

Update on the Kaon LT Experiment

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Hall A/C Summer Collaboration Meeting
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Outline

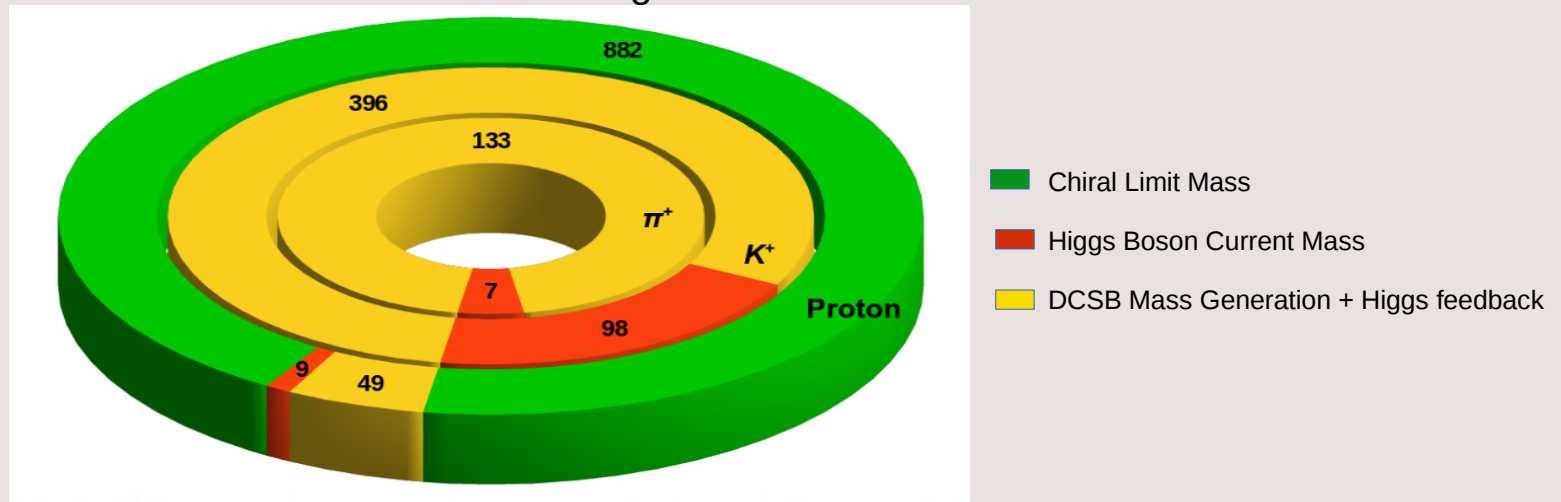
- Scientific overview
- Kaon LT experiment and its studies
- Rosenbluth separation overview
- Analysis updates
- Update on the Kaon LT PAC49
- Future perspectives

Dynamics of gluons in QCD

● **How does the mass of the nucleon arise?**

● **How does the spin of the nucleon arise?**

Hadron mass budget



● **Higgs mechanism is not sufficient !!!**

● **Dynamical Chiral Symmetry Breaking (DCSB)** is expected to provide most of the hadron mass.

● The **pion** and **kaon** are very important candidates to utilize in the **DCSB** study.

- π the lightest quark system. Responsible for the long range character of the strong interaction.
- K structure is involved with the strangeness.
- Both are connected to the Goldstone modes of **DCSB**.

Kaon LT experiment (E12-09-011)

- Kaon LT experiment is carried out in Hall C at Jefferson Lab over fall 2018 and spring 2019.

- Exclusive reaction of the experiment.
- Studies for the separation of cross-sections (σ_L , σ_T , σ_{LT} & σ_{TT}).
- Further studies for the “soft” and “hard” QCD factorization.
- Attempt to extract the kaon form factor above the resonance region for the first time.

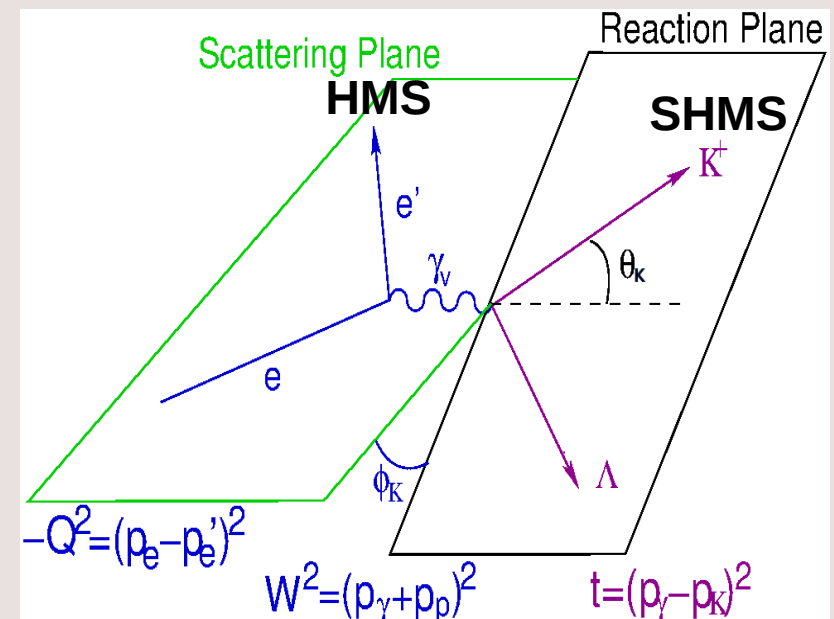
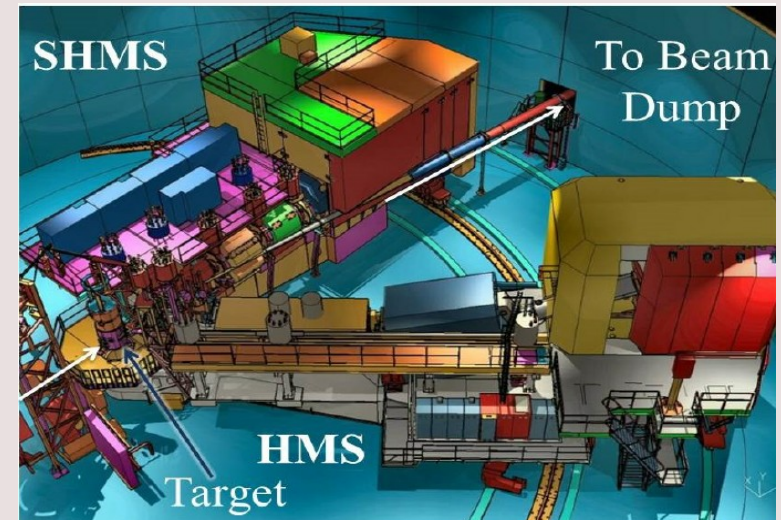
- The reaction system of the experiment is,

$$e + p \rightarrow e' + K^+ + \Lambda$$

$$M_\Lambda = 1115.68 \text{ MeV}^2/c^2$$

$$e + p \rightarrow e' + K^+ + \Sigma^0$$

$$M_{\Sigma^0} = 1192.64 \text{ MeV}^2/c^2.$$



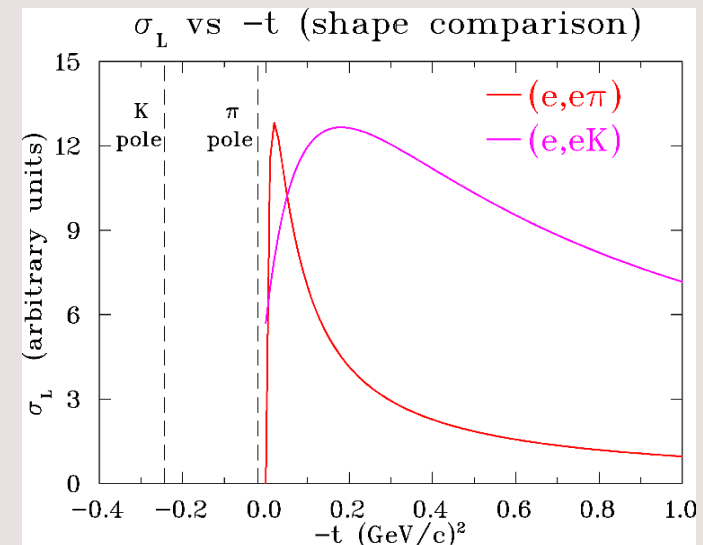
$e(p, e'K^+)\Lambda$ or Σ^0 L/T separation studies (search for the K^+ pole)

• The $-t$ dependence studies at constant Q^2 are important to search the kaon pole.

- The kaon pole is further away from the kinematically allowed region.
- It dominates σ_L at smallest $-t$. This test is very crucial to the kaon form factor and has never been tested before.

• Challenges:

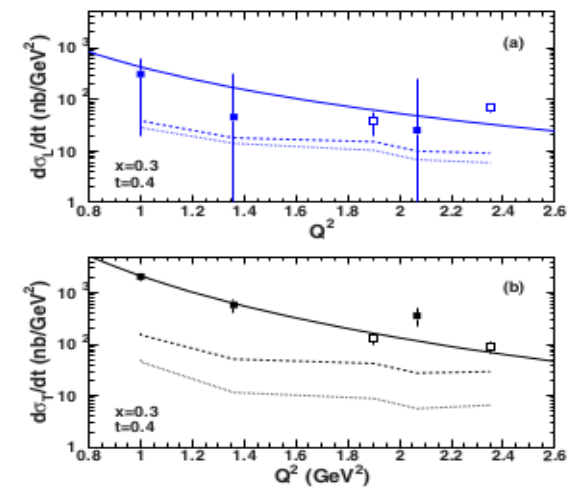
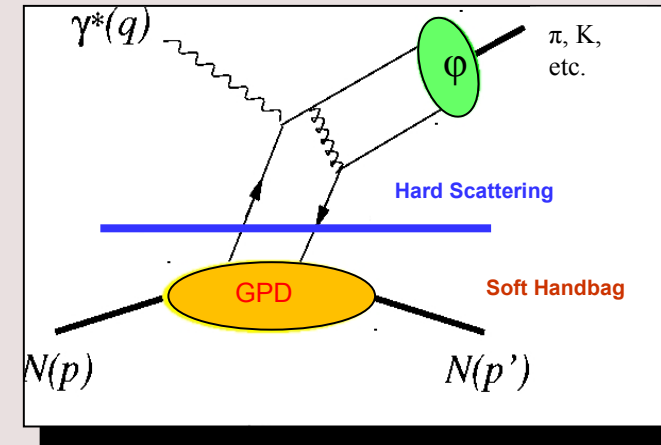
- The L/T separation studies require a deep understanding of each detector used in the experiment.
- The kaon PID is very challenging because the pion background is much higher in our data.
- Systematic and statistical uncertainties are also challenges to our analysis.



$e(p, e'K^+)\Lambda$ or Σ^0 L/T separation studies (QCD factorization)

● The Q^2 dependence studies at constant x_B are important to the QCD factorization test.

- Testing the factorization theorem and understand the dynamic effects in Q^2 and $-t$ kinematics.
- Understand the non-perturbative contributions in the experimentally accessible kinematics.
- One of the predictions, $\sigma_L \sim 1/Q^6$ and $\sigma_T \sim 1/Q^8$ at fixed x_B .
- 6 GeV data analysis had made an effort to the QCD factorization test earlier.
- We have collected data for the scaling studies at $x_B = 0.40$ and 0.25 .





Kaon LT data collected

- The completed data that have been collected in Hall C at Jefferson Lab over fall 2018 and spring 2019 run plans are shown in the table below.

E_b (GeV)	Q^2 (GeV ² /c ²)	W (GeV)	x_B	$\epsilon_{\text{High}}/\epsilon_{\text{Low}}$	Study Type
10.6/8.2	5.5	3.02	0.40	0.53/0.18	scaling
10.6/8.2	4.4	2.74	0.40	0.72/0.48	scaling
10.6/8.2	3.0	3.14	0.25	0.67/0.39	both
10.6/6.2	3.0	2.32	0.40	0.88/0.57	scaling
10.6/6.2	2.115	2.95	0.21	0.79/0.25	both
4.9/3.8	0.5	2.40	0.09	0.70/0.45	FF

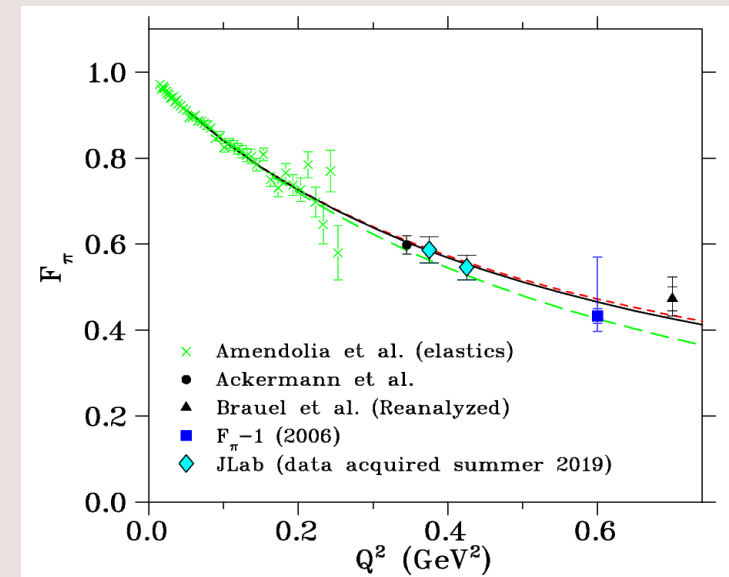


Pion LT (E12-19-006) data collected

- The Pion LT experimental data for low Q^2 have also been collected in Hall C at Jefferson Lab in summer 2019.

E_b (GeV)	Q^2 (GeV^2/c^2)	W (GeV)	x_B	ε
4.6/3.7/2.8	0.38	2.20	0.087	0.781/0.629/0.286
4.6/3.7/2.8	0.42	2.20	0.097	0.774/0.617/0.264

- We will use this data to compare the direct and indirect methods of extracting the pion form factor.
- The remaining experiment has been scheduled to be run from August to December and fall 2022.



Rosenbluth separation overview (simple version)

- Rosenbluth separation technique is one of the techniques that can be utilized to separate the cross-sections.

- Measure cross-section for at least two values of ϵ at fixed Q^2 , W and $-t$.
- Cross-section at two values of ϵ is then fitted to separate the cross-sections.

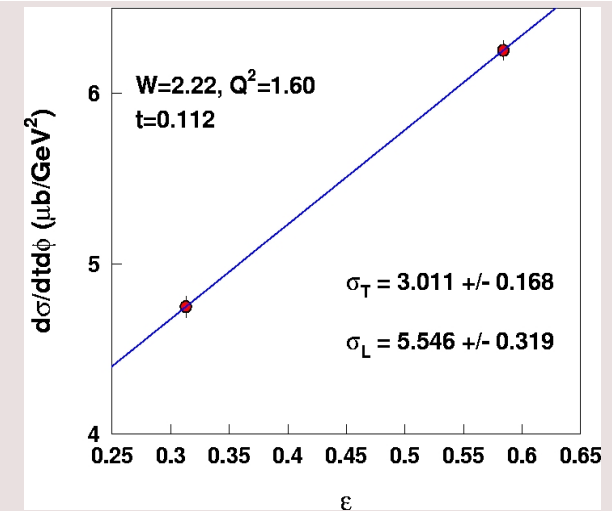
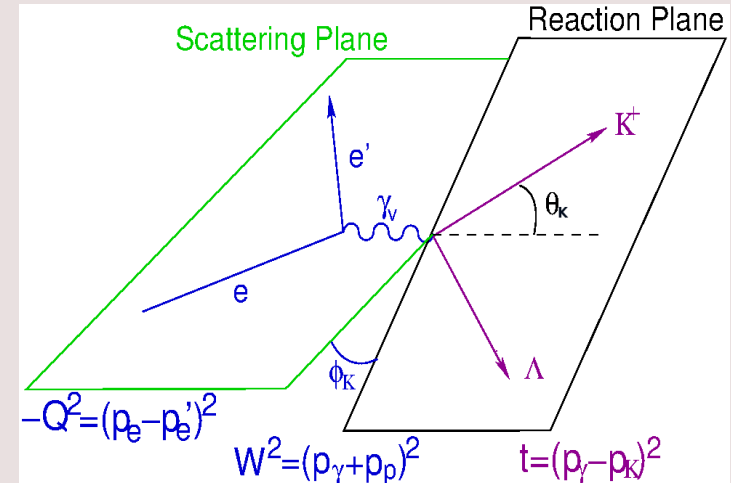
- In parallel kinematics, $\theta_K = 0$. (θ_K w.r.t \vec{q})

- Only the σ_L and σ_T terms contribute.
- The mathematical form is simple but requires uniform detector acceptance.

$$2\pi \frac{d^2\sigma}{dt d\phi} = \epsilon \frac{d\sigma_L}{dt} + \frac{d\sigma_T}{dt}$$

Virtual-photon polarization:

$$\epsilon = \left(1 + 2 \frac{(E_e - E_{e'})^2 + Q^2}{Q^2} \tan^2 \frac{\theta_{e'}}{2} \right)^{-1}$$



T. Horn, et al, PRL 97 (2006) 192001

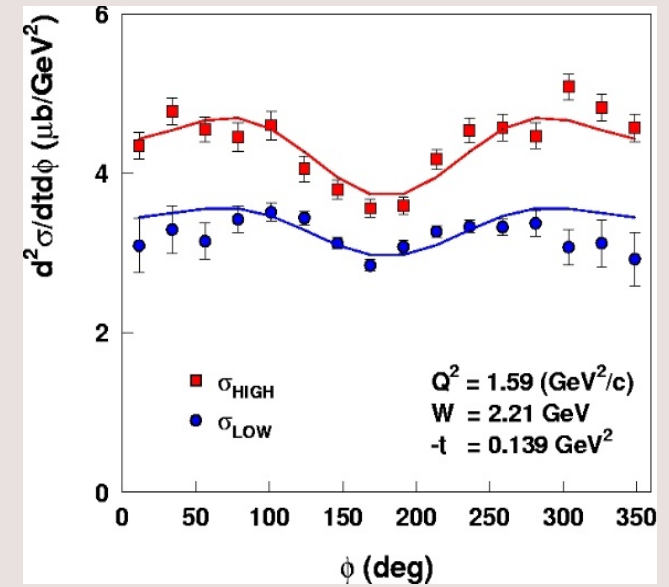
Rosenbluth separation overview (full version)

● In non-parallel kinematics, $\theta_K \neq 0$.

- Cross-section at two values of ε is simultaneously fitted with a four variable function to determine all of the cross-section terms.

$$2\pi \frac{d^2\sigma}{dtd\phi} = \varepsilon \frac{d\sigma_L}{dt} + \frac{d\sigma_T}{dt} + \sqrt{2\varepsilon(\varepsilon+1)} \frac{d\sigma_{LT}}{dt} \cos\phi + \varepsilon \frac{d\sigma_{TT}}{dt} \cos 2\phi$$

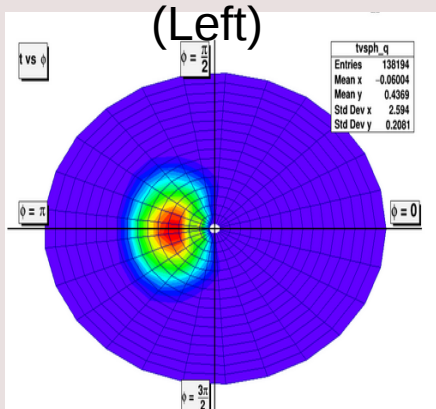
● We have collected the Kaon LT data for parallel and non-parallel kinematics settings.



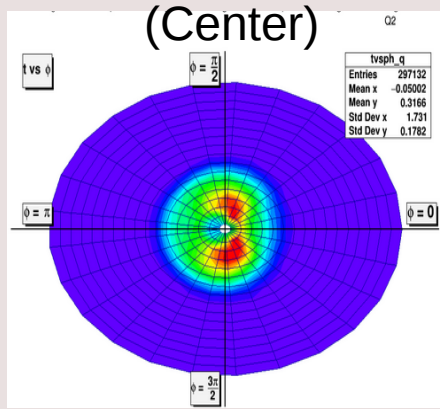
T. Horn, et al, PRL 97 (2006) 192001

$$E_b = 10.6 \text{ GeV} \quad Q^2 = 3.0 \text{ GeV}^2/\text{c}^2 \quad W = 3.14 \text{ GeV}$$

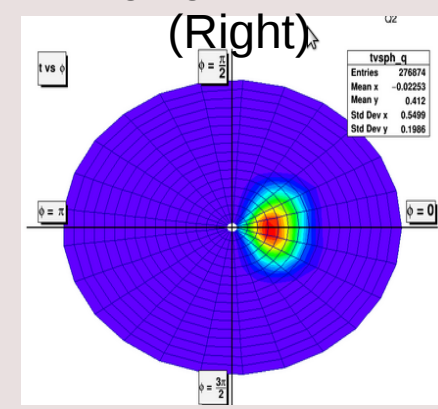
$$\theta_{\text{SHMS}} = 12.42^\circ$$



$$\theta_{\text{SHMS}} = 9.42^\circ$$



$$\theta_{\text{SHMS}} = 6.65^\circ$$

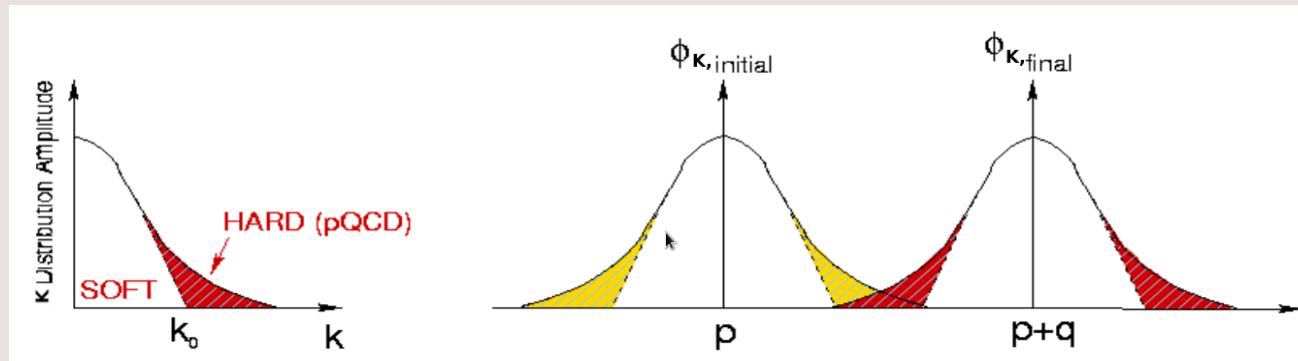


$-t$ is the radius

The meson wave function (form factor)

- The electromagnetic form factor is an important physical observable connected directly to the internal structure of mesons.
- In quantum field theory, the form factor is the overlap integral,

$$F_K(Q^2) = \int \phi_K^*(p) \phi_K(p+q) dp.$$



- The meson wave function can be separated into two parts.
 - ϕ_K^{soft} , ($k < k_0$) low momentum contributions only. Cannot be treated in pQCD.
 - ϕ_K^{hard} , hard tail can be treated in pQCD.
- If our data analysis indicates that the K^+ pole dominates σ_L then we will attempt to extract the kaon form factor.

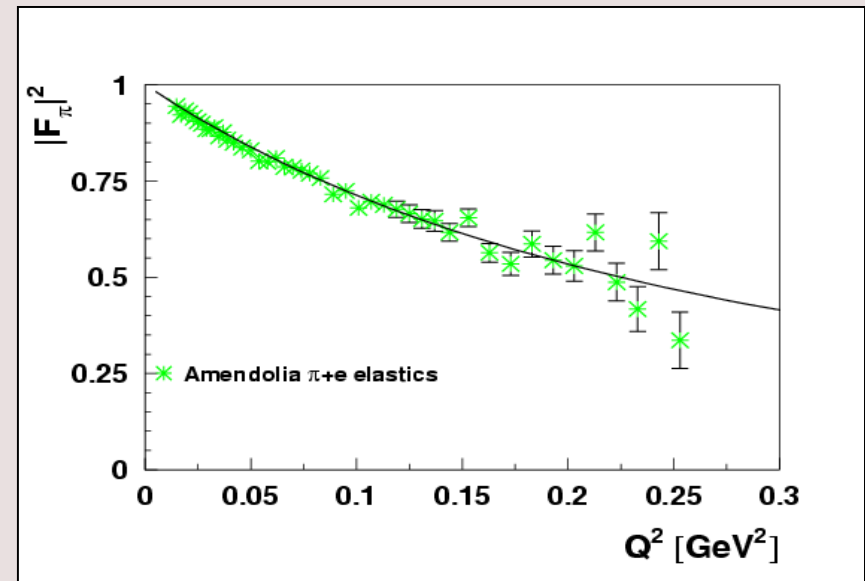
π and K form factors at low Q^2 (elastic scattering)

● $F_\pi(Q^2)$ and $F_K(Q^2)$ are known at low Q^2 , extracted by using π^- and K^- beams on LH_2 target.

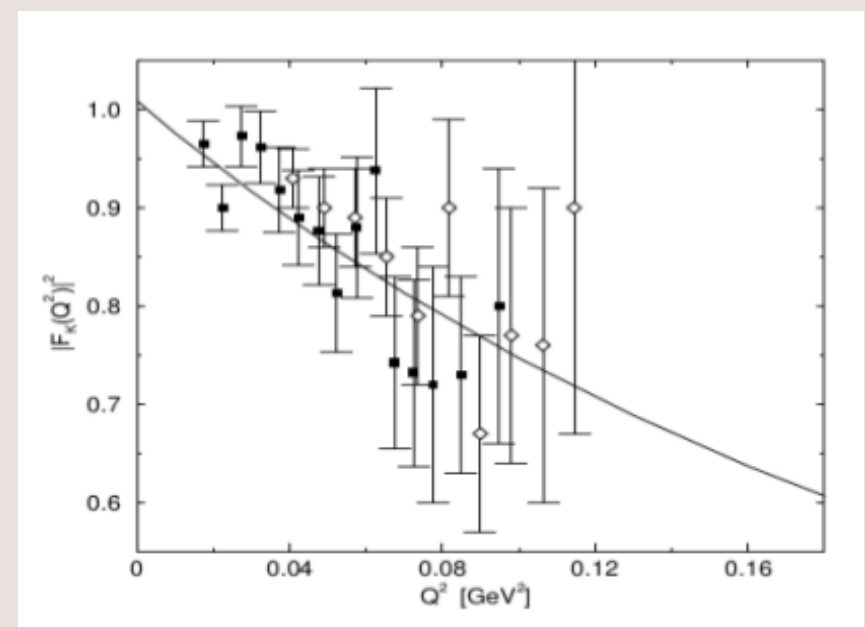
- π^- of 300 GeV, data were collected up to $Q^2 = 0.28 \text{ GeV}^2$.
- K^- of 250 GeV, data were collected up to $Q^2 = 0.13 \text{ GeV}^2$.

● These measurements were used to determine the charge radius of the π and K.

- The slope of the fitting function at $Q^2 = 0$ provides the radius.
- $\langle r^2 \rangle = -(6dF/dQ^2)_{Q^2=0}$.
- π charge radius $\langle r^2 \rangle^{1/2} = 0.657 \pm 0.012 \text{ fm}$.
- K charge radius $\langle r^2 \rangle^{1/2} = 0.340 \pm 0.050 \text{ fm}$.



[Amendolia, et al., NP B277 (1986) 168]



[Amendolia, et al., PL B178 (1986) 435]

π and K form factors at higher Q^2

At higher Q^2 , the direct scattering is not achievable.

- π and K beams of much higher energy which is not possible at the current experimental facilities.

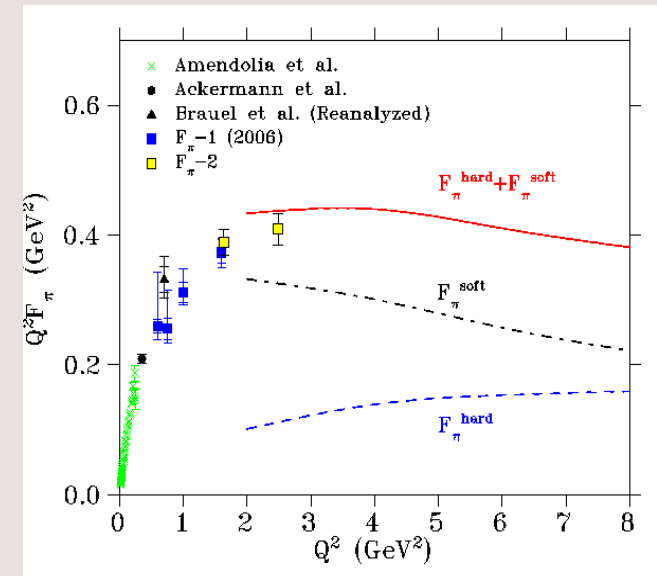
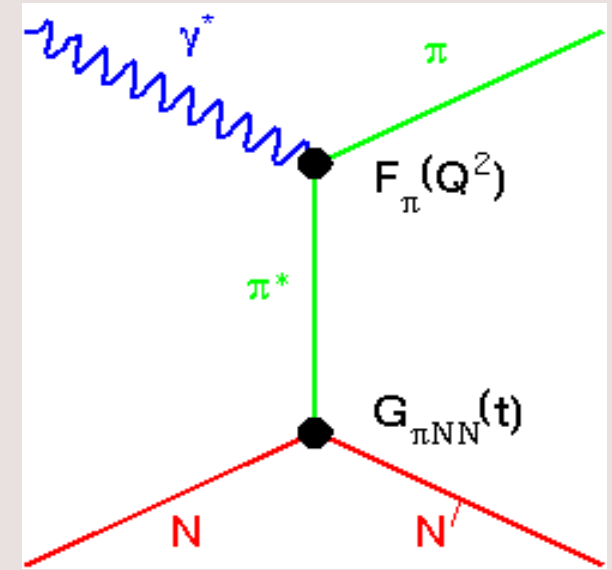
To access the form factors at higher Q^2 , must employ an alternative method.

- “Virtual cloud” of the π and K inside the proton makes the measurements possible.
- This attempt has been made at 6 GeV era and a few other experimental facilities.

In Born term model, the form factor appears as,

$$\frac{d\sigma_L}{dt} \propto \frac{-tQ^2}{(t - m_K^2)} g_{K\Lambda N}^2(t) F_K^2(Q^2, t)$$

- Indirect determination of the form factor is model dependent.

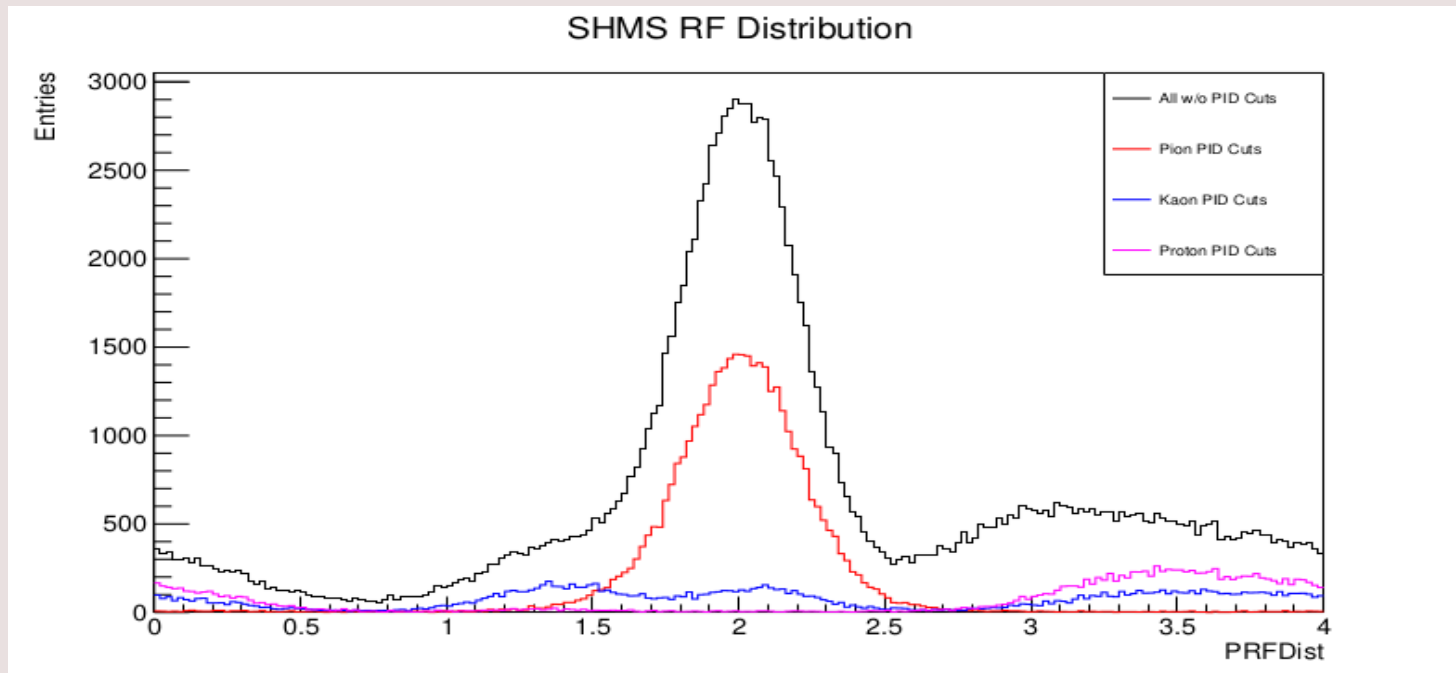


G.M. Huber et al., PRC 78 (2008) 045203.

Analysis update (RFTIME & missing mass branches)

● A new branch of the RFTIME variable has been added to hcana.

- In some settings the RFTIME variable is useful for PID.
- New variable of the RFTIME is calculated as,
- $\text{RFTIME} = \text{mod}((\text{BunchSpacing})(\text{P_RF_tdcTime} - \text{P_hod_fpHitsTime} + \text{RF_Offset}))$.
- **RF_Offset** is just to recentre the resulting distribution.



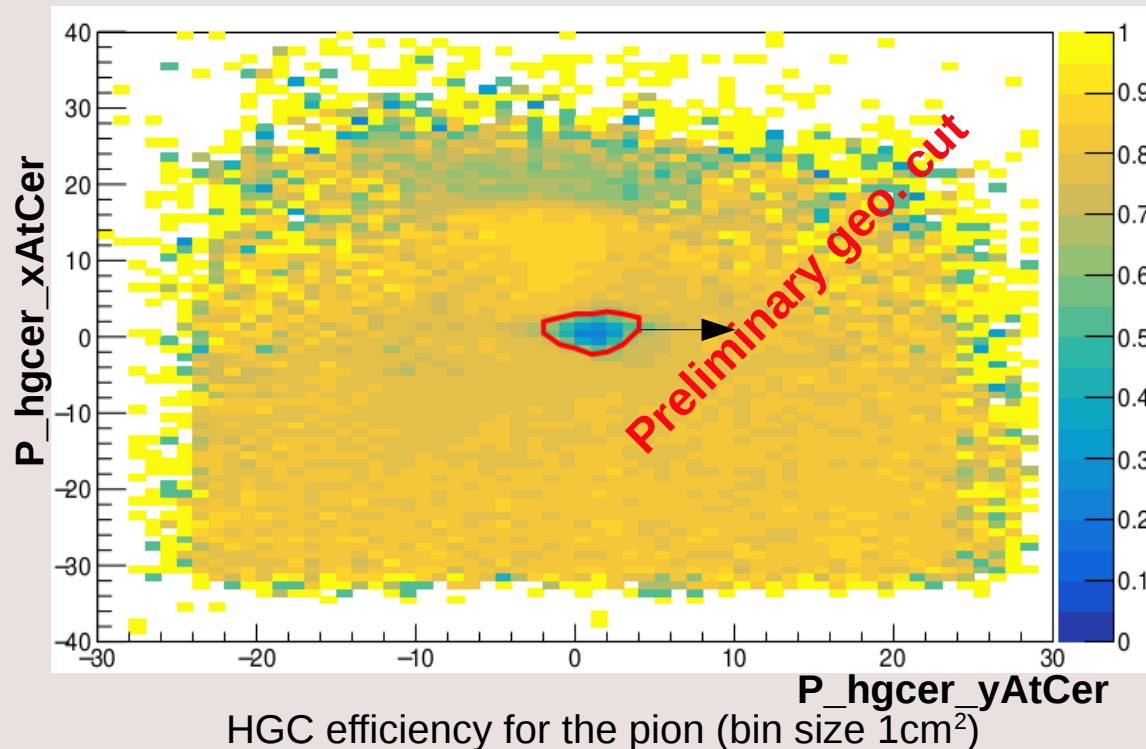
● New missing mass variables have also been added to hcana.

Stephen Kay (UofR), details of this work can be found at, [DocDB 1129-v2](#)

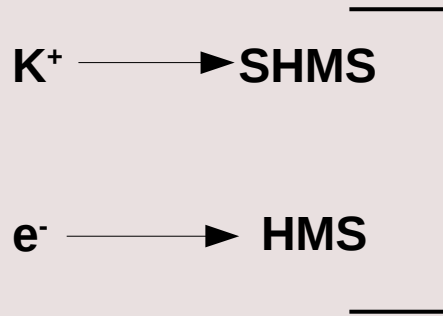
Analysis update (SHMS HGC efficiency & geometrical cut)

● Finalizing the HGC efficiency is challenging.

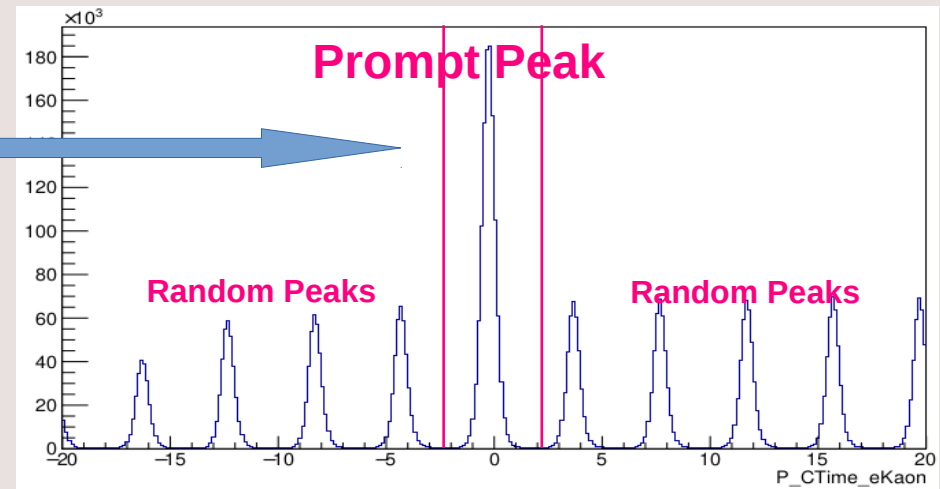
- The collection of Cherenkov lights in the HGC is not uniform.
- There is an inefficient region in the middle of the detector.
- Decided to cut off the very inefficient region.
- We are working to have **position-dependent bin-by-bin HGC efficiency** for regions outside the hole cut.



Analysis update (PID in SHMS)

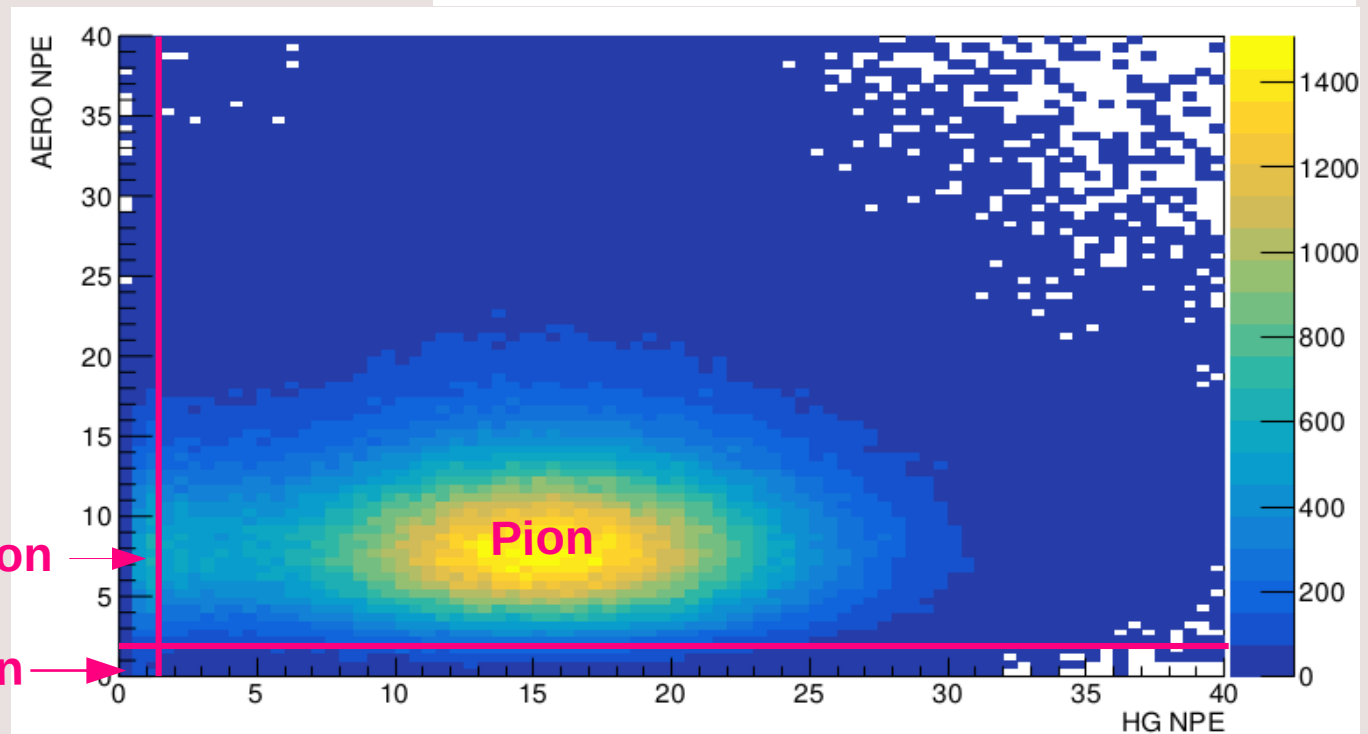


Real coincidence



- Aerogel refractive index **1.011** of current analysis setting.
- Heavy Gas refractive index **1.00139** at 1 atm.
- Y axis Aerogel Cherenkov light (NPE).
- X axis Heavy Gas Cherenkov light (NPE).

Kaon →
Proton →

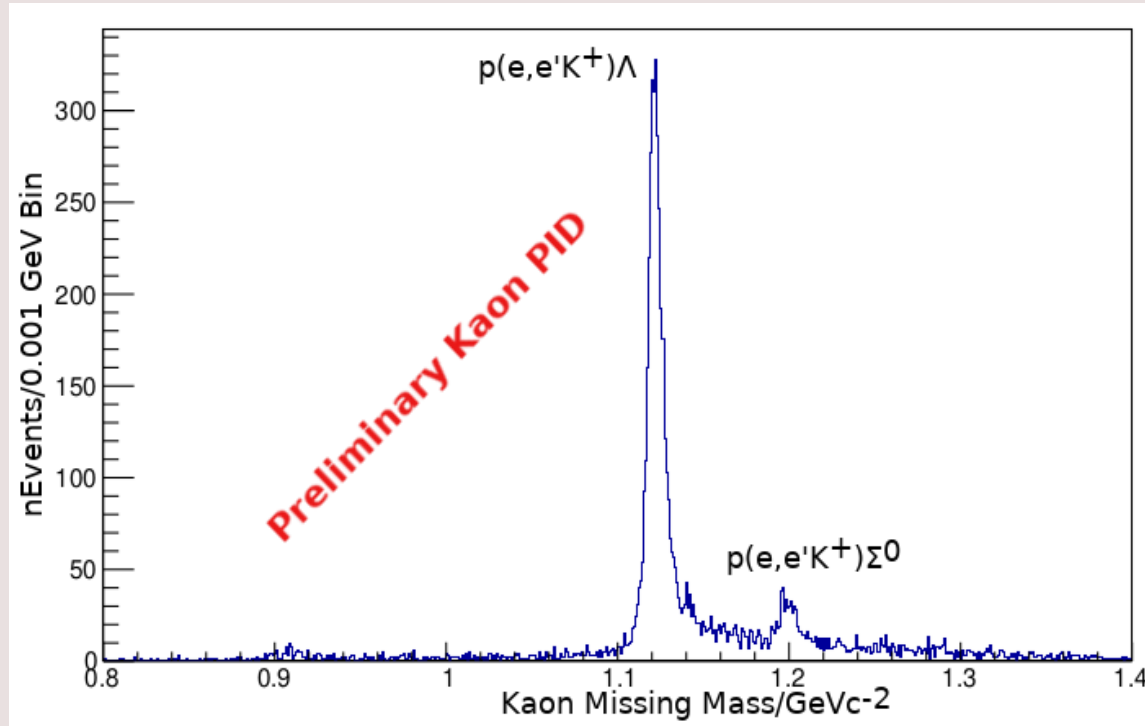


Analysis update (kaon PID)

● The preliminary kaon PID is shown in the kaon missing mass plot.

- The kaon sample looks clean but we need to optimize the kaon PID cuts.

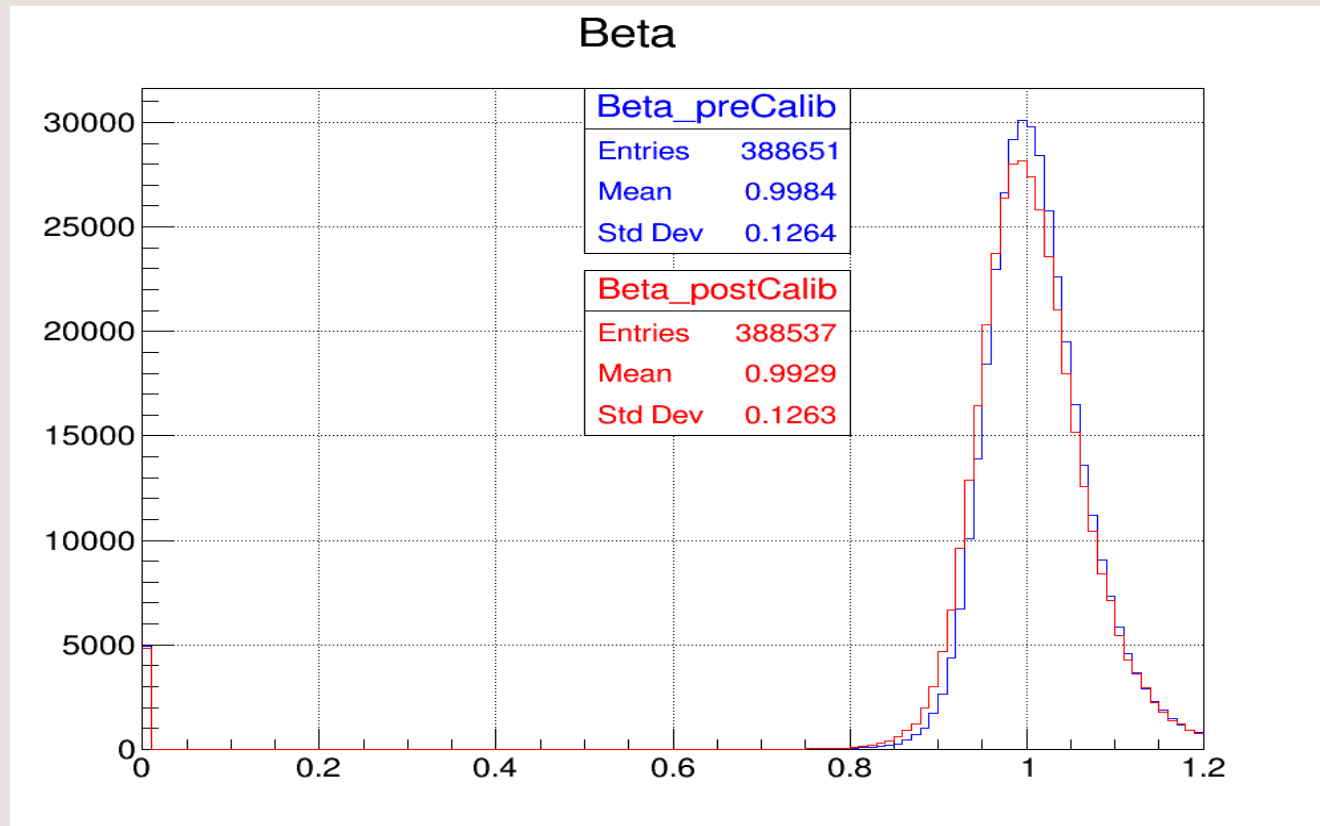
$$E_b = 6.19 \text{ GeV}, Q^2 = 3.0 \text{ GeV}^2/c^2 \quad W = 2.32 \text{ GeV}, P_{SHMS} = 3.48 \text{ GeV}/c$$



$$M_{miss} = [(E_b + m_p - E_{e'} - E_{K^+})^2 - (P_e - p_{e'} - p_{K^+})^2]^{1/2}$$

Analysis update (verifying the online SHMS Hodoscope calibration)

- The online Hodoscope calibration has been verified. Both the calibrations (online & offline) are almost equivalent.



Nathan Heinrich (UofR)

Analysis update (scaler luminosity yields)

- The preliminary scaler luminosity study has been completed.

6.2 GeV

$P_{\text{HMS}} = -3.939$

$\theta_{\text{HMS}} = 13.00$

$P_{\text{SHMS}} = -3.939$

$\theta_{\text{SHMS}} = 11.00$

65, 50, 30, 15, 5 μA

7841, 7846, 7847, 7864, 7865

HMS: Electron

$\text{Cer} > 0.5$

$2.0 > \text{Cal} > 0.6$

SHMS: Electron

$\text{Hgcer} > 1.5$

$\text{Aero} > 1.5$

$0.6 < \text{Cal} \leq 2.0$

8.2 GeV

$P_{\text{HMS}} = 5.745$

$\theta_{\text{HMS}} = 12.97$

$P_{\text{SHMS}} = 5.745$

$\theta_{\text{SHMS}} = 9.51$

65, 50, 30, 15, 5 μA

7948, 7949, 7950, 7951, 7952

HMS: Hadron

$\text{Cer} < 0.5$

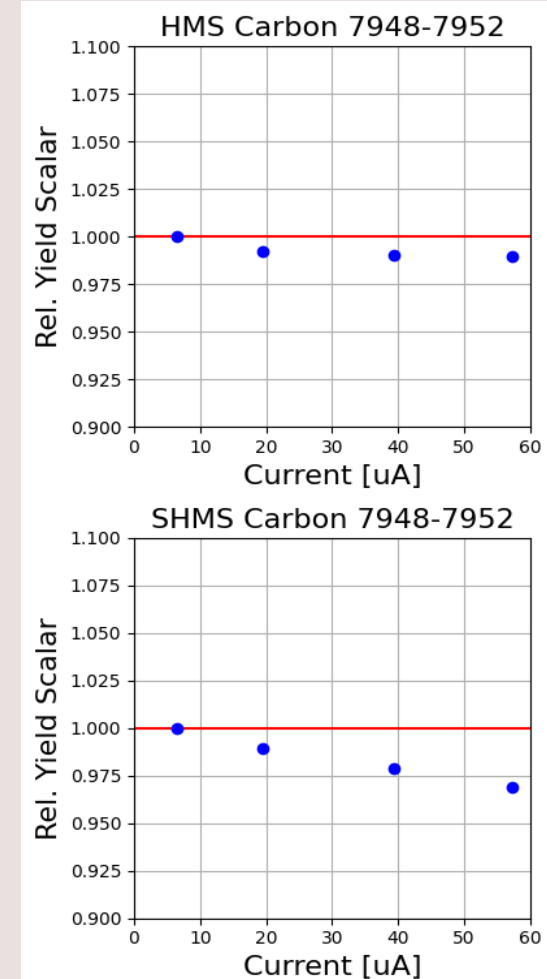
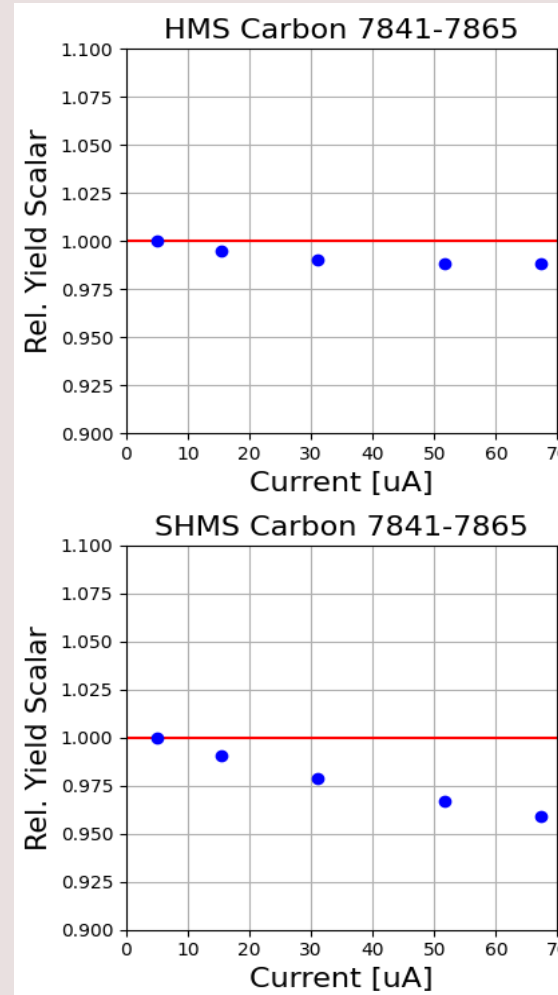
$0.0 < \text{Cal} \leq 0.6$

SHMS: Proton

$\text{Hgcer} < 1.5$

$\text{Aero} < 1.5$

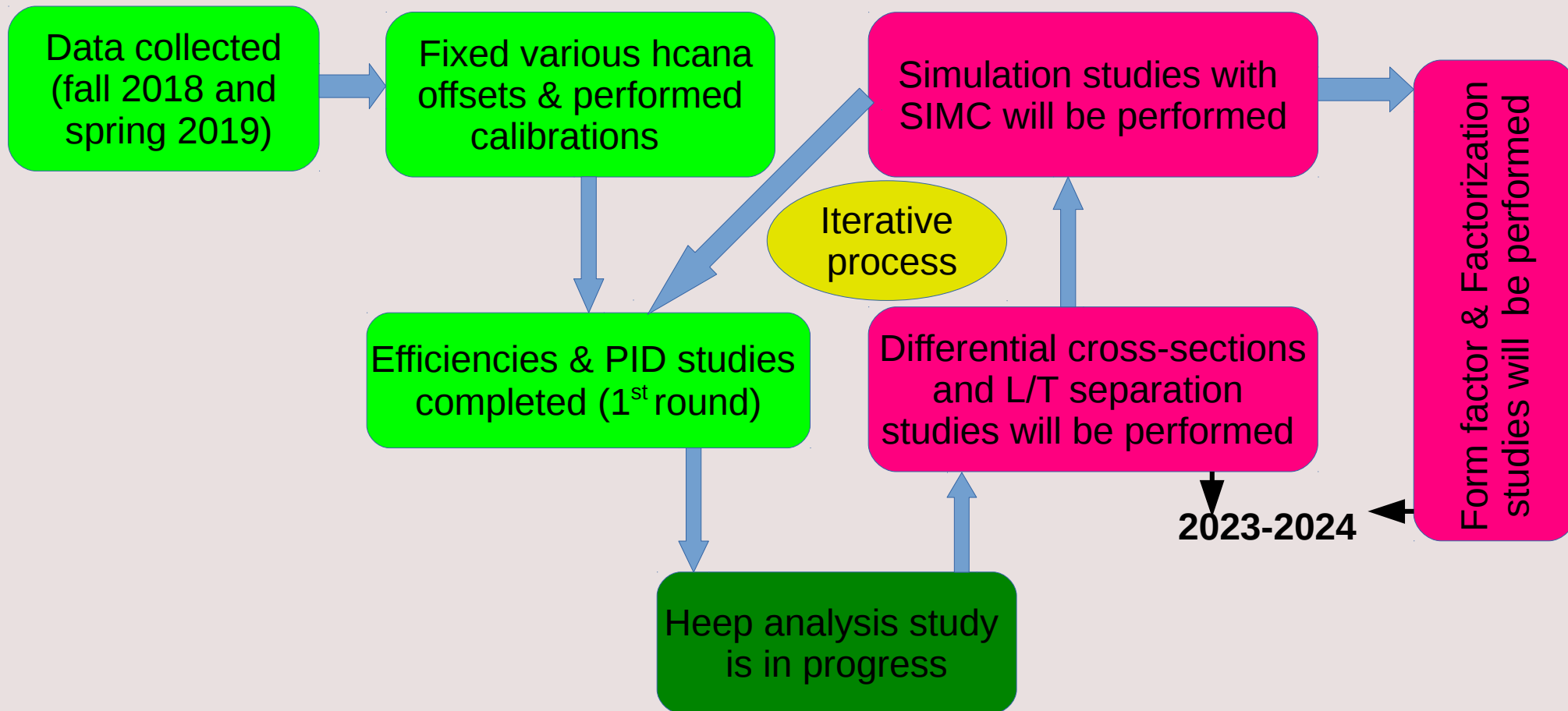
$0 < \text{Cal} \leq 0.6$



Richard Trotta (CUA)

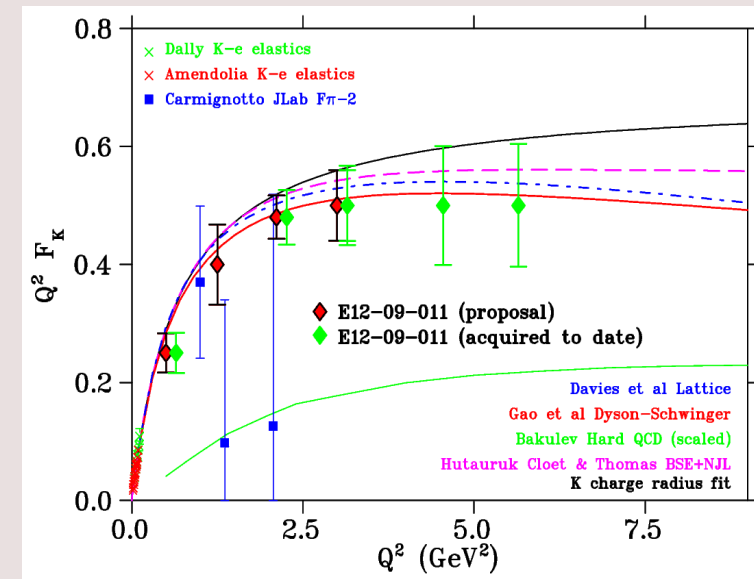
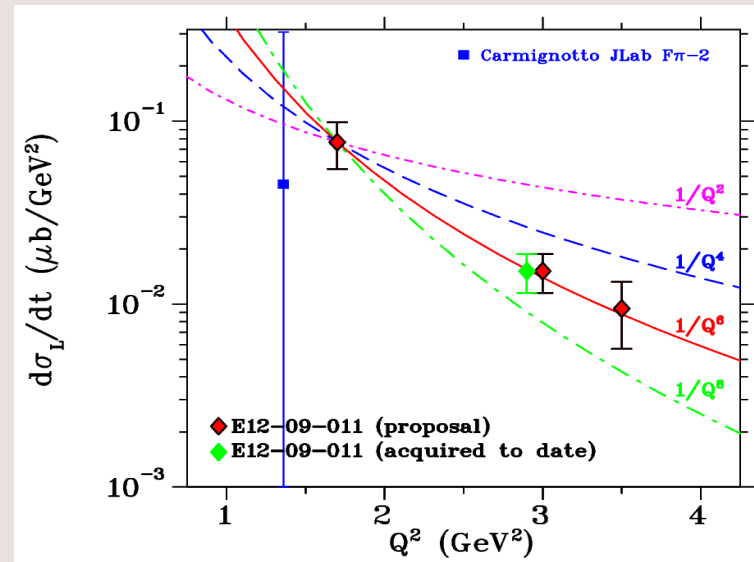
Analysis update (analysis flow chart)

- The data analysis of the Kaon LT experiment is an iterative process.



Update on the KaonLT PAC49

- Overall ~60% of our approved data were acquired in fall 2018 and spring 2019 run plans.
 - Q^2 scan at $x_B = 0.40$.
 - And $F_K(Q^2)$ points.
- The statistics at the higher beam energies are systematically lower except for one of the settings.
 - This increases the originally projected uncertainties and makes difficult to extract the science for the Q^2 scan at $x_B = 0.25$.
- A KaonLT PAC49 jeopardy document has been submitted, requesting a total of 18 days beam time.
 - Q^2 scan at $x_B = 0.25$.
 - And $F_K(Q^2)$ points.
- Spokespersons:
 - T. Horn, G. Huber and P. Markowitz



Summary and future perspectives

- The Kaon LT experiment is for the high precision L/T/LT/TT separation studies. The data analysis is intricate and the international journal publications are expected to be published in 2-3 years.
- I will start the $p(e, e'K^+)\Lambda$ L/T/LT/TT separation cross-sections studies at $Q^2 = 0.5 \text{ GeV}^2/c^2$ after completing the ongoing studies.
- I will also start the $p(e, e'\pi^+)n$ L/T/LT/TT separation cross-sections studies at $Q^2 = 0.5 \text{ GeV}^2/c^2$ from the Kaon LT experiment and at $Q^2 = 0.38 \text{ and } 0.42 \text{ GeV}^2/c^2$ from the Pion LT experiment.
- Other studies such as simulations (SIMC) will also be carried out with the physics analysis.
- Other students are also preparing to start the main physics analysis at higher Q^2 .



SAPIN-2021-00026

Thank You!

