E} \times \vec{E}) + \frac{1}{2} \gamma_{M1M1} \vec{\sigma} \cdot (\vec{H} \times \vec{H}) - \gamma_{M1E2} E_{ij} \sigma_i H_j + \gamma_{E1M2} H_{ij} \sigma_i E_j \right] \quad (4)$$

- These constants (γ) are spin (or vector) polarizabilities (e.g. γ_{M1E2} excited by electric quadrupole $E2$ radiation and decays by magnetic dipole $M1$ radiation.
- They describe the response of the proton spin to an applied electric or magnetic field, 'stiffness' of proton spin against E.M. induced deformations relative to the spin axis.

What do we know about Spin Polarizabilities

	K-mat.	HDPV	DPV	L_χ	$HB_\chi PT$	$B_\chi PT$
γ_{E1E1}	-4.8	-4.3	-3.8	-3.7	-1.1 ± 1.8 (th)	-3.3
γ_{M1M1}	3.5	2.9	2.9	2.5	2.2 ± 0.5 (st) ± 0.7 (th)	3.0
γ_{E1M2}	-1.8	-0.02	0.5	1.2	-0.4 ± 0.4 (th)	0.2
γ_{M1E2}	1.1	2.2	1.6	1.2	1.9 ± 0.4 (th)	1.1
γ_0	2.0	-0.8	-1.1	-1.2	-2.6	-1.0
γ_π	11.2	9.4	7.8	6.1	5.6	7.2

$$\gamma_0 = -\gamma_{E1E1} - \gamma_{E1M2} - \gamma_{M1M1} - \gamma_{M1E2} = (-1.0 \pm 0.08) \times 10^{-4} fm^4 \quad (5)$$

$$\gamma_\pi = -\gamma_{E1E1} - \gamma_{E1M2} + \gamma_{M1M1} + \gamma_{M1E2} = (-8.0 \pm 1.8) \times 10^{-4} fm^4 \quad (6)$$

- Spin-polarizabilities in units of $10^{-4} fm^4$.
- **K-matrix:** Kondratyuk et al., PRC 64, 024005 (2001), **HDPV and DPV (Dispersion Relation):** Holstein et al., PRC 61, 034316 (2000), Drechsel et al., Phys.Rep. 378, 99 (2003), Pasquini et al., PRC 76, 015203 (2007), L_χ (**Chiral Lagrangian**): Gasparyan et al., NP A866, 79 (2011), **$HB_\chi PT$ and $B_\chi PT$ (Heavy Baryon & Covariant Chiral PT) :** J. A. McGovern et al., Eur. Phys. J. A 49, 12 (2013).

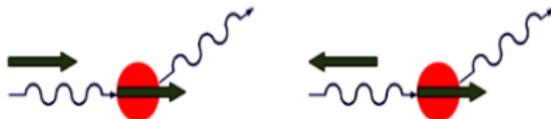
Best Way to extract Spin Polarizabilities

- **Spin polarizabilities appear in the effective interaction Hamiltonian at third order in photon energy**
 - It is in the Δ (1232) resonance region ($E_\gamma = 200 - 300$ MeV) where their effect becomes significant.
- **In this energy region, it is possible to accurately measure polarization asymmetries using a variety of polarized beam and target combinations**
 - The various asymmetries respond differently to the individual spin polarizabilities at different E and θ .
 - Measure three asymmetries at different E, θ .
- **Our plan is to conduct a global analysis:**
 - include constraints from “known” $\gamma_0, \gamma_\pi, \alpha_{E1}$ and β_{M1} .
 - extract all four spin polarizabilities independently with small statistical, systematic and model-dependent errors.

Three Polarization Asymmetry Experiment at A2

Circularly polarized beam, longitudinally polarized target

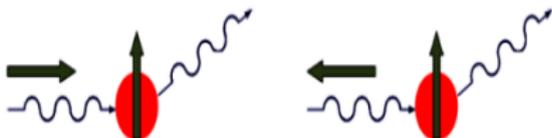
$$\sum_{2z} = \frac{\sigma_{+z}^R - \sigma_{+z}^L}{\sigma_{+z}^R + \sigma_{+z}^L} = \frac{\sigma_{+z}^R - \sigma_{-z}^R}{\sigma_{+z}^R + \sigma_{-z}^R}$$



- \sum_{2z} is sensitive to γ_{M1M1}

Circularly polarized beam, transversely polarized target

$$\sum_{2x} = \frac{\sigma_{+x}^R - \sigma_{+x}^L}{\sigma_{+x}^R + \sigma_{+x}^L} = \frac{\sigma_{+x}^R - \sigma_{-x}^R}{\sigma_{+x}^R + \sigma_{-x}^R}$$

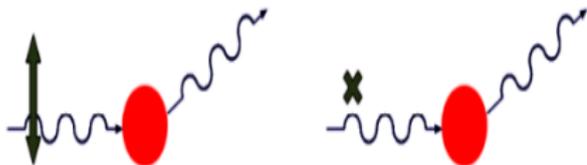


- \sum_{2x} is sensitive to γ_{E1E1}

Polarization Asymmetry (contd..)

Linearly polarized(\parallel and \perp to scattering plane) beam, unpolarized target

$$\Sigma_3 = \frac{\sigma_{\parallel} - \sigma_{\perp}}{\sigma_{\parallel} + \sigma_{\perp}}$$

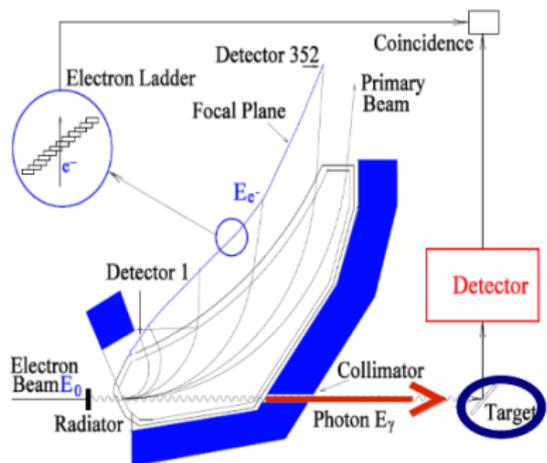


- Σ_{2z} is sensitive to γ_{M1M1}

Σ_{2x} , Σ_3 and Σ_{2z} Experimental data

- Transverse Target(Σ_{2x}): Sep 2010, Feb 2011 - 500 hrs (P. Martel)
- Unpolarized Target(Σ_3): Dec 2012- 150 hrs (C. Collicot)
- **Longitudinal Target** (Σ_{2z}): D. Paudyal (University of Regina) and A. Rajabi (University of Massachusetts)
 - First round of data in 2014 with Butanol (320 hours) and Carbon Target (180 hrs)
 - Second round of data in 2015 with Butanol (310 hours) and Carbon Target (60) hrs
 - Worked as a run coordinator for two weeks during 2015 beam time.

Experimental Apparatus at MAMI

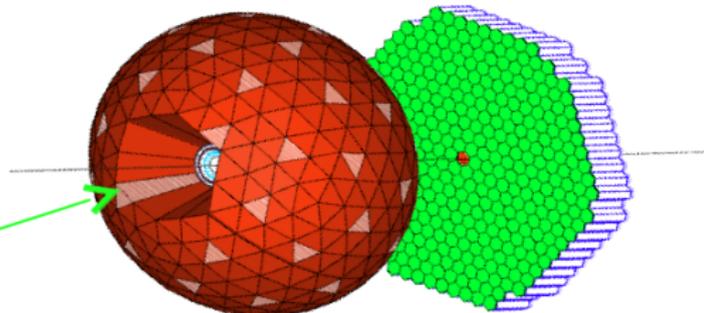


PID

- Cylindrical detectors, 24 thin plastic scintillator strips, Identification of charged particles

- ## Crystal Ball
- 672 NaI crystals, separate PMT and 94 % solid angle coverage

- ## TAPS
- 366 BaF₂, 72 PbWO₄ Crystals and 384 Veto Paddles



MWPC

MWPC between PID and CB for track reconstruction of charged particles

Σ_{2z} Data Analysis Status and Experimental challenges

- Calibration of first round of carbon and butanol target data has been completed
- **Tagging Efficiency and target polarization check has been completed**
- Σ_{2z} — **Experimental Challenges**
 - Small Compton scattering cross sections.
 - Coherent and incoherent reactions off of C, O, and He.
 - A source of polarized protons is not easy to come by (or to operate)
 - In Δ -region, proton tracks are required to suppress backgrounds, but energy losses in the frozen-spin cryostat, and CB-TAPS are considerable.
- What to do ???
 - π^0 photo-production cross section is about 100 times that of Compton scattering, work on π_0 photo-production Asymmetry.

π^0 Production - Background Reaction

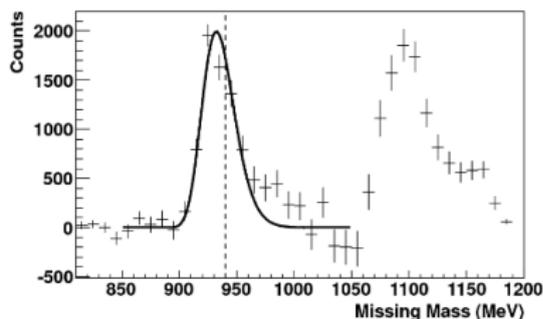
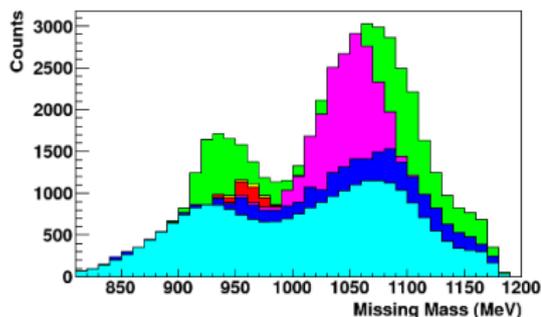
Compton Scattering $\gamma + P \rightarrow \gamma + P$

- Dominant Background for Compton Scattering Experiments

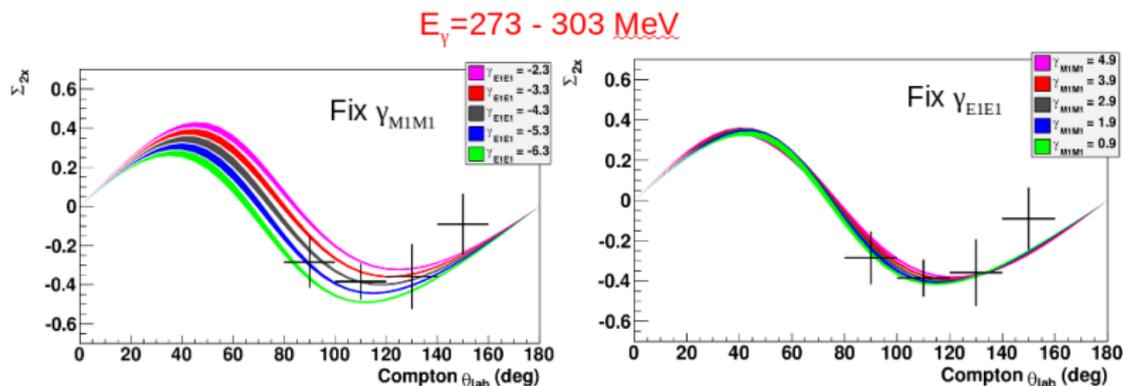


π^0 Production as a Systematic Check

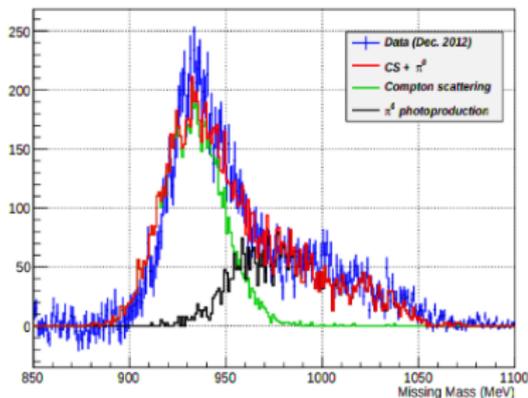
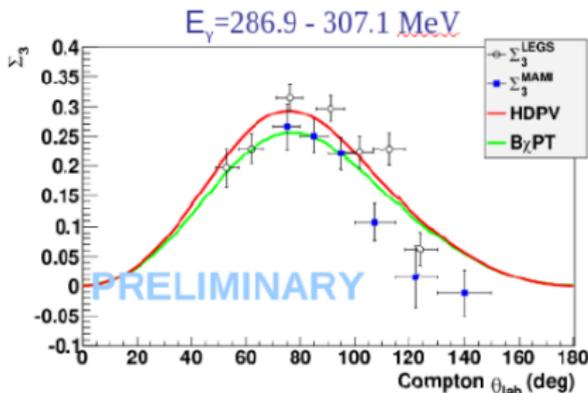
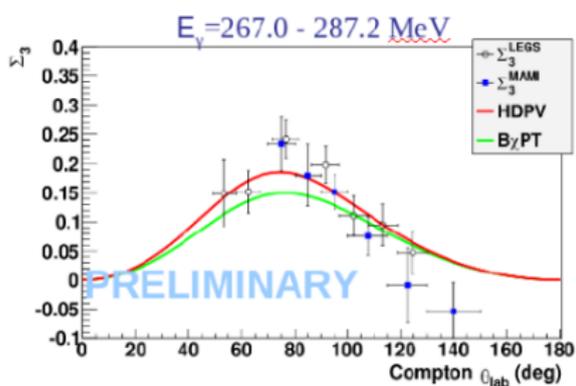
- Provides an excellent reaction for systematic checks and constraints. Due to the large σ (and clean reaction signal), π^0 production is an ideal reaction to perform systematic checks



- MM distribution for $E_\gamma = 273\text{--}303$ MeV, $\theta_{\gamma'} = 100\text{--}120$ degree (green)
- **Background contributions to MM:** accidental coincidences, (cyan) carbon/cryostat contributions (blue), reconstructed π_0 background where one decay γ escapes setup in: TAPS downstream hole (red) and CB upstream hole (magenta)
- **Right: Fully-subtracted MM spectrum with simulated Compton peak and conservative MM < 940 MeV cut is applied to exclude neutral pion production,**



- **New results!** Physical Review Letters 114, 112501 (2015), arXiv:1408.1576 [nucl-ex]
- **First measurement of a double-spin Compton scattering asymmetry on the nucleon.** Curves are from DR calculation of Pasquini et al., making use of constraints on " $\gamma_0, \gamma_\pi, \alpha_{E1} + \beta_{M1}, \alpha_{E1} - \beta_{M1}$ " (allowed to vary within experimental errors). Checks were done with $B_{\chi PT}$ calculation of Lensky & Pascalutsa.



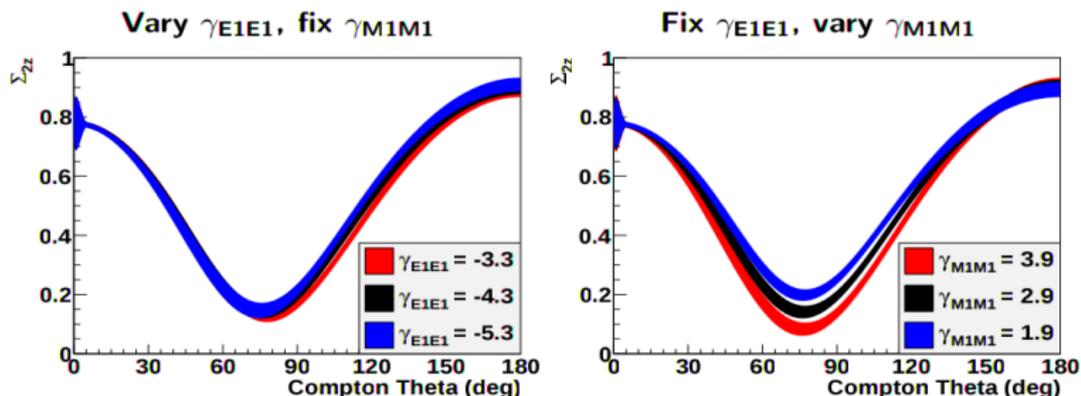
- **New MAMI and Older LEGS measurements along with two theoretical curves using their preferred polarizabilities**
- **Simulation of neutral pion photoproduction in Liquid hydrogen target matches background of the distribution quite well**

Preliminary Combined Spin Polarizabilities

	HDPV	B χ PT	Σ_{2x} and Σ_3^{LEGS}	Σ_{2x} and Σ_3^{MAMI}
γ_{E1E1}	-4.3	-3.3	-3.5 ± 1.2	-5.0 ± 1.5
γ_{M1M1}	2.9	3.0	3.16 ± 0.85	3.13 ± 0.88
γ_{E1M2}	-0.0	0.2	-0.7 ± 1.2	1.7 ± 1.7
γ_{M1E2}	2.2	1.1	1.99 ± 0.29	1.26 ± 0.43
γ_0	-0.8	-1.0	-1.03 ± 0.18	-1.00 ± 0.18
γ_π	9.4	7.2	9.3 ± 1.6	7.8 ± 1.8
$\alpha + \beta$			14.0 ± 0.4	13.8 ± 0.4
$\alpha - \beta$			7.4 ± 0.9	6.6 ± 1.7
χ^2/df			1.05	1.25

- Dispersion relation fits to Σ_{2x} along with either Σ_3^{MAMI} or Σ_3^{LEGS} (Note: Pion pole contribution has been subtracted)

Σ_{ZZ} – Estimated Experimental Precision (D. Paudyal, A. Rajabi)



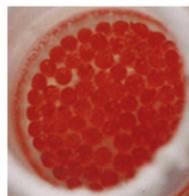
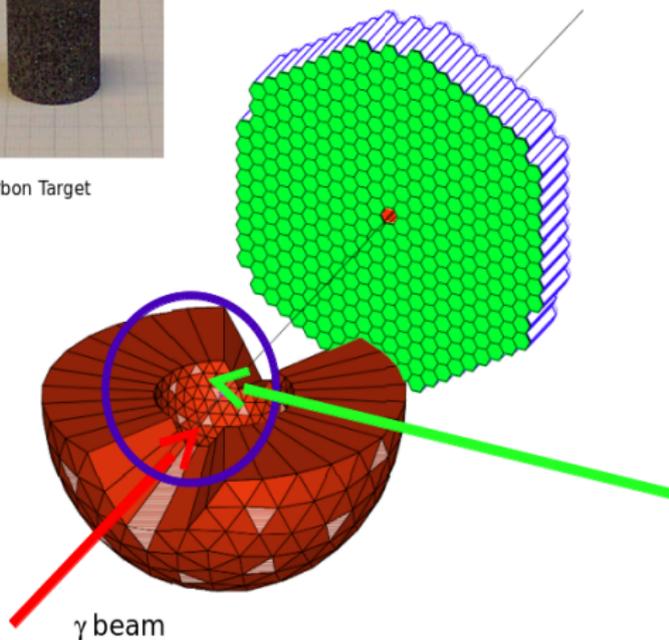
- To get a rough idea of the sensitivities, use a basis of γ_{E1E1} , γ_{M1M1} , γ_0 and γ_π , Produce event rates for nominal values of the SPs, using a dispersion theory calculation.
- Hold either γ_{E1E1} or γ_{M1M1} fixed, and perturb the other by a fixed amount and allow γ_0 , γ_π , α_{E1} and β_{M1} to vary by their experimental errors.
- **The bands represent the spread about these values by varying γ_0 , γ_π , α_{E1} and β_{M1} by their errors.**

- Σ_{2x} has been measured for the first time and published in recent PRL.
- **To further reduce the Σ_{2x} error bars, we have planned to acquire more data in January/February 2016.**
 - Σ_3 data analysis has been completed and planned for publication.
- **Planned to finish data analysis and have Compton double polarization asymmetries results Σ_{2z} before the end of 2016.**
- **Extract proton spin polarizabilities combining the Σ_{2z} results from first round of Σ_{2z} data taken in 2014 and second round of data taken in 2015**

Back Up slides

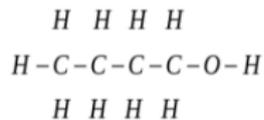


Carbon Target



Frozen Butanol

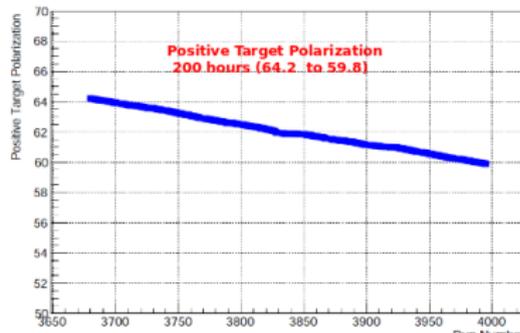
- 2 cm long and 2 cm diameter



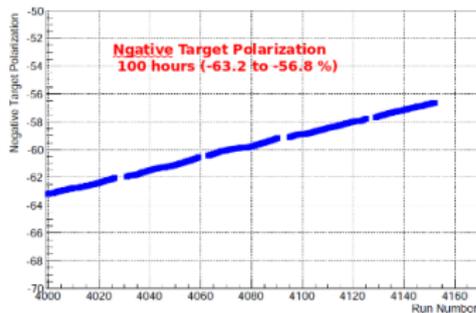


- **DNP**: Cool target to 0.2 K, use 2.5 Tesla magnet to align electron spins, pump 70 GHz microwaves, causing spin-flips between the electrons and protons.
- Cool target to 0.025 K, 'freezing' proton spins in place, remove polarizing magnet, energize 0.6 Tesla 'holding' coil in the cryostat to maintain the polarization, Relaxation times > 1000 hours, Polarizations up to 90%.

First Round of 2014 butanol data



First Round of 2014 butanol data



π^0 Beam Asymmetry: 288.3 ± 3.9 MeV

