



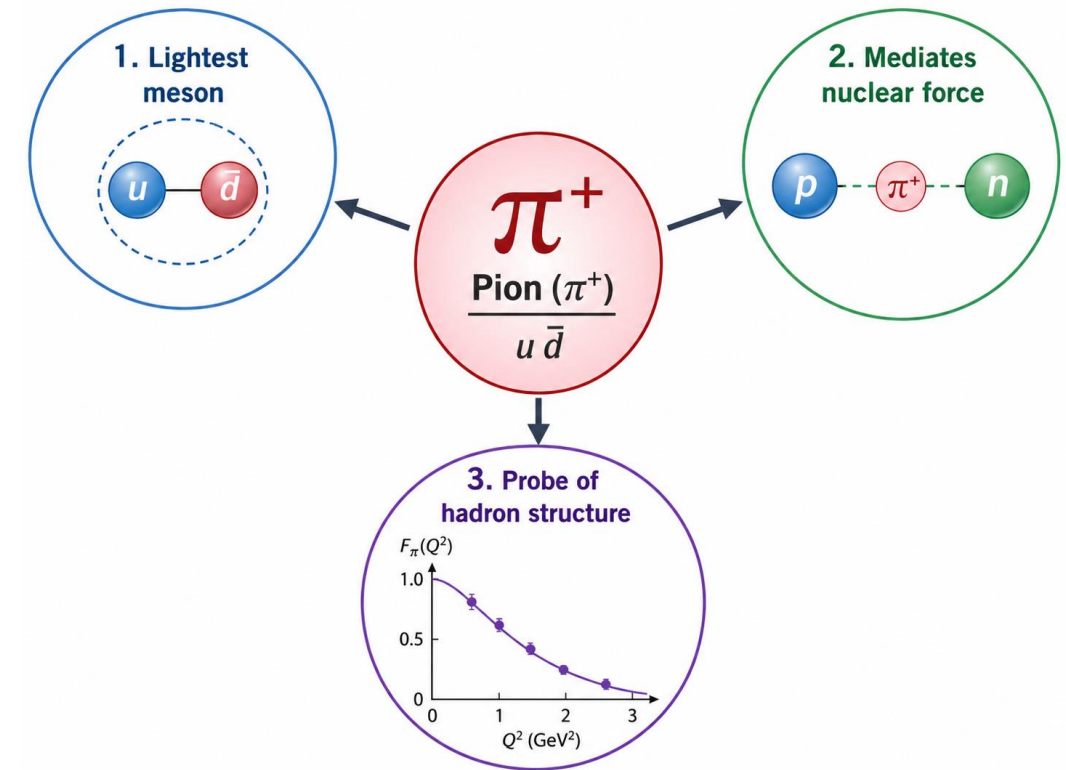
Exploring Hadronic Structure: Precision Rosenbluth Separation for the Pion Form Factor at Jefferson Lab

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Scientific Motivation

- ❑ The interaction of quarks and gluons is successfully described by **Quantum Chromodynamics (QCD)**.
- ❑ **But unable to construct the quantitative description of hadrons in terms of the underlying constituents, quarks and gluons.**
- ❑ **Pion** is the lightest meson and provides an ideal testing ground for our understanding of the bound $q\bar{q}$ hadronic system.
- ❑ **Pions mediate the residual strong force that helps bind nucleons inside the nucleus.**
- ❑ **Form factor ($F(Q^2)$)** is an important observable that can be studied to understand the internal structure of hadrons by describing the transverse spatial position of partons within hadrons.
- ❑ **Measuring the pion form factor at various Q^2 (up to 8.5GeV^2) checks the validity of QCD-based theories, including the transition region between perturbative and non-perturbative approaches.**

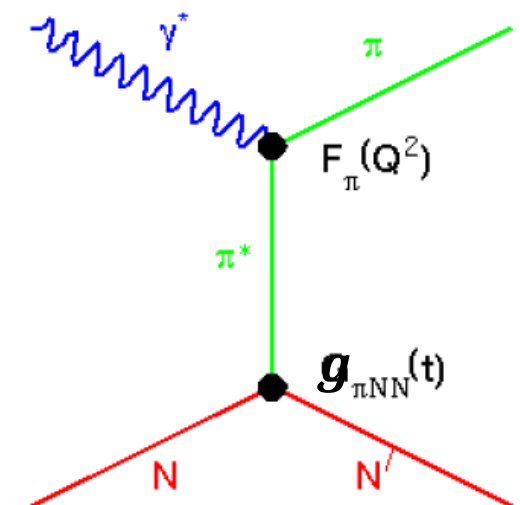
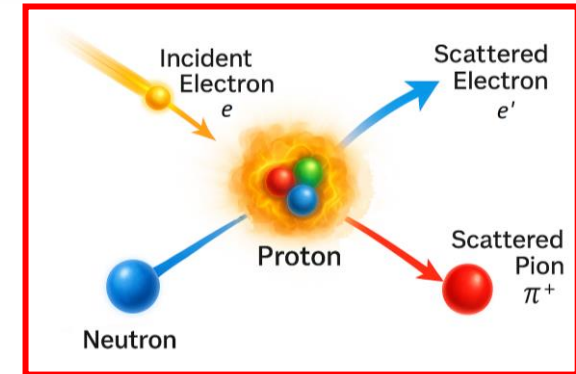


Pion Form Factor Measurement

- ❑ Above $Q^2 > 0.3 \text{ GeV}^2$, F_π is measured indirectly using the “pion cloud” of the proton via pion electroproduction $p(e, e' \pi^+) n$
- ❑ Indirect measurement – Form factor extraction requires a model.
- ❑ Need to extract σ_L from the total cross-section.
- ❑ As an illustration of how σ_L connects to $F_\pi^2(Q^2, t)$, we consider a simple **Born Term Model**;

$$\frac{d\sigma_L}{dt} \propto \frac{-tQ^2}{(t - m_\pi^2)} g_{\pi NN}^2(t) F_\pi^2(Q^2, t)$$

- ❑ In reality, we use Regge base model such as VGL, CKY and PKT Models for $F_\pi^2(Q^2, t)$ extraction.



Rosenbluth Separation Technique

- Rosenbluth separation required to isolate σ_L for L/T separation.

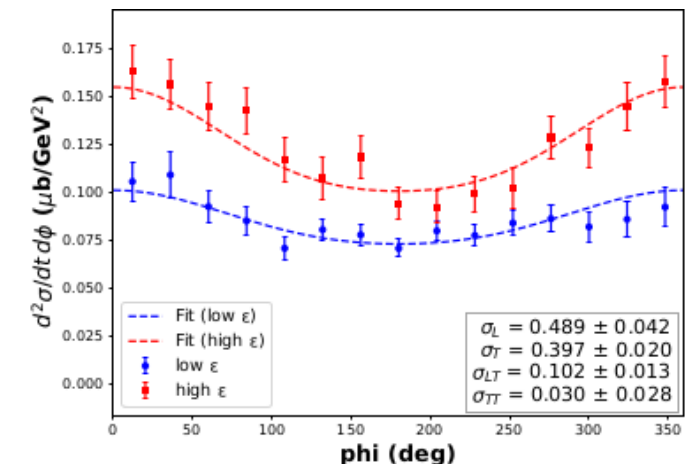
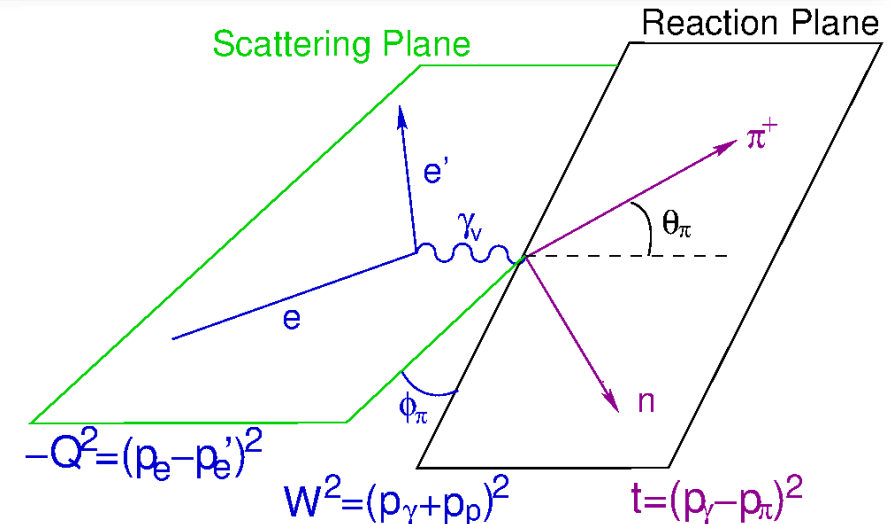
- The physical cross-section for the electroproduction process is given by;

$$2\pi \frac{d^2\sigma}{dt d\phi} = \epsilon \frac{d\sigma_L}{dt} + \frac{d\sigma_T}{dt} + \sqrt{2\epsilon(\epsilon+1)} \frac{d\sigma_{LT}}{dt} \cos \phi + \epsilon \frac{d\sigma_{TT}}{dt} \cos 2\phi$$

- Here “ ϵ ” is polarization of virtual photon.

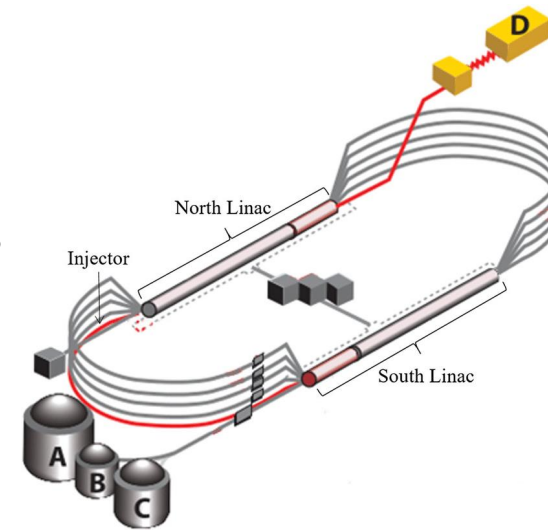
$$\epsilon = \left[1 + 2 \frac{(E_e - E_{e'})^2 + Q^2}{Q^2} \cdot \tan^2 \frac{\theta_{e'}}{2} \right]^{-1}$$

- Perform two scattering measurements with different beam energies “ E_e ” to vary “ ϵ ” and separate different cross-section terms.
- Careful attention must be paid to systematic studies such as spectrometer acceptance, kinematics, efficiencies, etc.



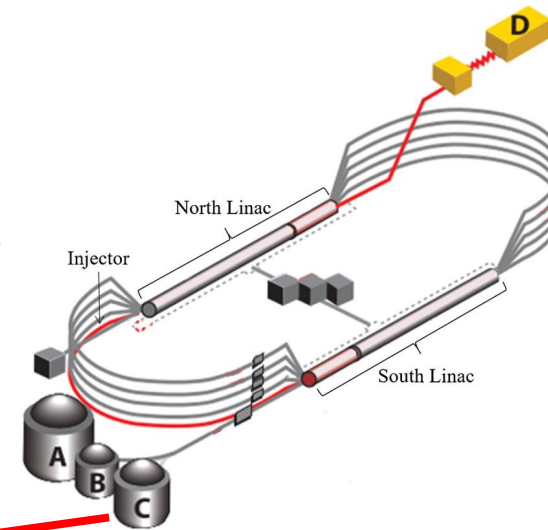
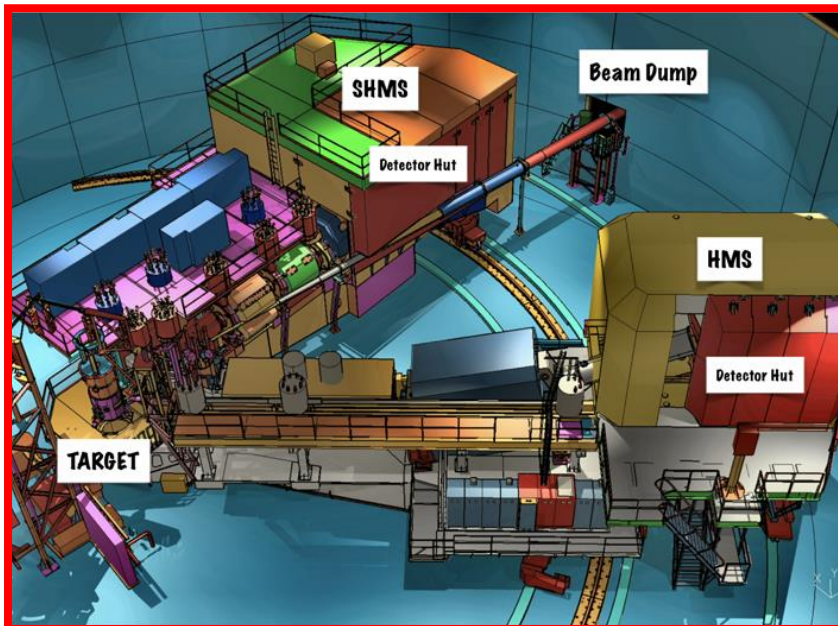
Pion-LT Experiment at JLab

- Located at Newport News, Virginia, USA.
- 2 Superconducting LINACs configured as “Racetrack”.
- Continuous Electron Beam Accelerator Facility (CEBAF).
- Capable of 12GeV beam up to $200\mu\text{A}$.
- Four experimental halls, all running unique experiments and capable of running simultaneously.



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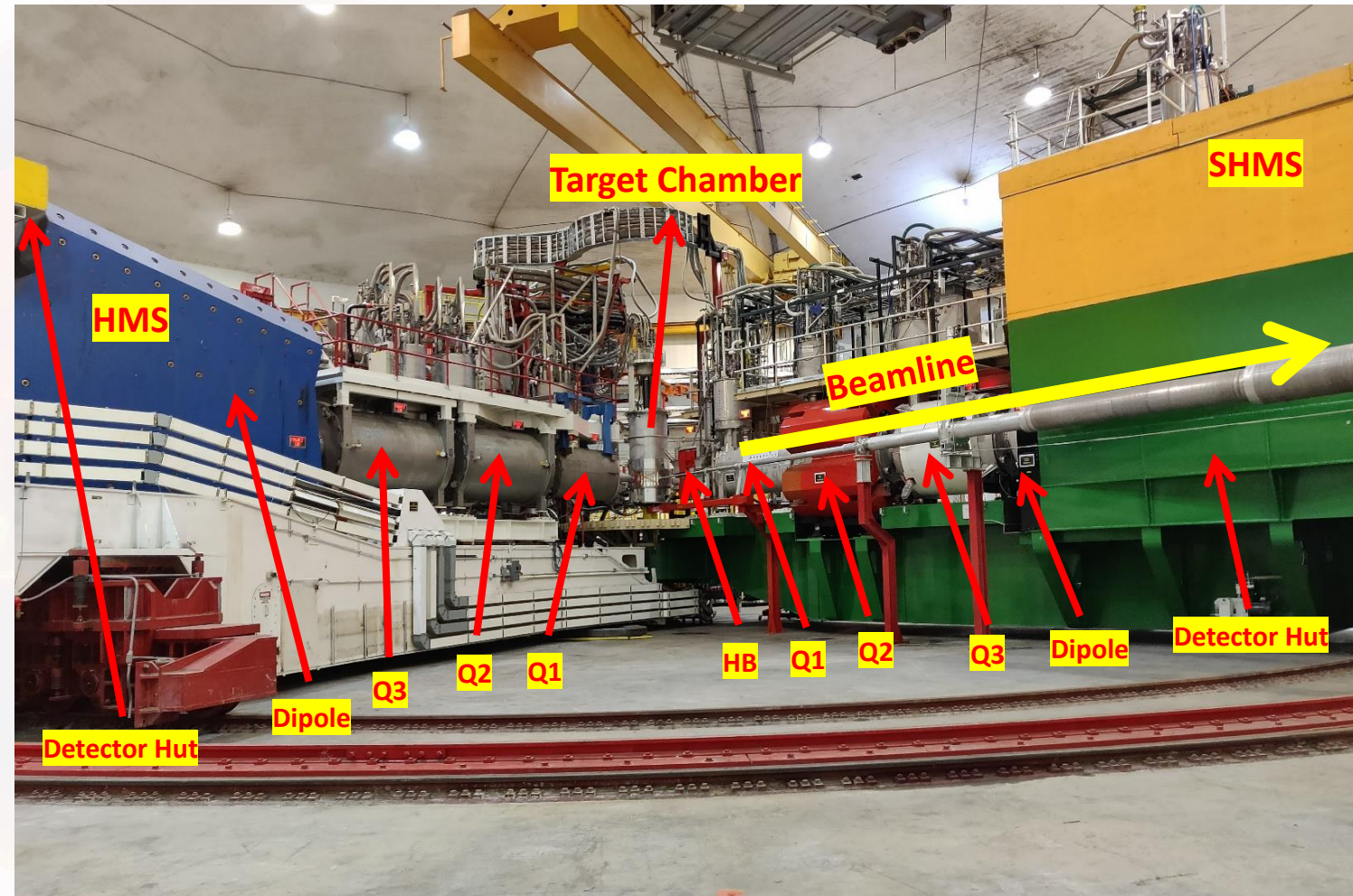
- Specially designed to measure precise cross-sections and form factors for mesons.
- Hall C has a target and two detector arms.
- Target can be Liquid H_2 , Liquid D_2 or solid targets.
- Two detectors are;
 - High Momentum Spectrometer (HMS)
 - Super High Momentum Spectrometer (SHMS)

Hall C Spectrometers

- Both spectrometers are movable and have sets of quadrupole (Q), dipole (D) superconducting magnets and a Detector Hut.
- Both spectrometers are rotatable and contain similar detector packages.

Spec	Angle range (deg)	Momentum range
HMS	10.5 – 90	0.5 – 7 GeV
SHMS	5.5 – 40	0.5 – 11 GeV

- Took coincidence data; simultaneously detected **electrons in HMS** and **pions in the SHMS**

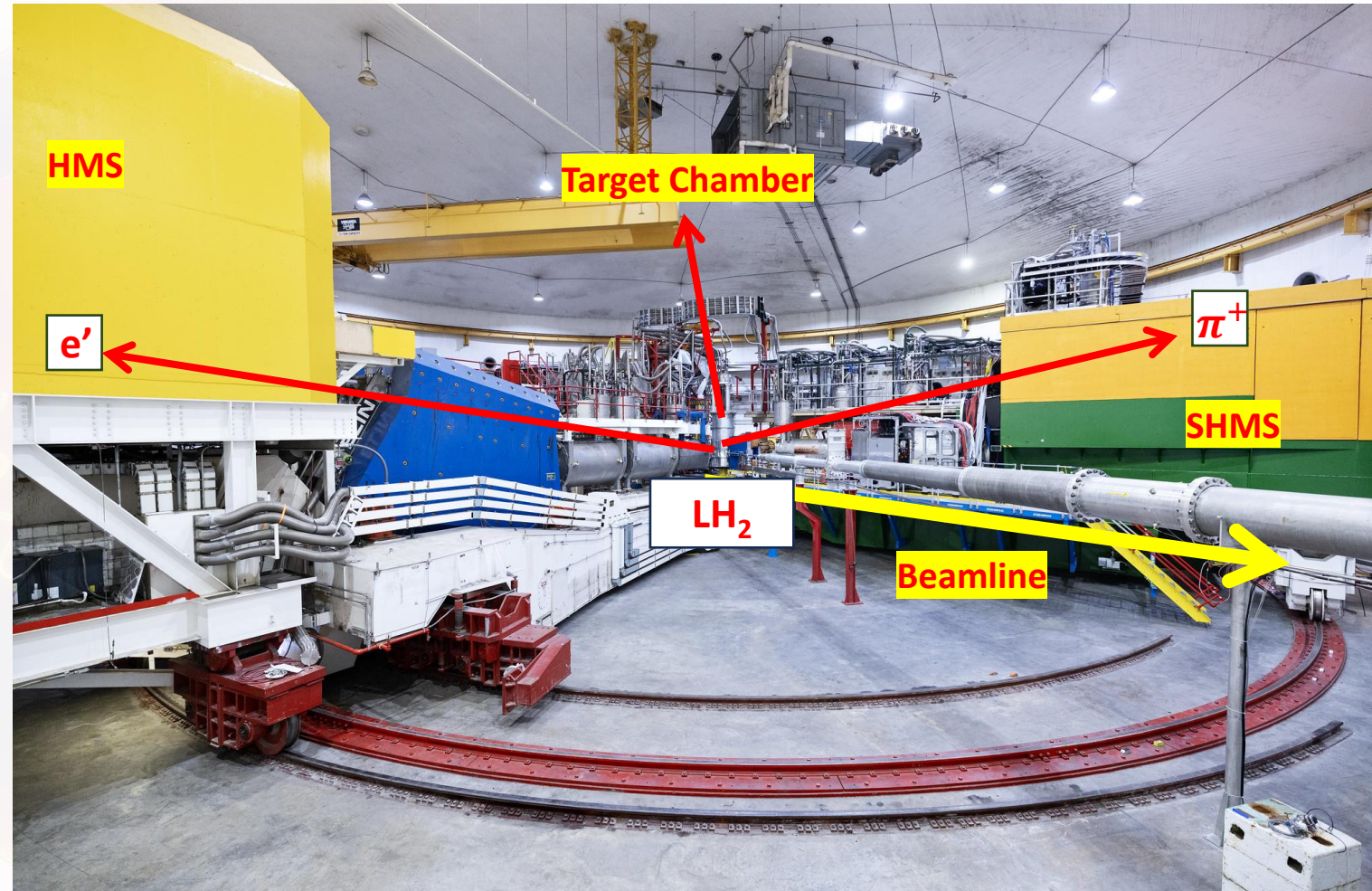


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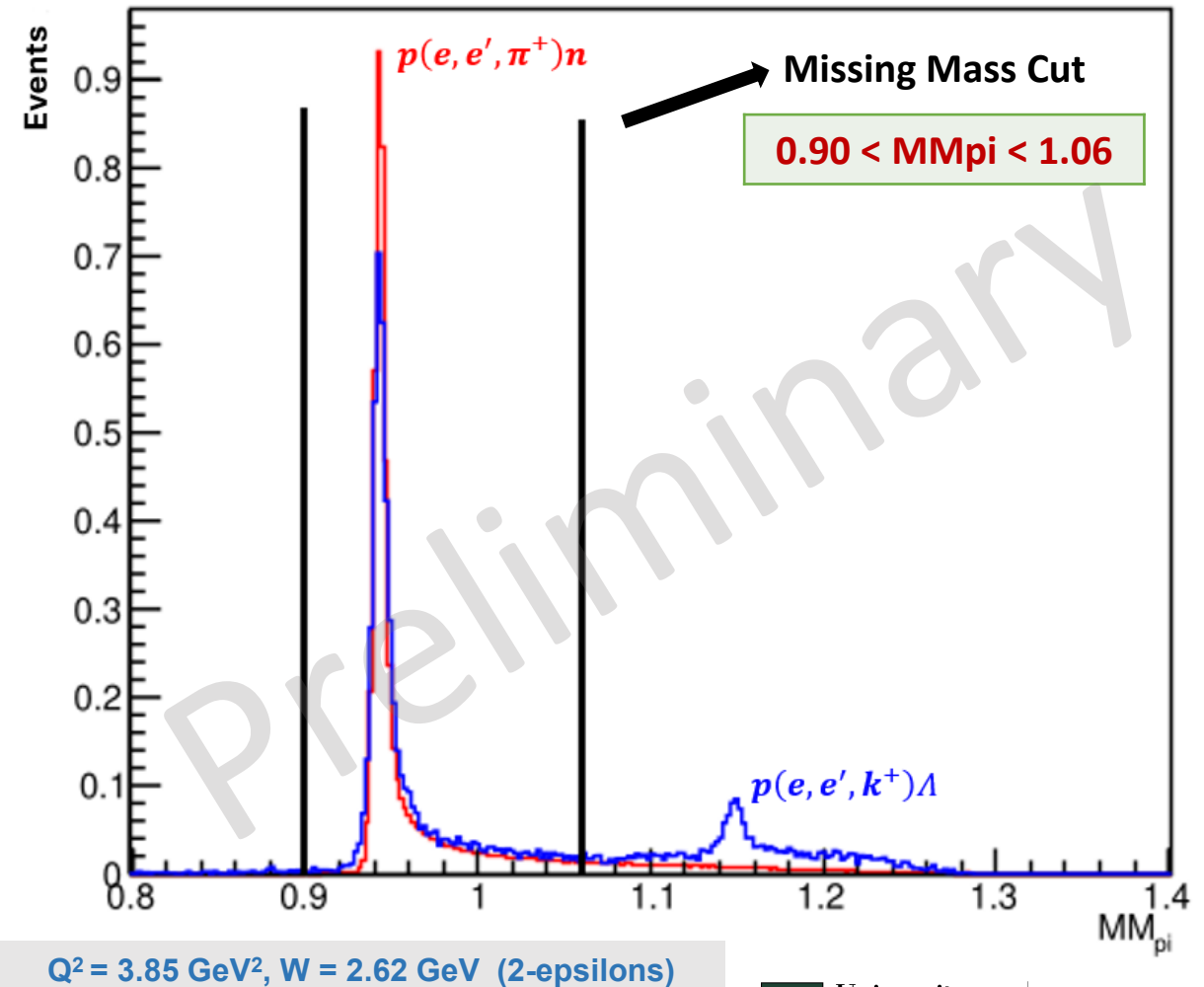
Selection of Event Sample

- Selected electrons in HMS in coincidence with pions in the SHMS.

- Missing particle reconstruction (neutron).

$$M_m = \sqrt{(E_e + m_p - E_{e'} - E_{\pi^+})^2 - (\mathbf{p}_e - \mathbf{p}_{e'} - \mathbf{p}_{\pi^+})^2}$$

- Blue represents the experimental data missing mass plot.
- Red represents the Monte Carlo missing mass plot.
- To eliminate the remaining background, a missing-mass cut was applied to both the Monte Carlo and experimental data.



Cross-section Measurements

- ❑ The goal is to calculate the experimental cross-sections using the ratio method.

$$\frac{d^2\sigma}{dtd\phi_{EXP}} = \left(\frac{Y_{EXP}}{Y_{SIMC}} \right) \frac{d^2\sigma}{dtd\phi_{SIMC}}$$

- ❑ The same cuts and kinematic selection criteria were applied to both the data and the Monte Carlo.
- ❑ This technique is model-dependent.
- ❑ Requires the Monte Carlo empirical model to reproduce data.
- ❑ Only reliable if Monte Carlo reproduces the data well in both shape and normalization.

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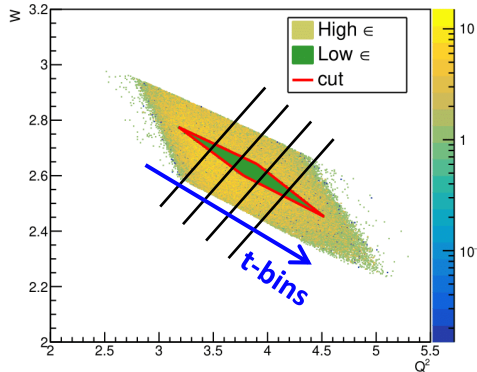
- ❑ The same cuts and kinematic selection criteria were applied to both the data and the Monte Carlo.
- ❑ This technique is model-dependent.
- ❑ Requires the Monte Carlo empirical model to reproduce data.
- ❑ Only reliable if Monte Carlo reproduces the data well in both shape and normalization.
- ❑ Fit the Rosenbluth equation to extract the cross-section components.

$$2\pi \frac{d\sigma}{dtd\phi} = \varepsilon \frac{d\sigma_L}{dt} + \frac{d\sigma_T}{dt} + \sqrt{2\varepsilon(\varepsilon + 1)} \frac{d\sigma_{LT}}{dt} \cos \phi + \varepsilon \frac{d\sigma_{TT}}{dt} \cos 2\phi$$

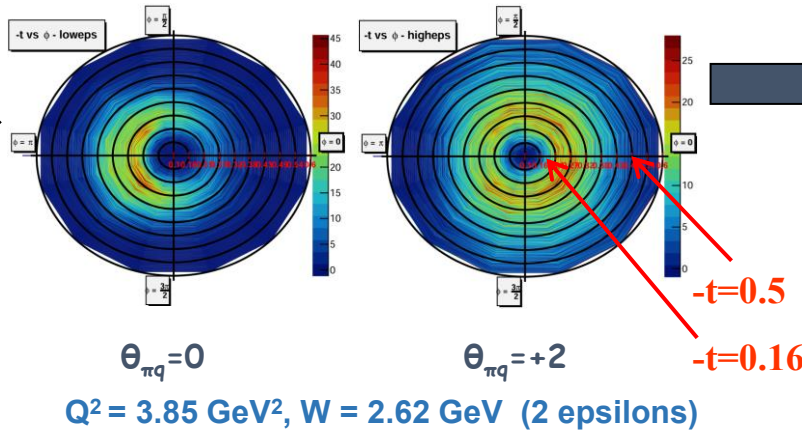
- ❑ Need to iteratively tune L/T/LT/TT empirical model until Monte Carlo reproduces experimental data.

L/T Separation Iteration Procedure

Diamond cut



Improve ϕ coverage by taking data at multiple π (SHMS) angles, $-2^\circ < \theta_{\pi q} < 2^\circ$.

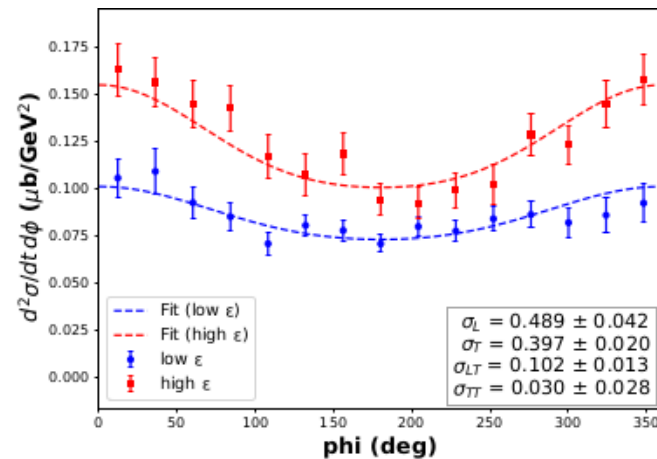
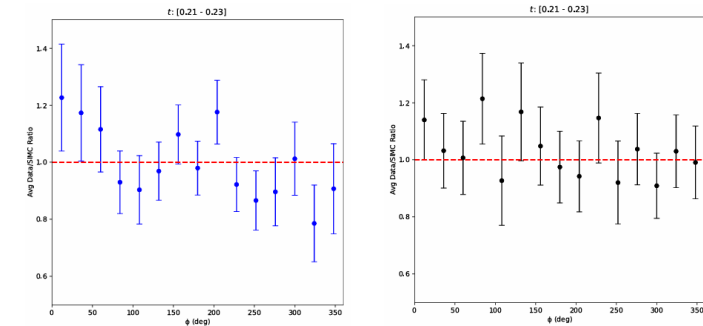


For each π HMS setting, form ratio:

$$R = \frac{Y_{EXP}}{Y_{SIMC}}$$

Combine ratios for π settings together, propagating errors accordingly.

Ratios



Extract via simultaneous fit of L,T,LT,TT

$$\frac{d^2\sigma}{dt d\phi}_{EXP} = \left(\frac{Y_{EXP}}{Y_{SIMC}} \right) \frac{d^2\sigma}{dt d\phi}_{SIMC}$$

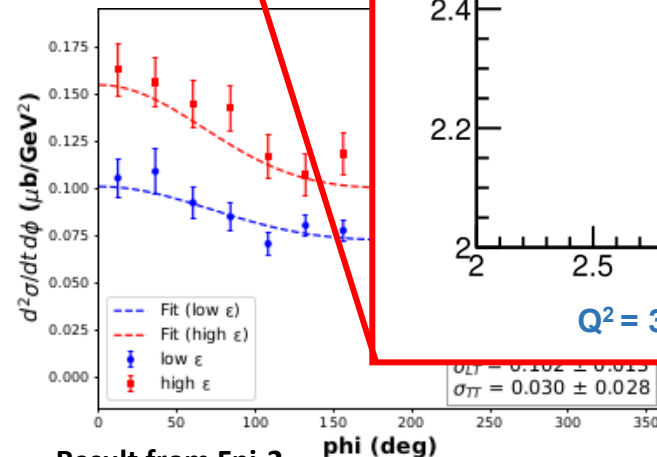
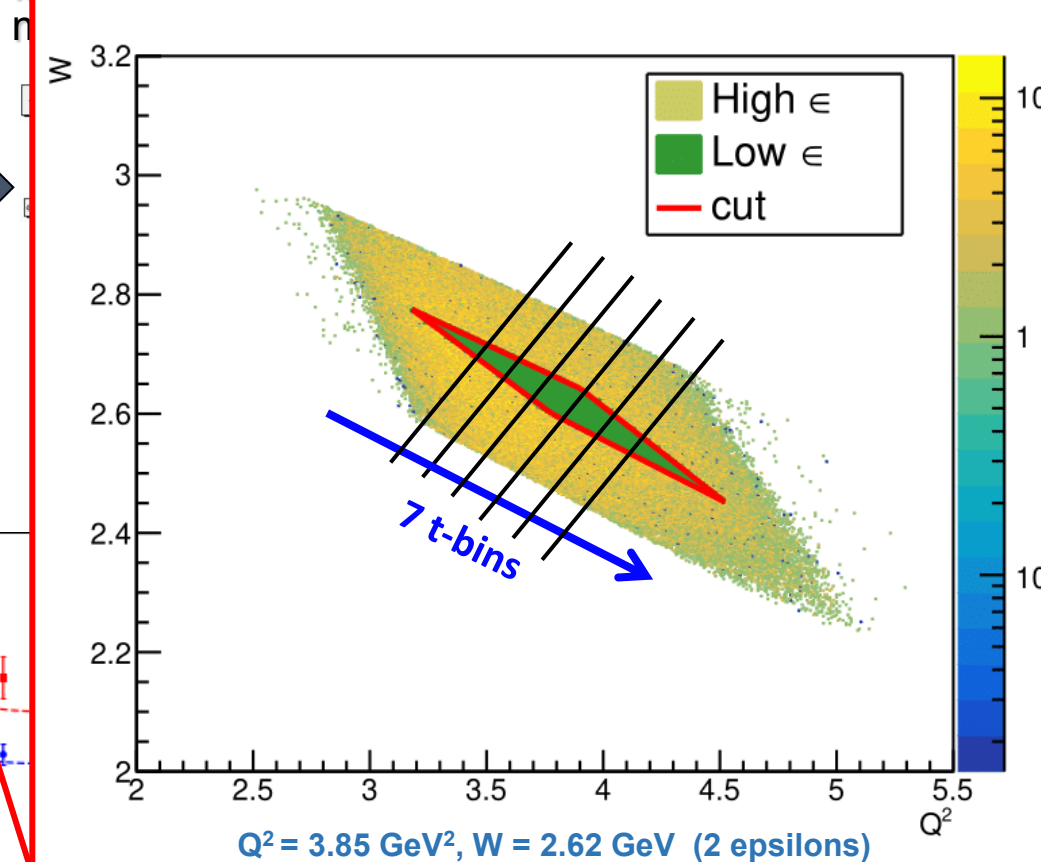
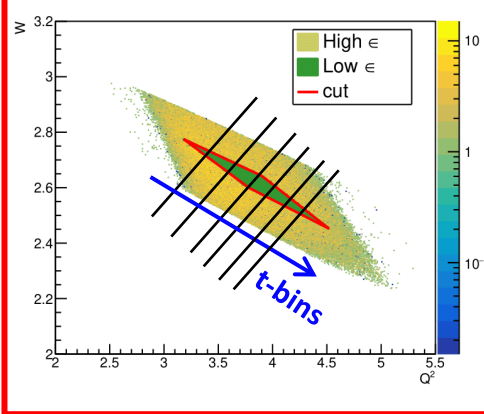
$$2\pi \frac{d\sigma}{dt d\phi} = \epsilon \frac{d\sigma_L}{dt} + \frac{d\sigma_T}{dt} + \sqrt{2\epsilon(\epsilon+1)} \frac{d\sigma_{LT}}{dt} \cos\phi + \epsilon \frac{d\sigma_{TT}}{dt} \cos 2\phi$$

Importance of Diamond Region

Improve ϕ coverage by taking data at

For each π HMS setting, form ratio:

Diamond cut



Result from Fpi-2

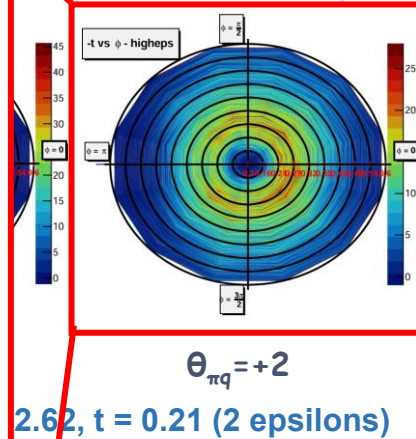
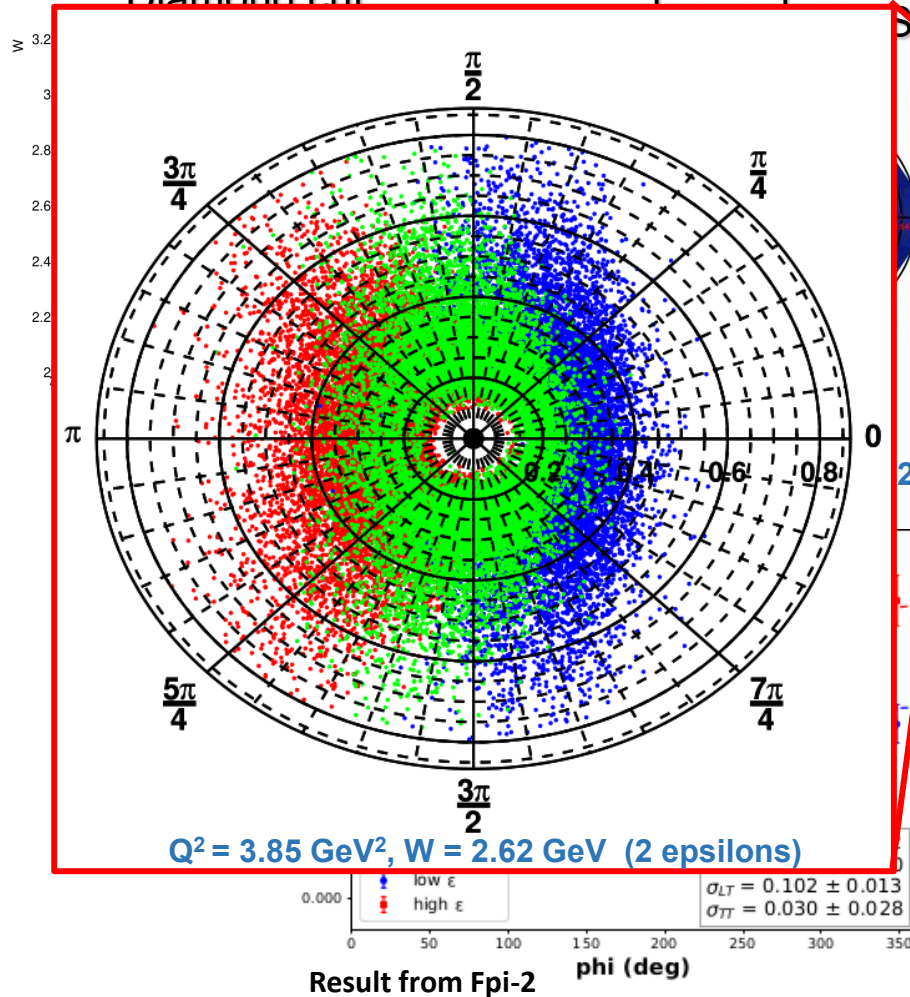
- Electron spectrometer acceptance is larger for high ϵ .
- Selected an overlapped phase-space region.
- Divided data into 7 t-bins based on data statistics.
- Purpose is to ensure consistency across different kinematic settings and measure the t-dependence

$$\frac{2\pi}{dtd\phi} = \epsilon \frac{d^2\sigma}{dt} + \frac{d^2\sigma}{dt} + \sqrt{2\epsilon(\epsilon+1)} \frac{d^2\sigma}{dt} \cos\phi + \epsilon \frac{d^2\sigma}{dt} \cos 2\phi$$

Full ϕ -Coverage

Diamond cut

Improve ϕ coverage by taking data at
small angles, $-2^\circ < \theta_{\pi q} < 2^\circ$.

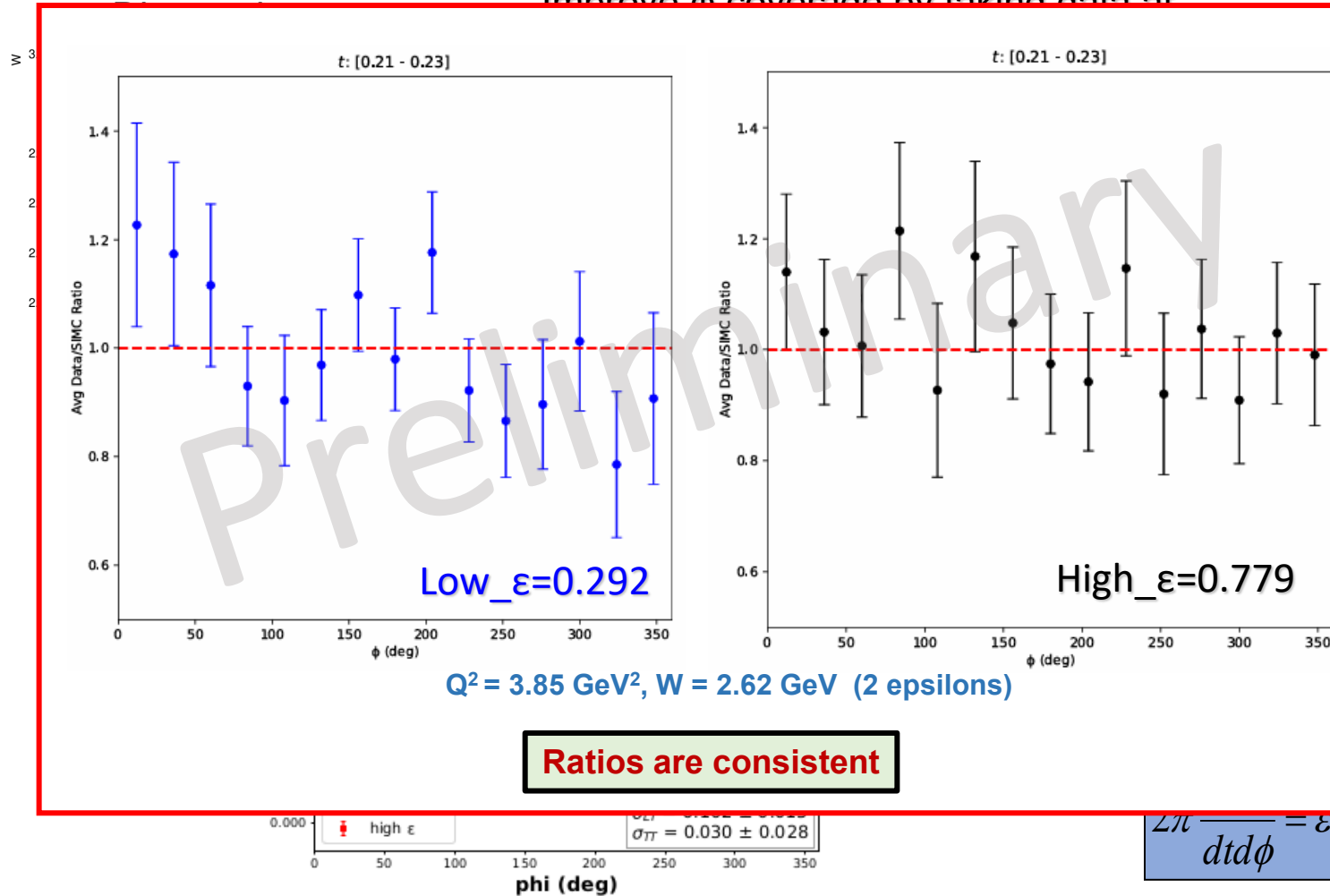


- To get full- ϕ coverage, data is taken on two degrees on the right and left of the central angle by rotating the pion arm.
- **Red corresponds to the right angle pion arm setting**
- **Green corresponds to the central angle pion arm setting**
- **Blue corresponds to the left angle pion arm setting**
- Divided data into 15 ϕ -bins to measure the ϕ dependence.

$$2\pi \frac{d\sigma}{dt d\phi}$$

Low and High Epsilon Data/MC Ratios

Improve ϵ coverage by taking data at

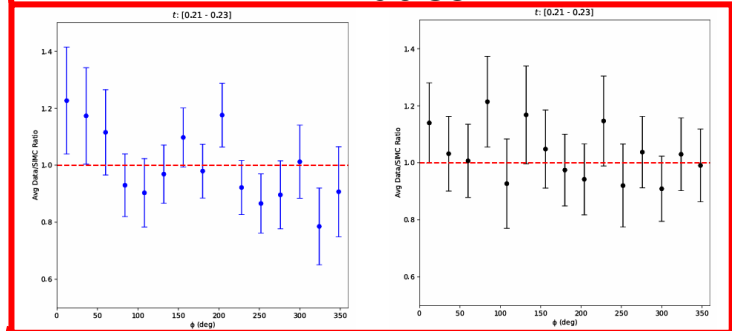


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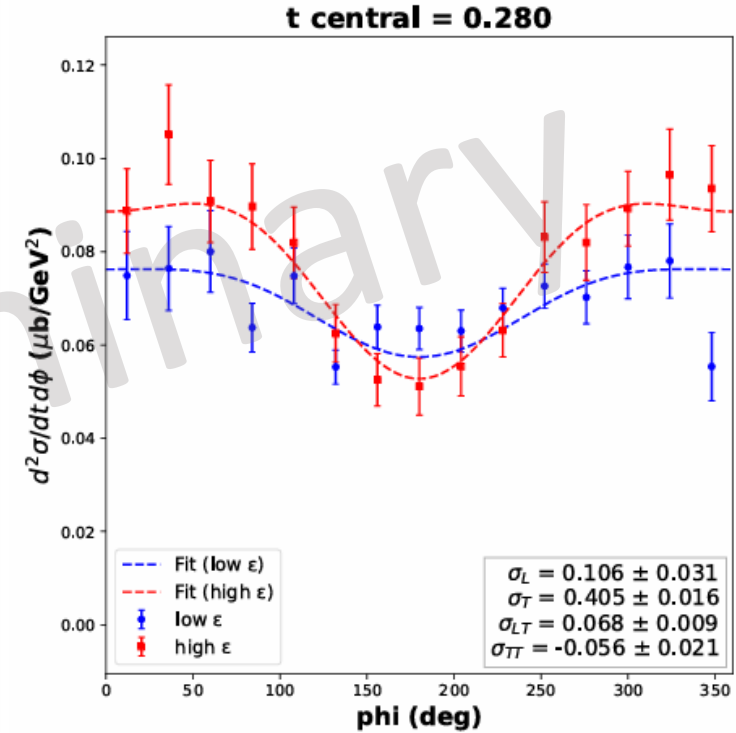
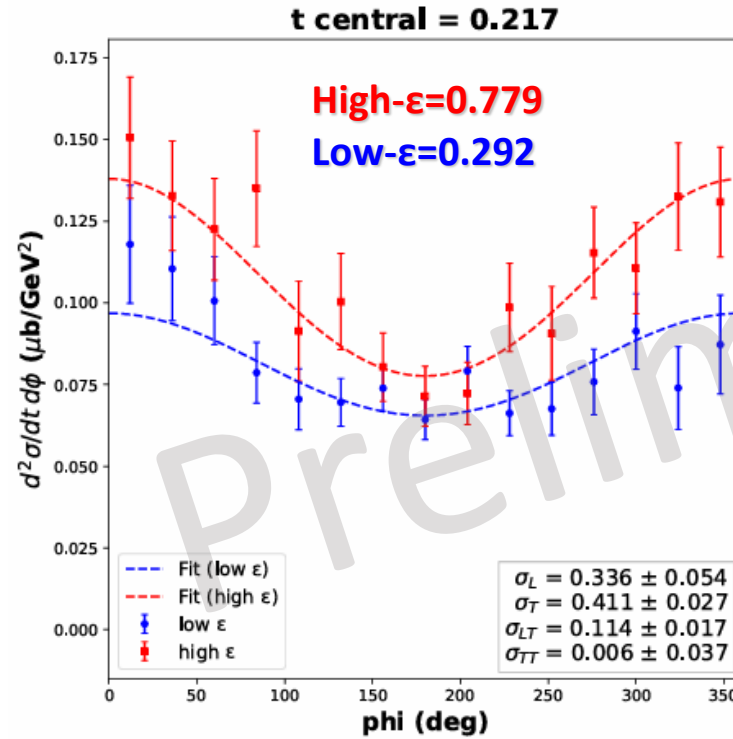
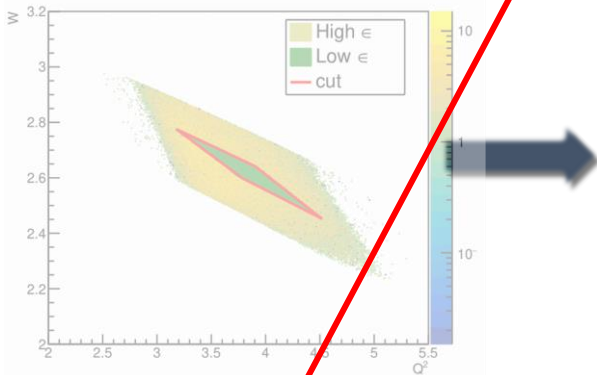
Ratios



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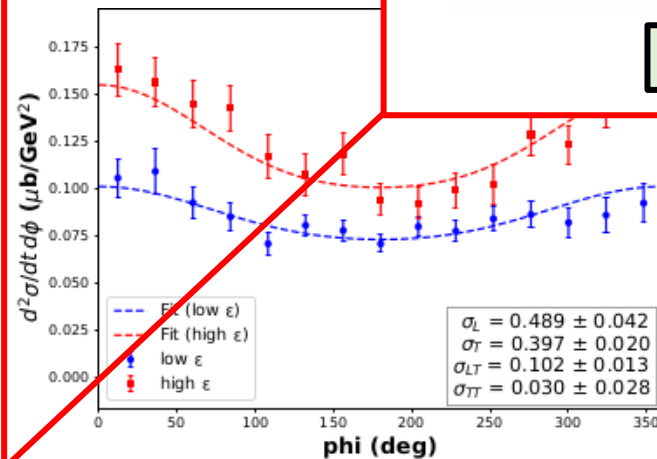
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L/T Separation

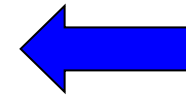


σ_L drops as we go to higher t

$Q^2 = 3.85 \text{ GeV}^2, W = 2.62 \text{ GeV}$ (2 epsilons)



Extract via
simultaneous fit
of L,T,LT,TT



$$\frac{d^2\sigma}{dtd\phi}_{EXP} = \left(\frac{Y_{EXP}}{Y_{SIMC}} \right) \frac{d^2\sigma}{dtd\phi}_{SIMC}$$

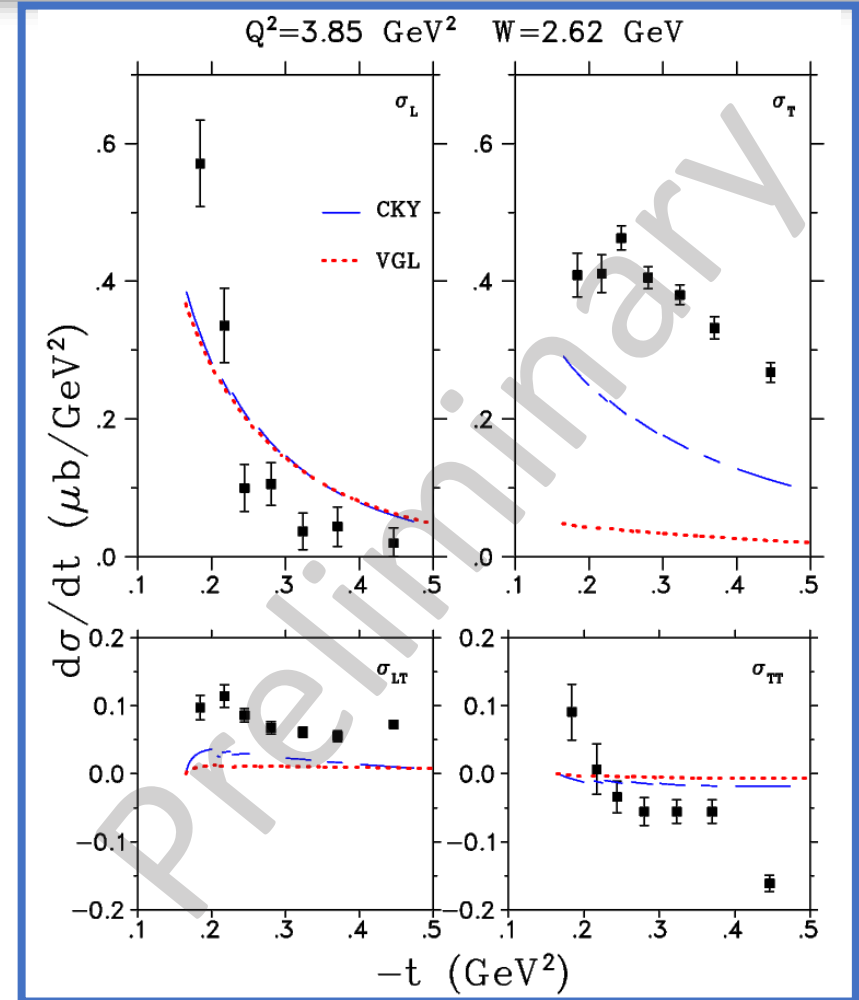
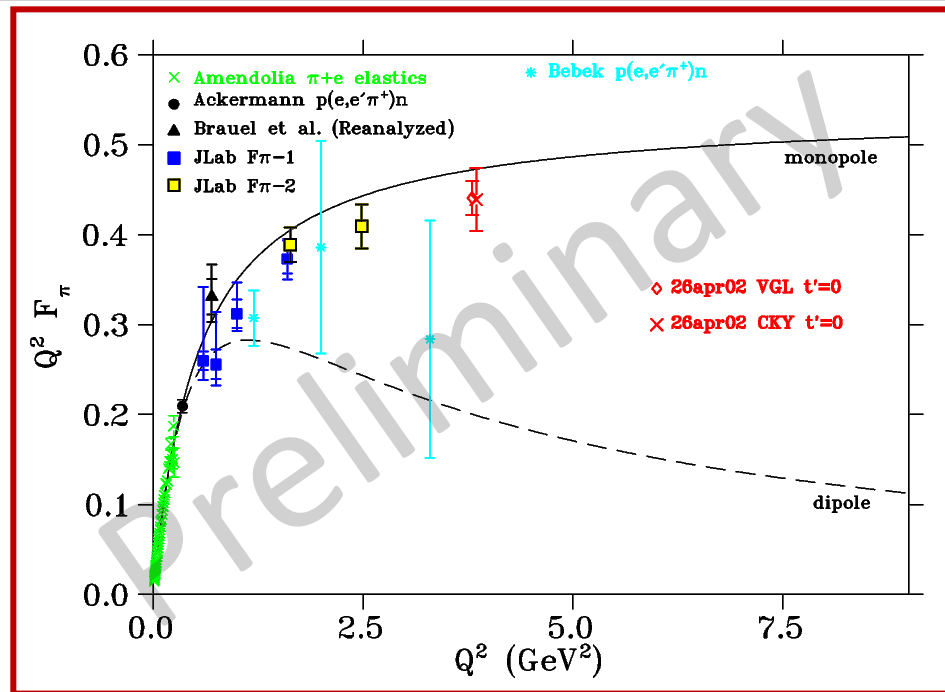
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ts



Results from Moderate- Q^2 Pion-LT Experiment

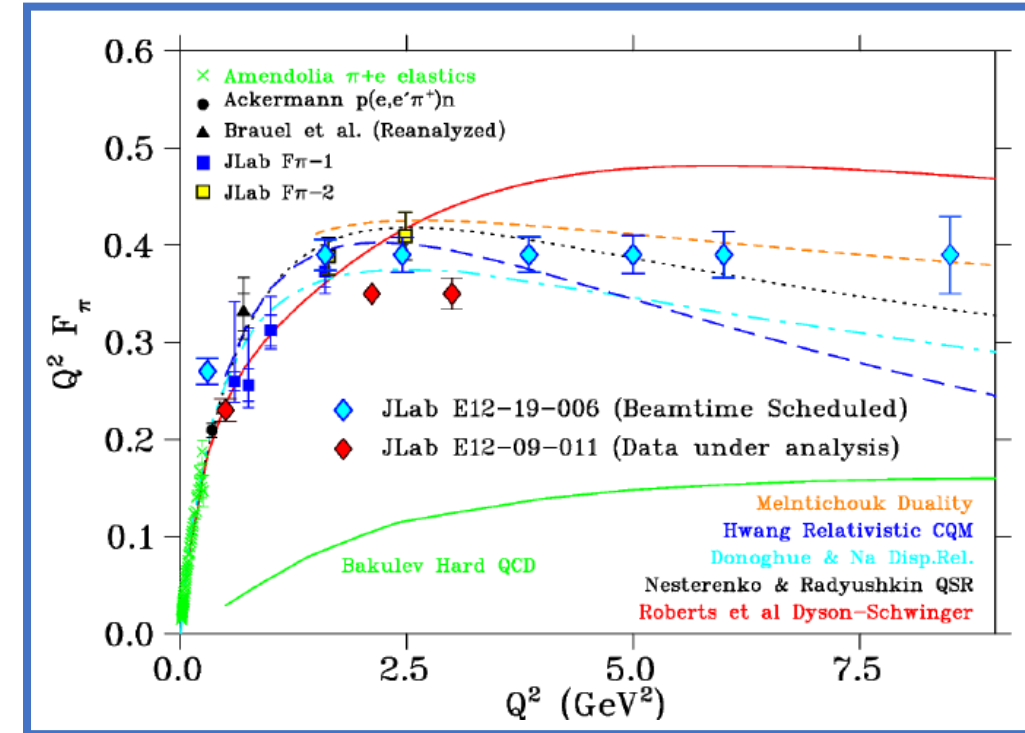
- The VGL model struggles to describe the t -dependence of σ_L well, and dramatically underestimates T/L ratio.
- The CKY model also struggles to describe the t -dependence of σ_L well, but has a better T/L ratio.
- Comparison with PKT Model is in progress.



Model is evaluated at precise kinematics of data. Discontinuities indicate change in (Q^2, W) for each t -bin.

Summary and Future Plans

- ❑ E12-19-006 (12 GeV Flagship Experiment) is expected to provide the definitive $p(e, e' \pi^+)n$ L/T-separation data set and will remain important for decades to come.
- ❑ Preliminary L/T separation is completed for $Q^2=3.85 \text{ GeV}^2$, $W=2.62 \text{ GeV}$ physics setting.
- ❑ Systematic uncertainty studies still need to be done.
- ❑ Detailed comparison with existing VGL, CKY, and PKT theoretical models.
- ❑ Results will help to understand the dependence of the Form factor and in validating theoretical models.



Projected Results

Acknowledgement

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University
of Regina



OHIO
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- **Graduate Students:**

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- **Pion-LT/Kaon-LT Collaboration:**

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THANKS



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BACKUP

Literature

Scientific Motivation

- Understanding the origin of the masses of protons and neutrons, as **> 99% of the visible mass of the universe resides in atomic nuclei.**

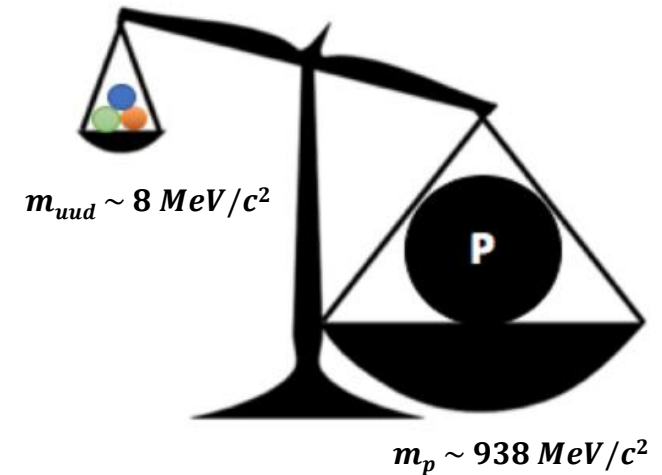
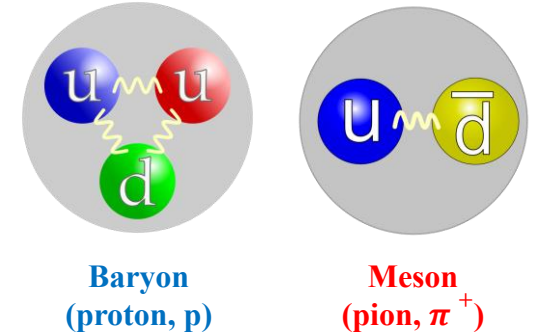
- Proton consists of two up-quarks (**u**) and one down-quark (**d**).

$$m_p \approx 938 \text{ MeV}/c^2$$

$$m_u \approx 2 \text{ MeV}/c^2 \text{ and } m_d \approx 4 \text{ MeV}/c^2$$

$$\text{Total mass} = (2 \times 2 + 4) \text{ MeV}/c^2 = 938 \text{ MeV}/c^2 ?$$

- The Higgs mechanism, through which fundamental particles acquire mass, can explain only a small fraction of the nucleon mass.
- The majority of the mass comes from the strong interactions that bind the quarks and gluons together.
- We do not adequately understand the mass generation mechanisms of the strong interaction.
- To address the emergence of hadronic mass, one can examine internal structure of the lightest mesons i.e pions, kaons.**
- Form factor** can be studied to understand the internal structure of hadrons.



What is Form Factor?

- Form Factor ($F(Q^2)$) describes the transverse spatial position of partons within hadrons.
- $F(Q^2)$ is the Fourier transform of the particle's spatial charge density ($\rho(r)$).

$$F(Q^2) = \int \rho(r) e^{iQ \cdot r} d^3r$$

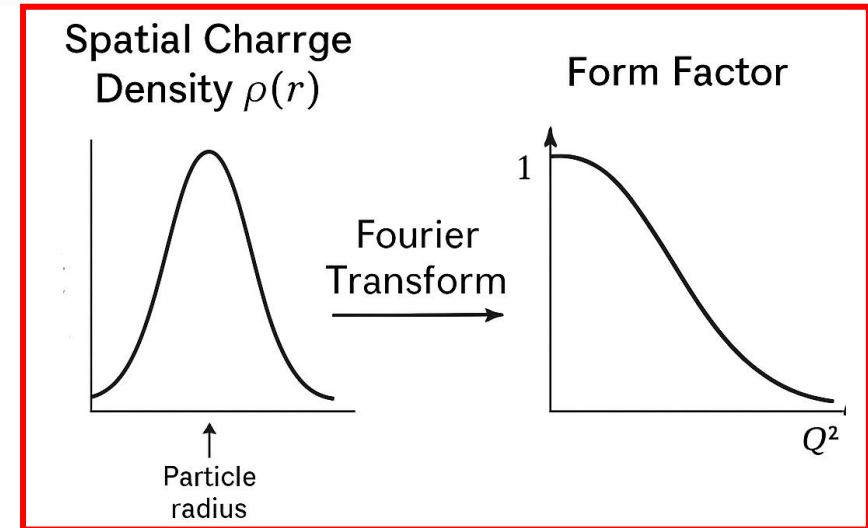
- At low Q^2 , the form factor directly relates to particle's charge radius $\langle r^2 \rangle$.

$$\langle r^2 \rangle = -6 \left. \frac{dF(Q^2)}{dQ^2} \right|_{Q^2=0}$$

- At high Q^2 , form factor follows predictions from perturbative QCD and helps to understand nature of the quark-gluon coupling constant renormalization.

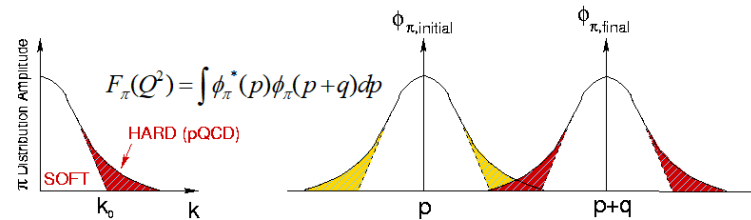
$$Q^2 F_\pi(Q^2) \rightarrow 16\pi\alpha_s(Q^2) f_\pi^2 \quad (Q^2 \rightarrow \infty)$$

- Measuring the pion form factor at various Q^2 checks the validity of QCD-based theories, including the transition between region between perturbative and non-perturbative approaches.

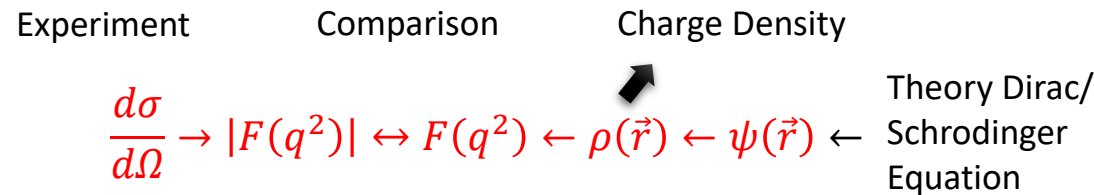


Pion Form Factor

- ❑ Mesons give an ideal testing ground for our understanding of bound quark-antiquark system.
- ❑ **Form Factor** describes transverse spatial position of partons within hadrons.
- ❑ The Form Factor of a particle is defined in Quantum Field Theory as an overlap integral;



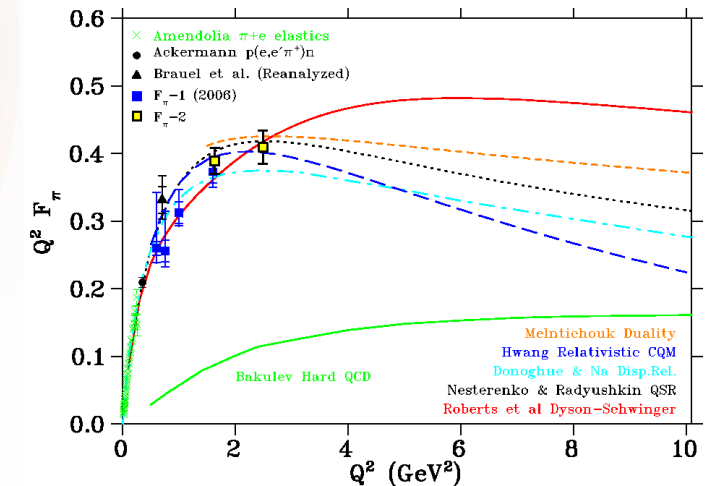
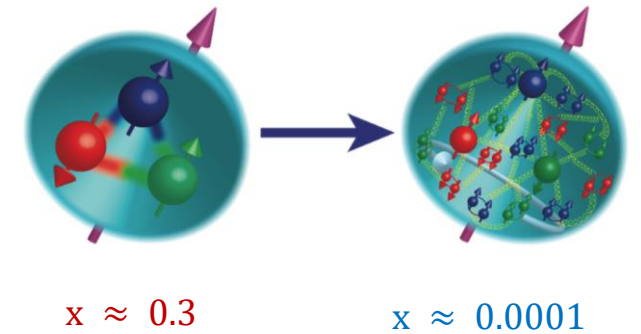
- ❑ where φ is the wave function is an important variable to study as it allows for comparison with theory:



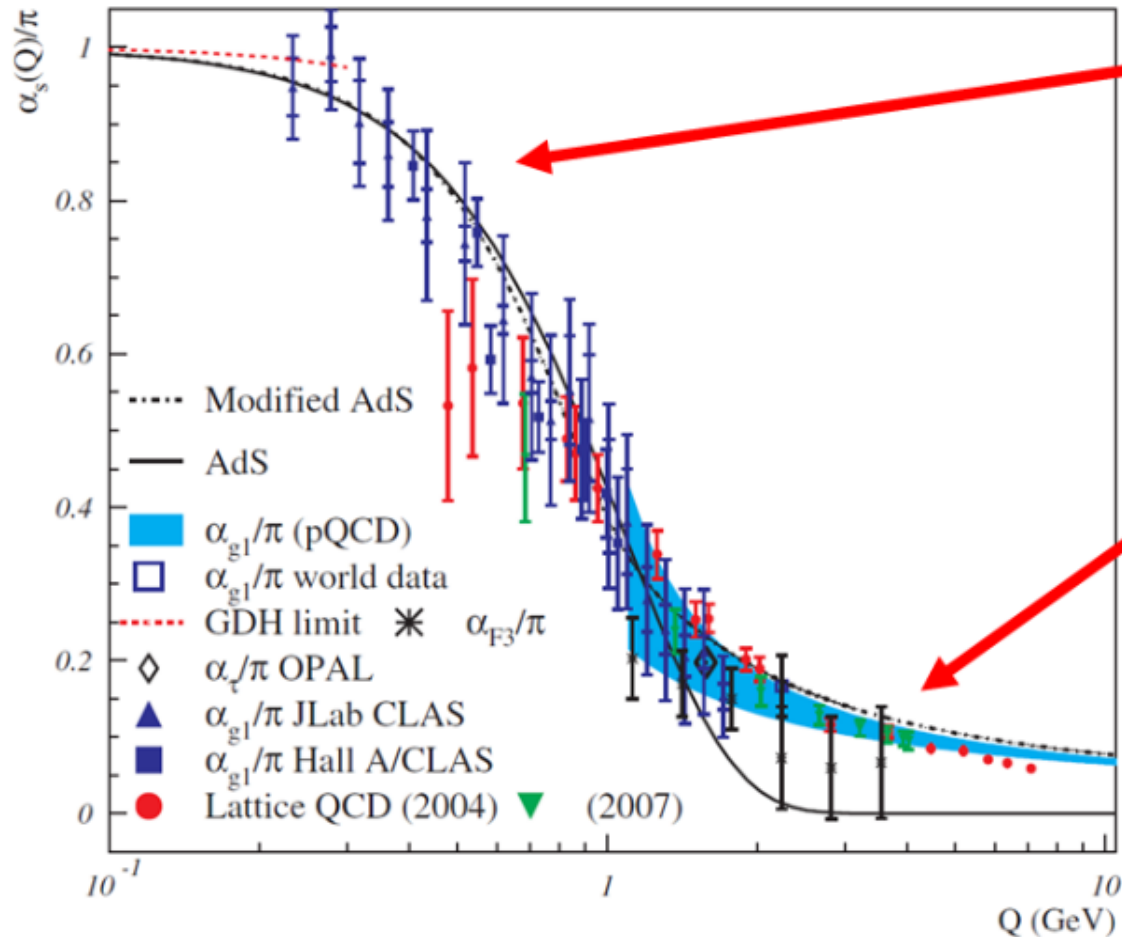
- ❑ The meson wave function can be separated into $\varphi_{\pi}^{\text{soft}}$ with only low momentum contributions ($k < k_0$) and a hard tail $\varphi_{\pi}^{\text{hard}}$. While $\varphi_{\pi}^{\text{hard}}$ can be treated in pQCD, $\varphi_{\pi}^{\text{soft}}$ cannot.
- ❑ From a theoretical standpoint, the study of the Q^2 -dependence of the form factor focuses on finding a description for the hard and soft contributions of the meson wave-function.

Scientific Motivation

- ❑ QCD is an extremely successful theory, but it struggles to predict the internal structure of hadrons.
- ❑ In QCD, there are two main domains of the strong interaction.
 - ❑ Perturbative Quantum Chromodynamics (pQCD) corresponds to short-range quark-gluon interaction part (asymptotic freedom)
 - ❑ Non-Perturbative Quantum Chromodynamics (non-pQCD) corresponds to long-range quark-gluon interaction part (confinement).
- ❑ pQCD predicts the F_π at very high Q^2 (i.e. $Q^2 \rightarrow \infty$) and it is well understood.
- ❑ In non-pQCD region, various theoretical models struggle to predict F_π at low Q^2 (i.e. $Q^2 \approx 1 - 10 \text{ GeV}^2$).
- ❑ These QCD-based models disagree on how quickly the transition between pQCD and non-pQCD should occur.
- ❑ **Pion Form factor** is an important observable that can be studied to understand the transition region between pQCD and non-pQCD.



Dual Nature of QCD



Long Distance (Low Energy) Interaction:

- Quarks strongly bound within hadrons
- **Color Confinement**
- Quantitative QCD description of hadron properties remains a puzzle
- Unfortunately, this is the regime under which we live!

Short Distance (High Energy) Interaction:

- Short range interaction is weak
- **Quarks inside hadrons behave as though they are nearly unbound**
- Can understand behaviour using perturbative QCD (pQCD)
- Asymptotic Freedom

F_π sensitivity to Emergent Mass

At empirically accessible energy scales, F_π is sensitive to emergent mass scale in QCD.

Two dressed-quark mass functions distinguished by amount of DCSB

- DCSB emergent mass generation is 10% stronger in system characterized by solid green curve, which is more realistic case.

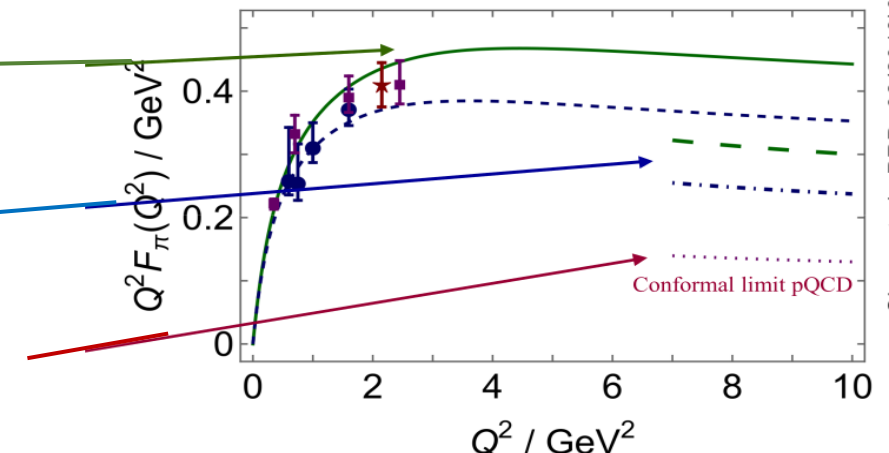
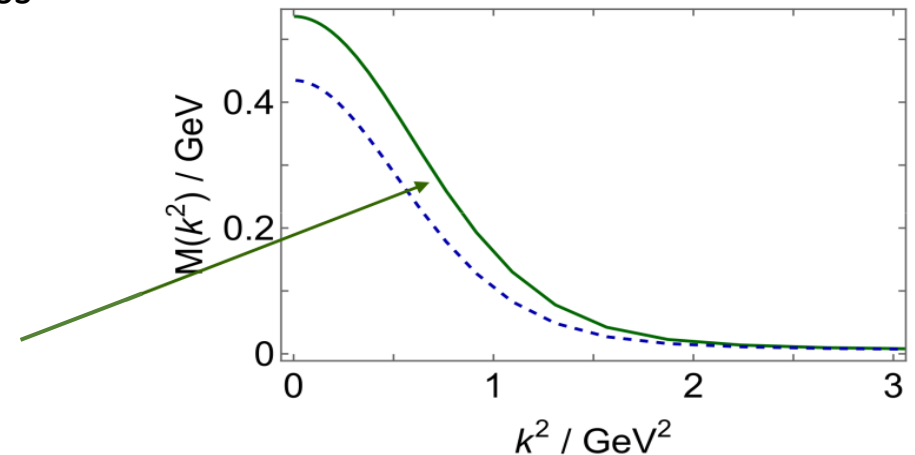
$F_\pi(Q^2)$ obtained with these mass functions.

- $r_\pi = 0.66$ fm with solid green curve
- $r_\pi = 0.73$ fm with solid dashed blue curve

$F_\pi(Q^2)$ predicts from QCD hard scattering formula, obtained with related, computed pion PDAs

QCD hard scattering formula, using conformal limit of pion's twist-2 PDA

$$\phi_\pi(x) = 6x(1-x)$$

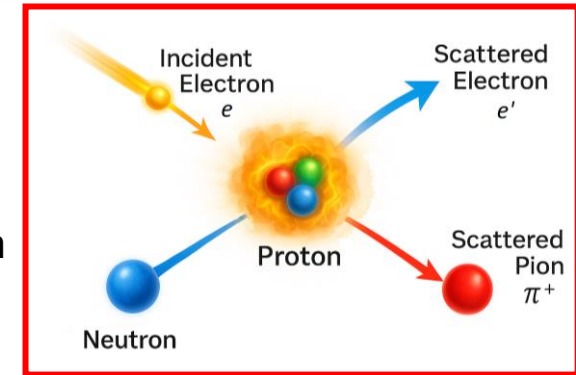


Chen, et al., PRD 98(2018)091505(R); Aguilar et al., EPJA 55(2019)190

Pion Form Factor Measurement

□ Direct Measurement:

- Elastic Scattering of electrons from pions gives
$$e + \pi^+ \rightarrow e' + \pi^{+'}$$
- Limitation:** Pion ($\pi^\pm$) targets not possible due to short lifetime ($\sim 2.6 \times 10^{-8}$ s)
- Even scattering high energy pion beam (1 TeV, if some facility could be constructed) can access only $Q^2 \sim 1 \text{ GeV}^2$



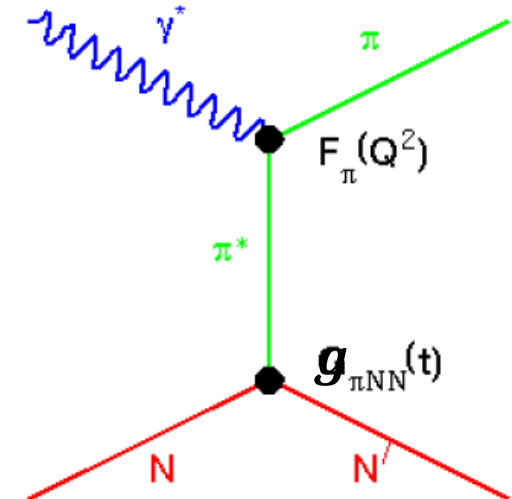
□ Indirect Measurement:

- Above $Q^2 > 0.3 \text{ GeV}^2$, F_π is measured indirectly using the “pion cloud” of the proton via pion electroproduction $p(e, e' \pi^+) n$
- Indirect measurement – Form factor extraction requires a model.

- As an illustration of how σ_L connects to $F_\pi^2(Q^2, t)$, we consider a simple **Born Term Model**;

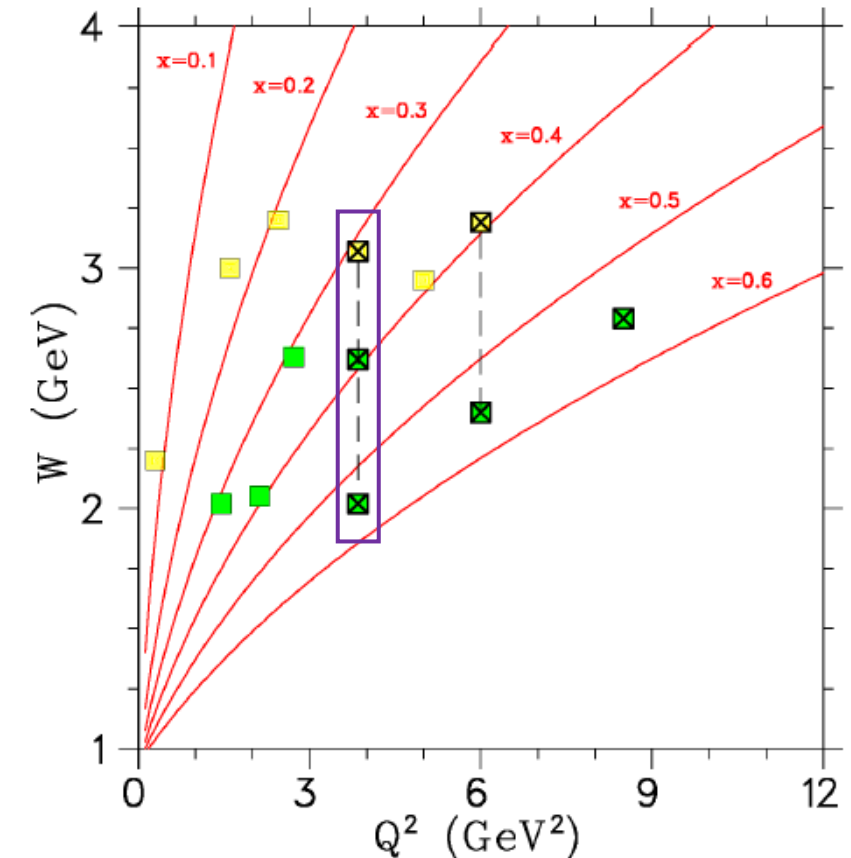
$$\frac{d\sigma_L}{dt} \propto \frac{-tQ^2}{(t - m_\pi^2)} g_{\pi NN}^2(t) F_\pi^2(Q^2, t)$$

- In reality, we use Regge base model such as VGL, YCK and PKT Models for $F_\pi^2(Q^2, t)$ extraction.



Pion-LT Experiment

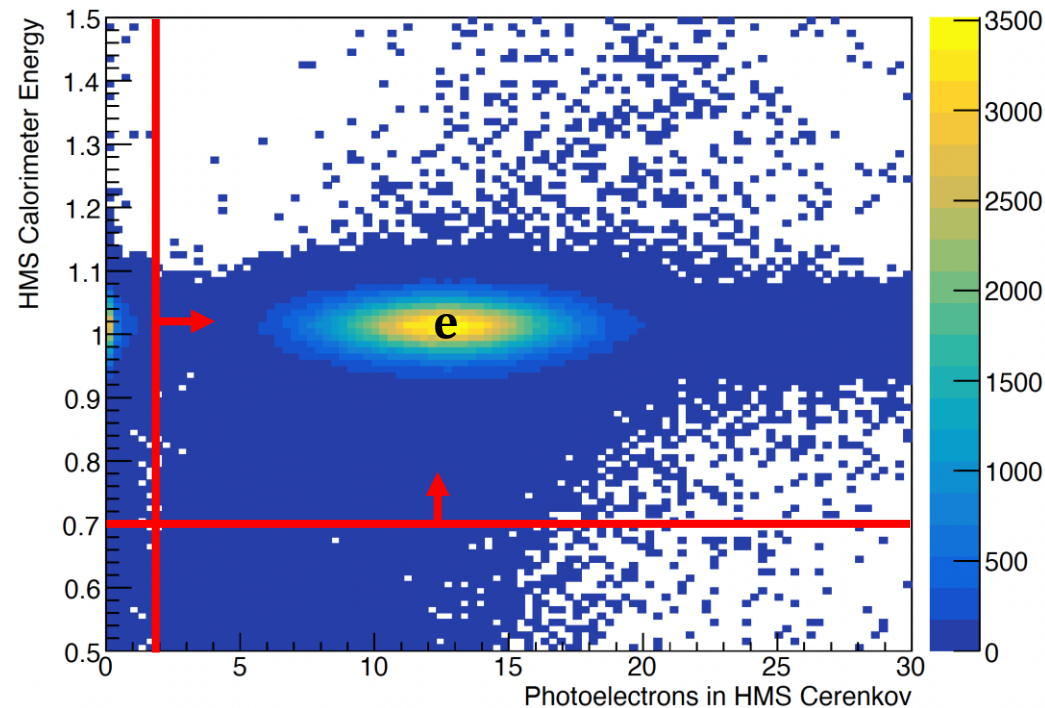
- ❑ The Pion-LT experiment was conducted over a wide kinematic range at the Hall C facility, JLab.
- ❑ Green points represent the Pion Form Factor Study.
- ❑ Yellow points represent the Pion Scaling Study.
- ❑ Vertical black dashed lines shows W -scan at fixed Q^2 .
- ❑ Points marked with an 'x' are instrumental in higher Q^2 , F_π extraction
 - ❑ $Q^2 = 8.5 \text{ GeV}^2$ is highest achievable extraction at JLab
- ❑ Red lines allow for $1/Q^n$ scaling study at fixed $x = 0.31, 0.39, 0.55$.
- ❑ Focusing on the physics settings highlighted in purple box.



PID Study

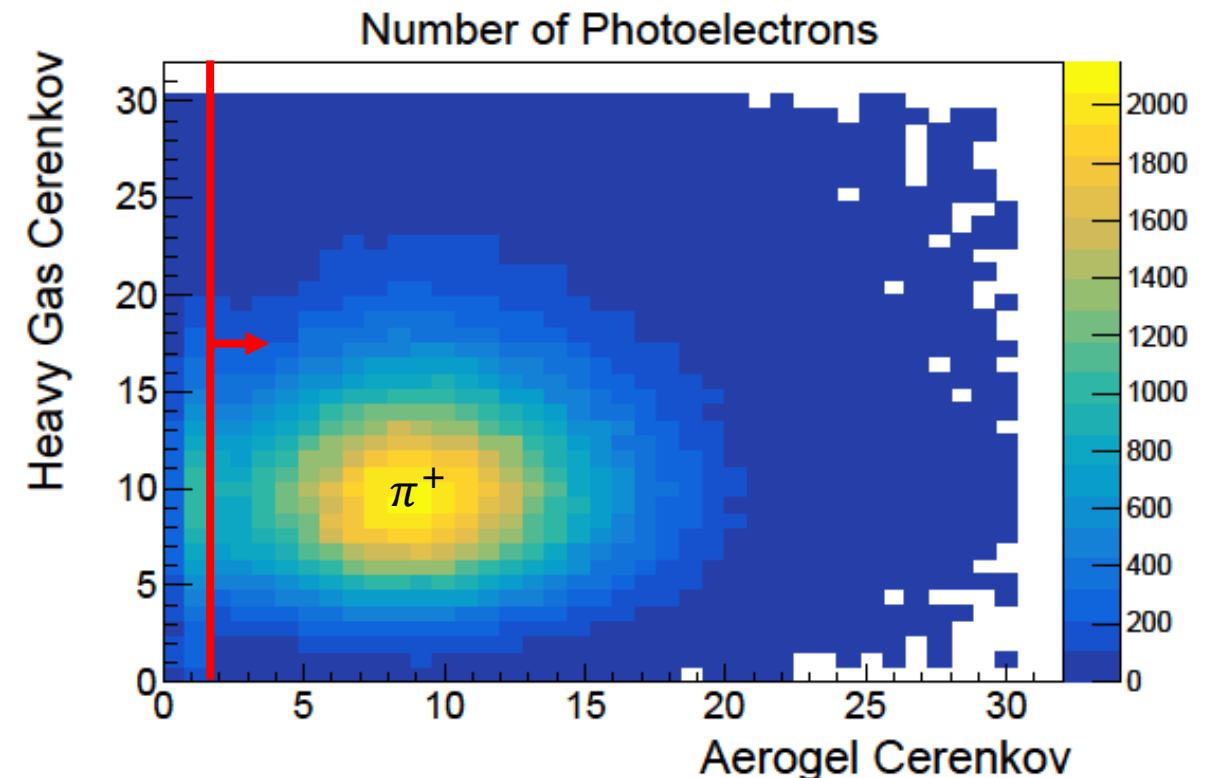
Event Selection Criteria

- Purpose is to select the clean pion events sample in coincidence with the electrons to measure the cross-sections.
- Cuts must be tight, but not so close as to sacrifice efficiency across the acceptance.
- Overall, detector efficiencies are greater than ~98%.



HMS PID

$Q2 = 3.85, W = 2.62, t = 0.21$ (2 epsilons)



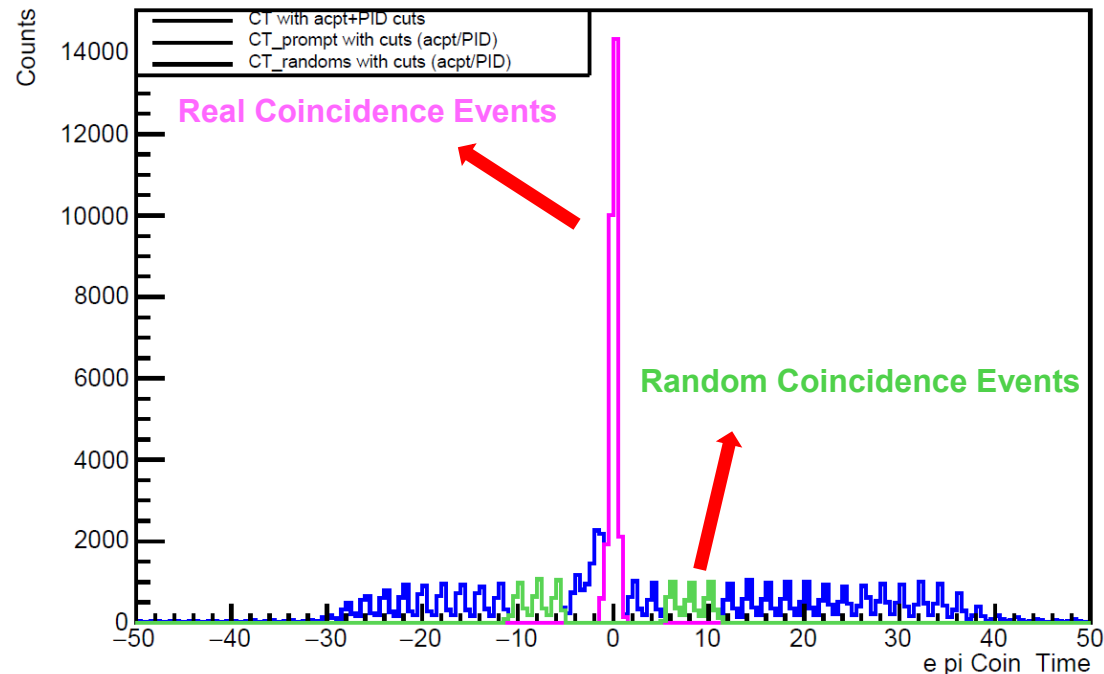
SHMS PID

Event Selection Criteria

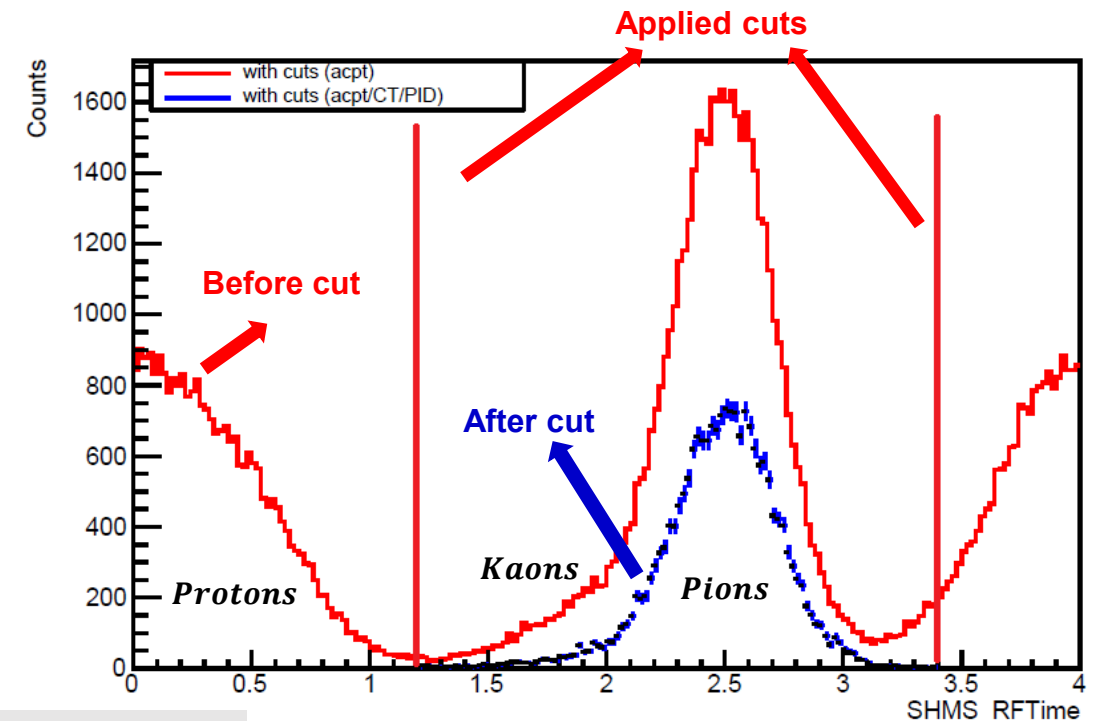
- CoinTime (CT) is given as:

$$t_{\text{coin}} = t_{\text{HMS}} - t_{\text{SHMS}}$$

- Used to eliminate random background caused by random coincidences.



- RF time refers to the time difference between a detected event and the nearest preceding RF signal of the electron linac.
- Enables the differentiation between pions, kaons, and protons due to their distinct velocities.



Q2 = 3.85, W = 2.62, t = 0.21 (2 epsilons)

Efficiency Study

Tracking Efficiency Study

- The purpose of this study is to determine the tracking efficiency to account for the loss in the experimental normalized yields due to the bad track reconstruction by the Hall C analyzer (HCANA).

- The tracking efficiency can be expressed as;

$$\epsilon_{track} = \frac{N_{did,track}}{N_{should,track}}$$

- $N_{should,track}$ is the events where at least one track was expected after cuts (Acceptance + HMS PID + SHMS PID)
- $N_{did,track}$ is the events where at least one track formed by tracking process after cuts ($N_{should,track} + DC \text{ tracks} > 0$).

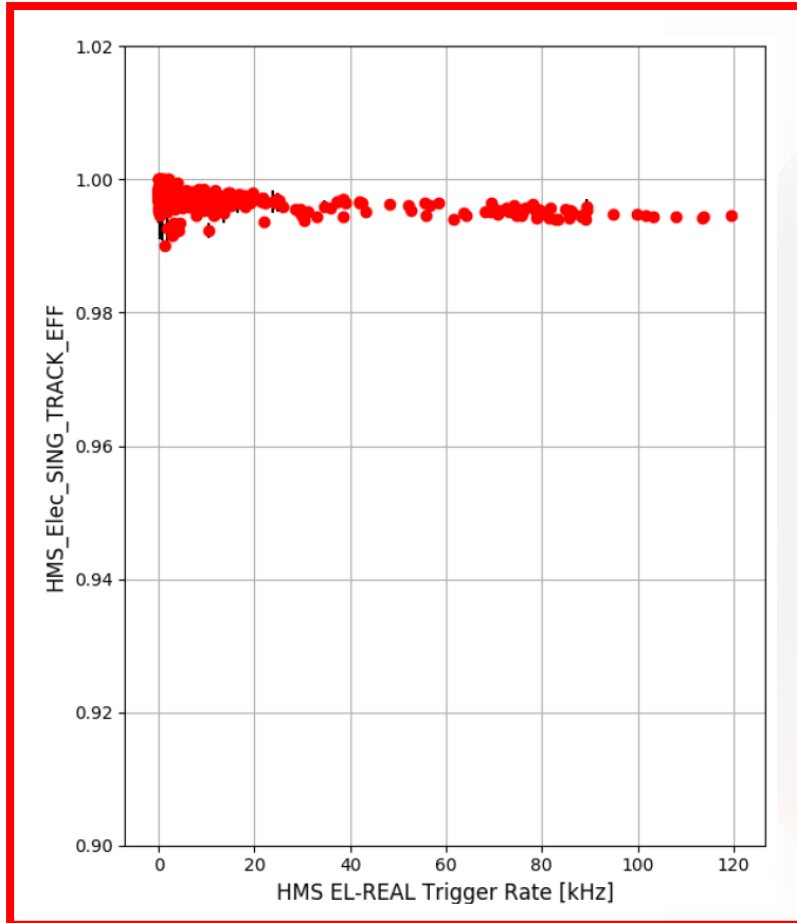
- The tracking efficiency error calculation can be expressed as;

$$\delta\epsilon_{track} = \sqrt{\frac{(N_{should,track} \times N_{did,track}) - (N_{did,track} \times N_{did,track})}{(N_{should,track})^3}}$$

HMS Acceptance Cuts*	SHMS Acceptance Cuts
$-8.0\% \geq \delta \geq 8.0\%$	$-10.0\% \geq \delta \geq 20.0\%$
$-0.08 \geq x'_{tar} \geq +0.08$	$-0.08 \geq x'_{tar} \geq +0.08$
$-0.045 \geq y'_{tar} \geq +0.045$	$-0.04 \geq y'_{tar} \geq +0.04$
HMS PID Cuts (Electrons)	
HMS Gas Cerenkov (NPE) > 1.5	
HMS Calorimeter (Energy) > 0.7	
SHMS PID Cuts (Pions)	SHMS PID Cuts (Protons)
HGC Cerenkov (NPE) > 1.0	HGC Cerenkov (NPE) > 1.0
Aerogel (NPE) > 1.5	Aerogel (NPE) > 1.5
Calorimeter (Energy) > 0.05	Calorimeter (Energy) > 0.05

- δ is the deviation of particle momentum from the spectrometer central momentum
- x'_{tar} and y'_{tar} are the angle of the particle trajectory at the target.
- NPE (Number of Photoelectrons) corresponds to the number of photoelectrons detected by a photomultiplier tube (PMT) when a charged particle passes through a detector.

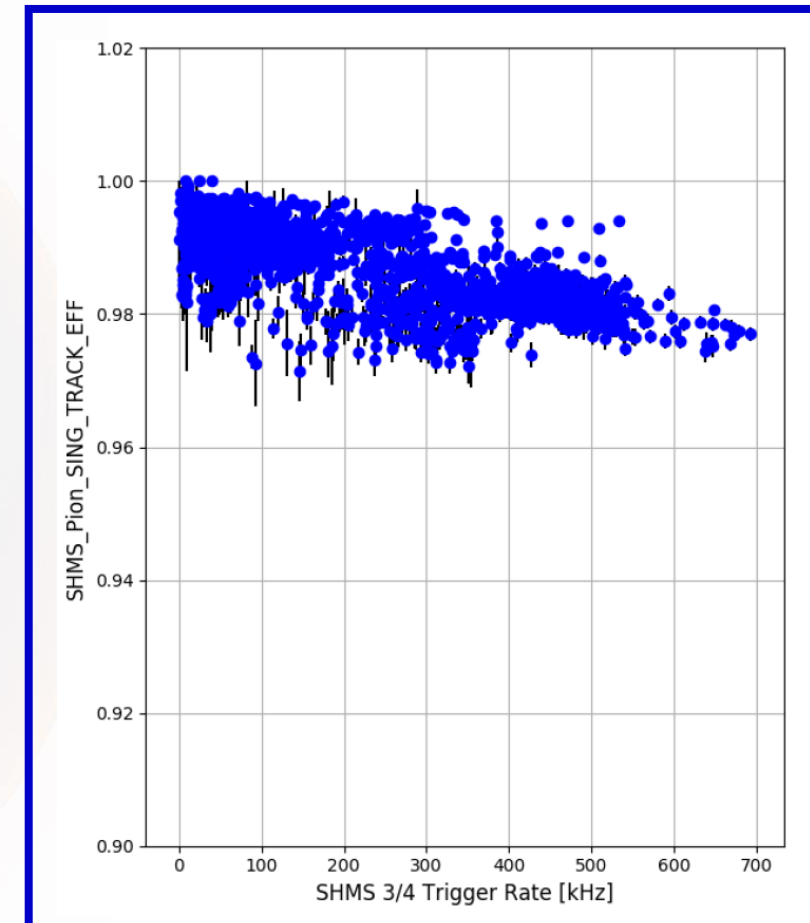
Tracking Efficiency Study



HMS Electron Tracking Efficiency

Tracking Efficiency plots from physics data.

- Both electron and pion tracking efficiencies are plotted against the detector trigger rates to check the rate dependence.
- Overall, the pion and electron tracking are looking good.
- The electron tracking is consistently near 100% ($\geq 99\%$).
- The pion tracking has a bit of rate dependence and maintained above $\sim 97\%$.

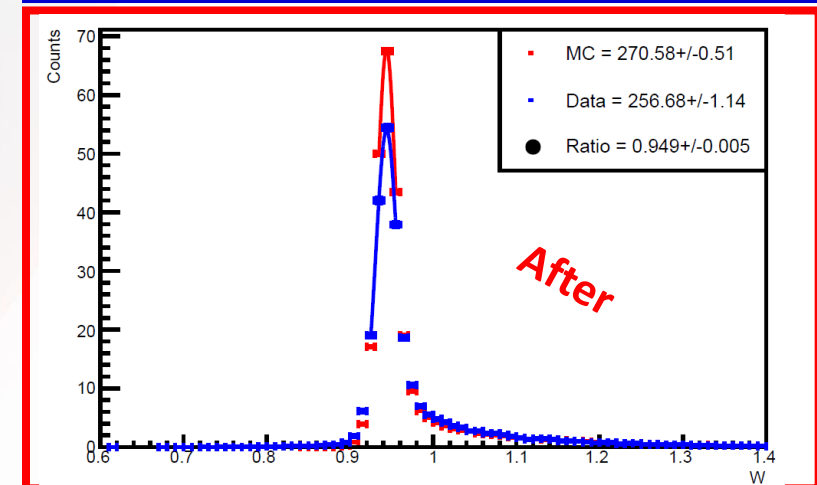
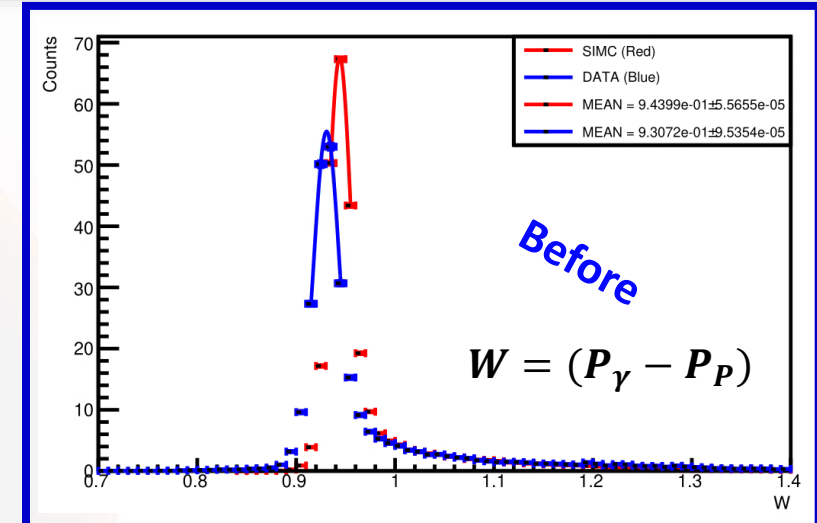


SHMS Pion Tracking Efficiency

HeeP Study

HeeP Offset Study

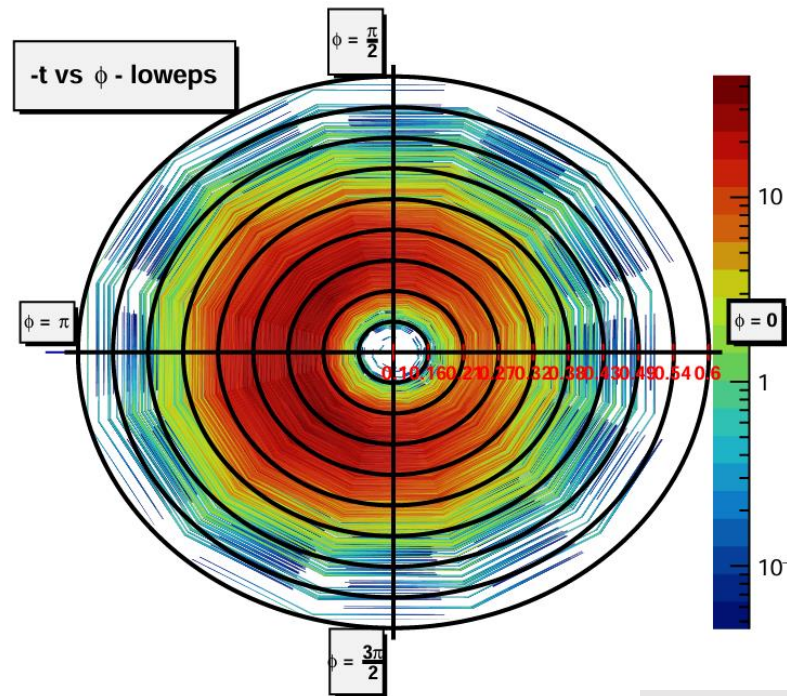
- ❑ looked at distributions of the missing variables ($E_M, \vec{P}_M, W$) to find deviations from expected values based on kinematics.
- ❑ The E_M , the components of $\vec{P}_M$, and W are correlated to different spectrometer quantities (momentum, angle, beam energy), and thus offsets can be extrapolated from them.
- ❑ By adjusting the beam energy, momentum and angles, missing variables ($E_M, \vec{P}_M, W$) can be moved to better agree between data and SIMC.
- ❑ Red is SIMC and Blue is data.
- ❑ Offsets are doing a good job and peaks are matching with each other.
- ❑ Applied the determined HeeP offsets to the physics data for cross-section measurement.



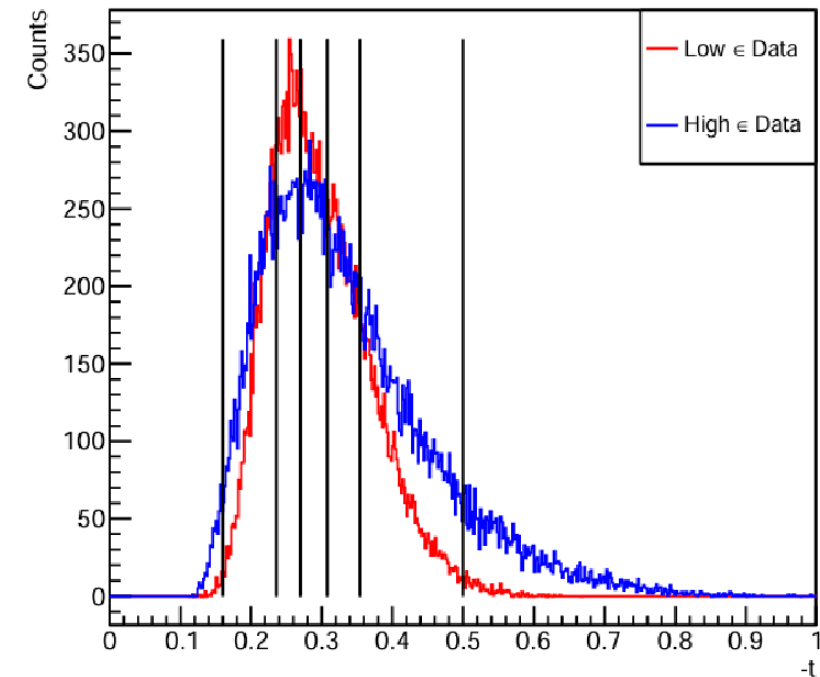
L/T-Separation Analysis

Determination of t-binning

- Purpose of t-bining is to maintain both statistical consistency and physical reliability in the extraction of the separated cross-sections.
- t-range is determined by looking at tvsphi polar plot of $Q^2=3.85 \text{ GeV}^2$, $W=2.62 \text{ GeV}$ (low-epsilon) physics setting.
- Divided “t” into 5 bins for $Q^2=3.85 \text{ GeV}^2$, $W=2.62 \text{ GeV}$ (2-epsilon) physics setting.
- Tried to keep the number of events consistent between each bin (>2500).
- Red represents the low-epsilon experimental data.
- Blue represents the low-epsilon experimental data.



$Q^2 = 3.85$, $W = 2.62$, $t = 0.21$ (2-epsilons)



t-binning range 0.16 – 0.50

Yield Calculations

- Calculated normalized bin-by-bin data yield.

$$Y_{EXP} = \frac{N}{Q_{eff}}$$

where,

Q_{eff}
= Charge × Tracking Eff × Detector Eff × EDTM Live Time
× Boiling Corr × other normalization factors ...

- Calculated normalized bin-by-bin MC yields.
- Calculated ratios (DATA/MC) for each t & phi-bin setting-by-setting, separately.

$$R(t, \varphi) = \frac{Y_{EXP}}{Y_{SIMC}}$$

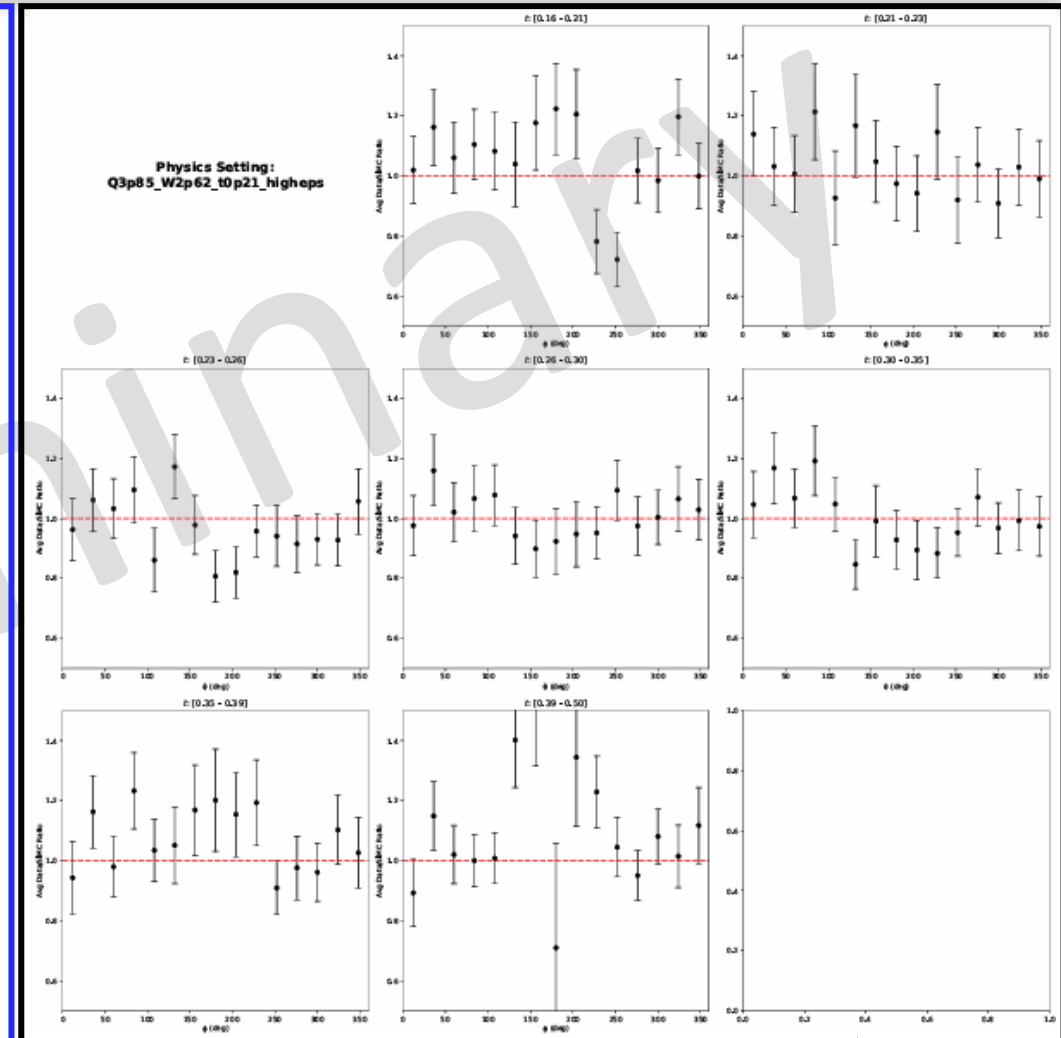
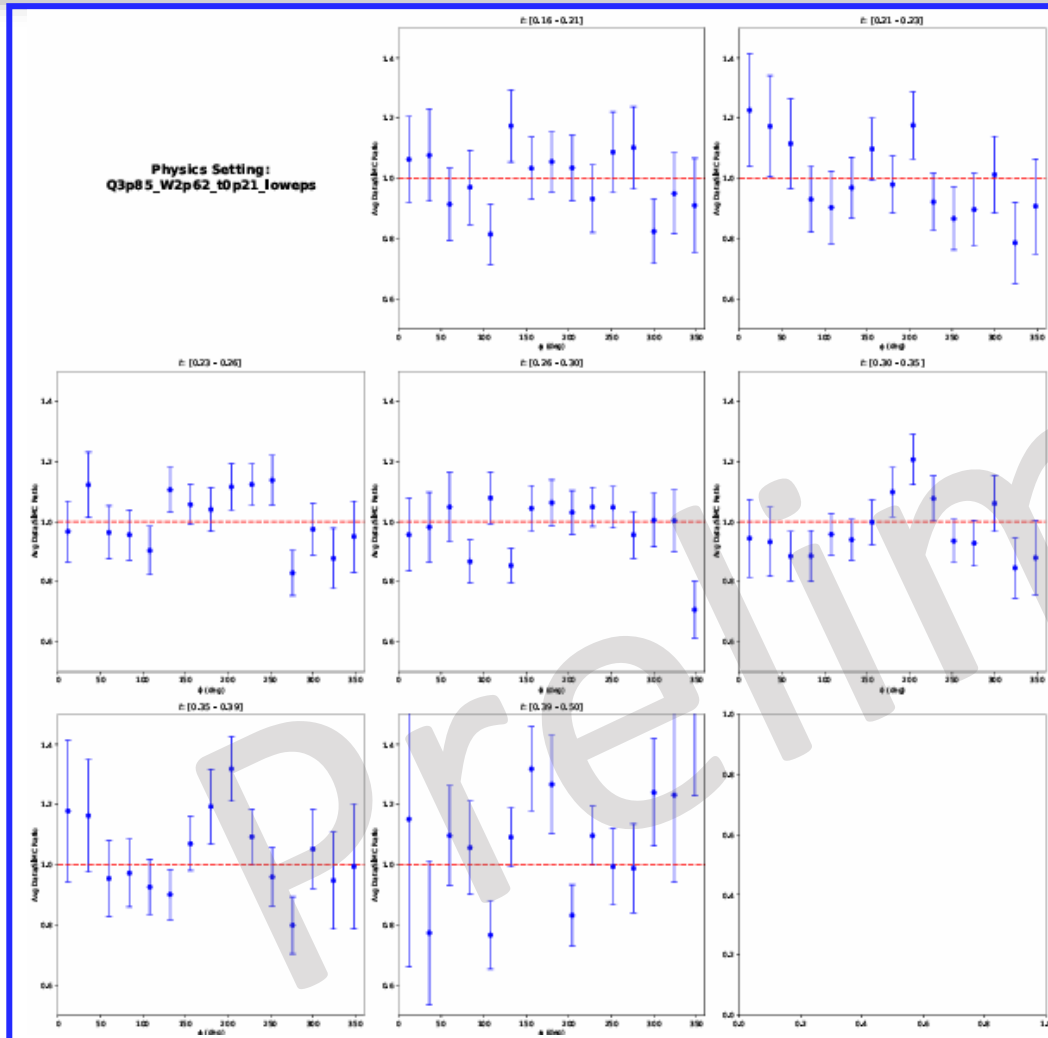
- Combined data from each pion arm angle setting per ϵ by calculating their error-weighted average.

Normalization Factors	Comments
Charge	Calculated run-by-run
HMS & SHMS Tracking Efficiencies	> 98%
Live Time Correction	>98%
HMS Cerenkov Efficiency	>99%
HMS Calorimeter Efficiency	>99%
SHMS Aerogel Efficiency	>98%
HMS & SHMS Hodoscope Efficiency	>98%
RF Efficiency	>99%
Boiling Correction Factor	Calculated run-by-run
Coin Blocking Correction	Calculated run-by-run
Pion Absorption Correction	~97%

High & Low-epsilon Data/MC Ratios

$$R(t, \varphi) = \frac{Y_{\text{EXP}}}{Y_{\text{SIMC}}}$$

Low_ε=0.292
High_ε=0.779

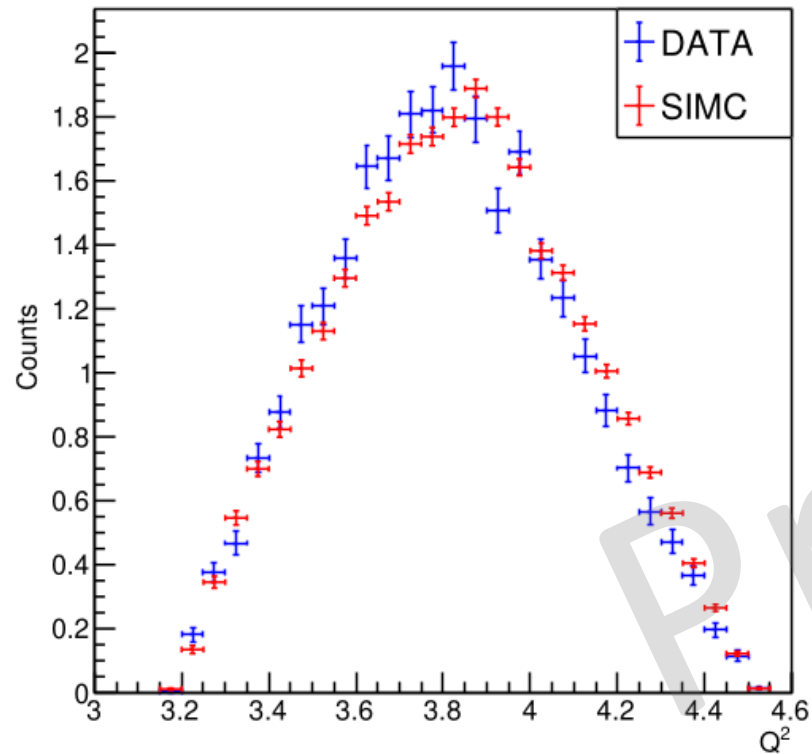


$Q^2 = 3.85 \text{ GeV}^2$, $W = 2.62 \text{ GeV}$ (low-epsilons)

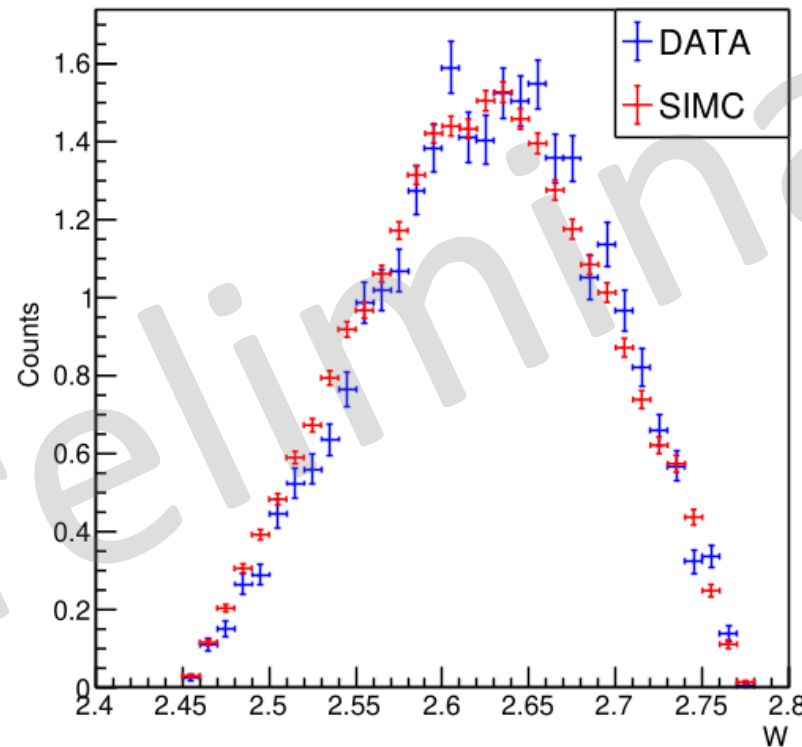
Data/MC Comparison Plots

- Compared kinematic and spectrometer variables between data and MC to verify that the SIMC model reliably reproduces the measured distributions.

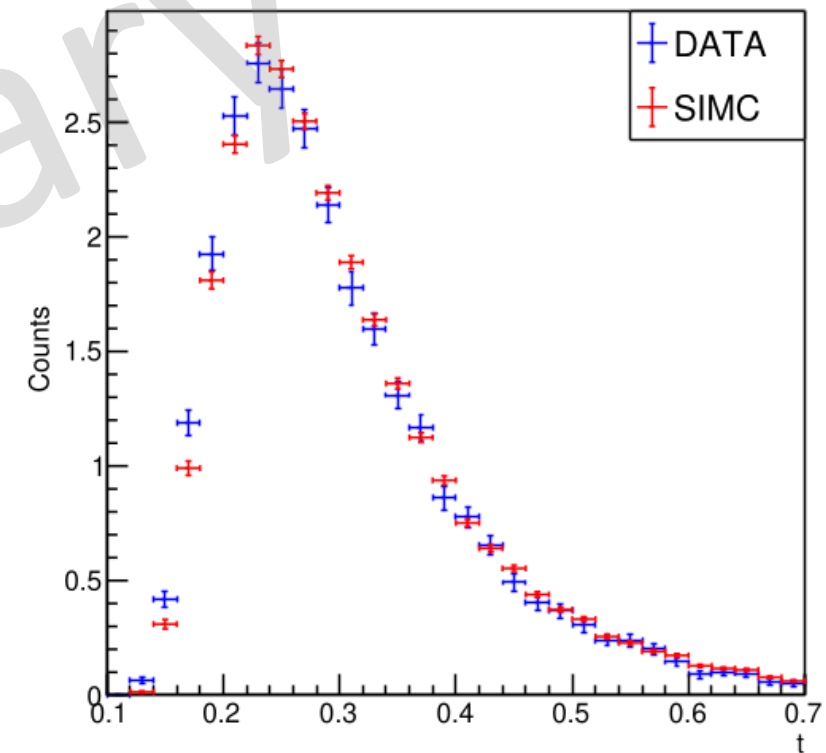
Q² Distribution



W Distribution



t Distribution



$Q^2 = 3.85 \text{ GeV}^2$, $W = 2.62 \text{ GeV}$ (low-epsilons Center Setting)

SIMC Empirical Model for Cross-sections

□ Final SIMC L/T/LT/TT functional forms:

$$\frac{d\sigma_T}{dt} = \left(\frac{p1}{Q^2} \right) \cdot e^{(p2 Q^2)} \cdot e^{(p3 |t|)}$$

$$\frac{d\sigma_L}{dt} = (p4 + p5/Q^2) \cdot \frac{|t|}{(|t| + m_\pi^2)^2} \cdot Q^2 e^{(p6|t|)} F_\pi^2$$

Where, $F_\pi = \frac{1}{(1+p7 \cdot Q^2 + p8 \cdot Q^4)}$

$$\frac{d\sigma_{LT}}{dt} = \left(\frac{p9}{Q^2} + e^{(p10/|t|)} \cdot \frac{p11}{|t|^{p12}} \right) \cdot \sin(\theta^*)$$

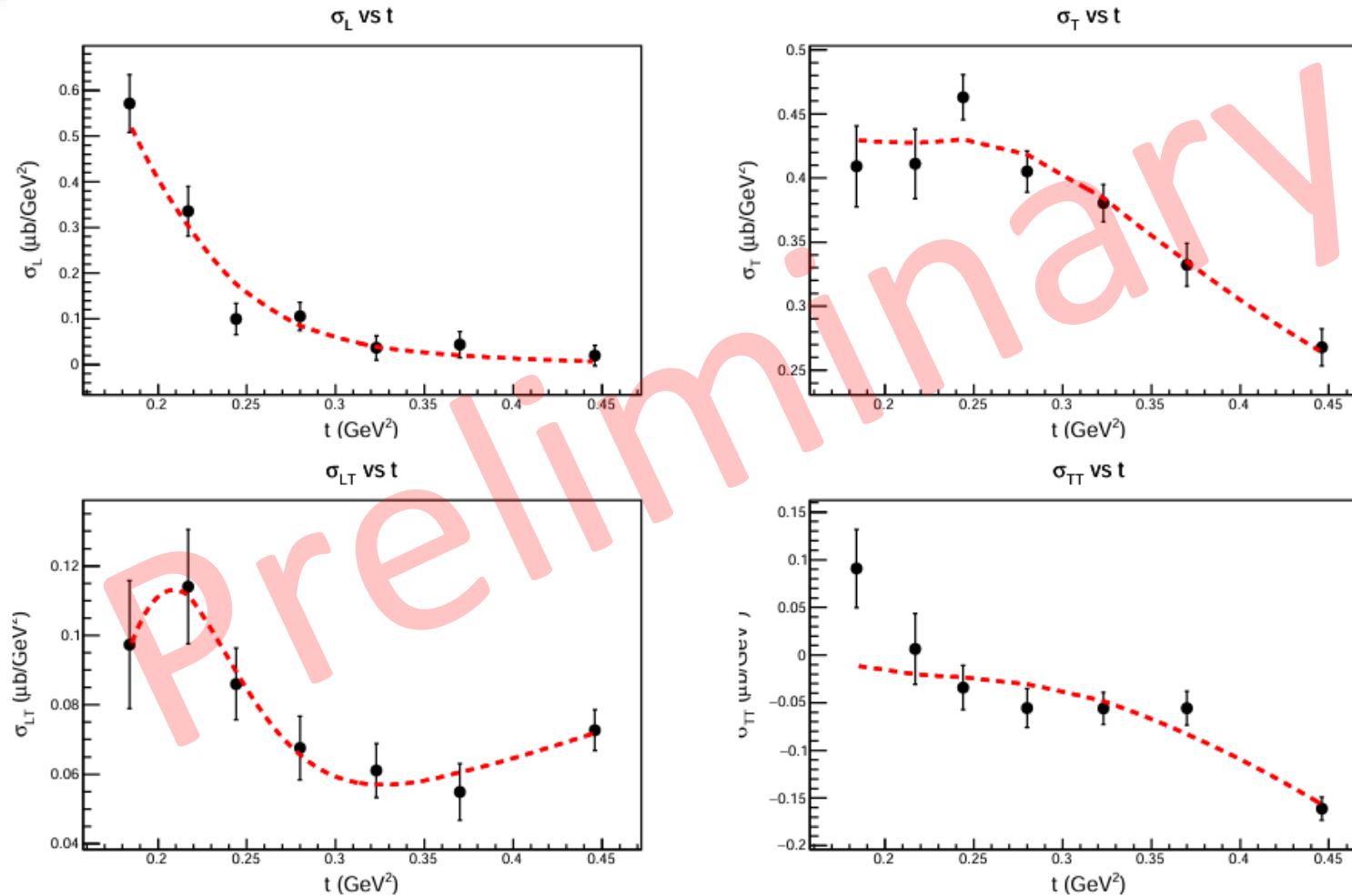
$$\frac{d\sigma_{TT}}{dt} = \left(\frac{p13}{Q^2} + \frac{p14}{|t|^{p15}} \right) \cdot \sin(\theta^*)^2$$

Parameter	Final Values	Parameter	Final Values
p1	20.93	p9	41.42
p2	0.11	p10	-4.8
p3	-2.55	p11	1.1e-7
p4	-2861.81	p12	26.98
p5	12153.41	p13	2696.06
p6	-9.22	p14	-661.19
p7	1.77	p15	0.14
p8	0.05		

□ Did 13 iterations to get optimized parameters.

Physics Setting: $Q^2=3.85 \text{ GeV}^2$, $W=2.62 \text{ GeV}$

L/T Separated Cross-sections



Physics Setting: $Q^2=3.85 \text{ GeV}^2$, $W=2.62 \text{ GeV}$

Model Comparison & F_π Extraction

Model Dependence & Future Extraction of F_π

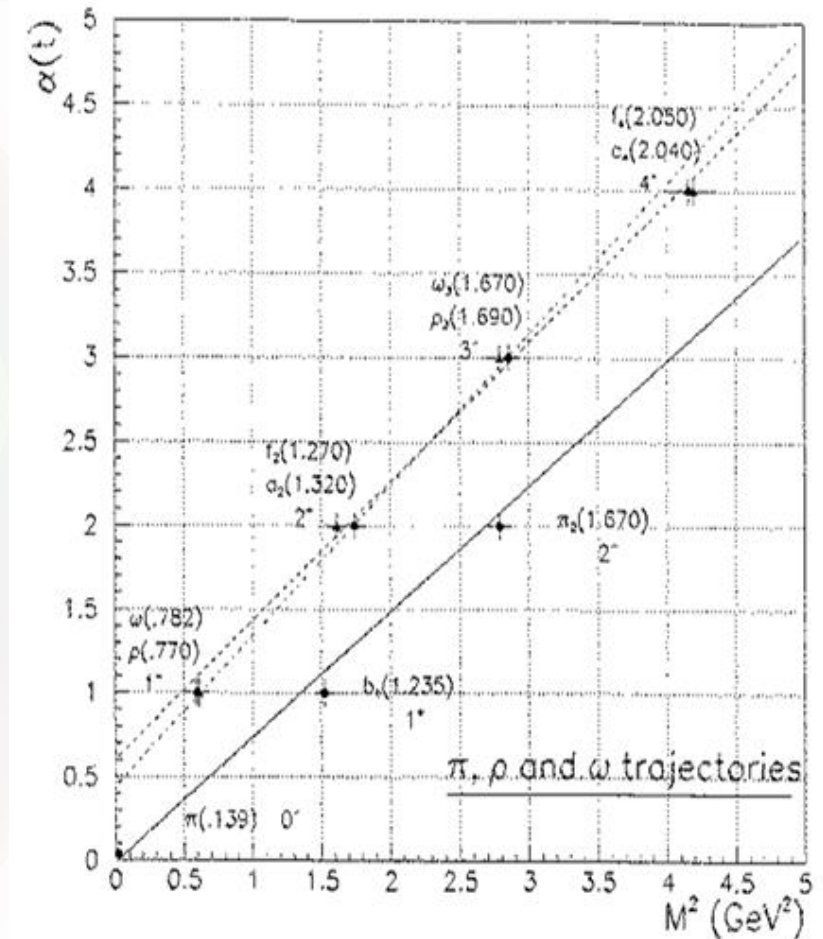
- ❑ Our philosophy is to publish our experimentally measured $\frac{d\sigma_L}{dt}$, so that updated values of $F_\pi(Q^2)$ can be extracted as improved theoretical models become available.
- ❑ Jefferson Lab F_π experiments have primarily used three models for form-factor extraction:
 - **VGL Regge Model** — Vanderhaeghen, Guidal, Laget, PRC 57, 1454 (1998)
 - VGL model has proven to give a **reliable description of σ_L** over a wide kinematic domain.
 - However, VGL **significantly underestimates σ_T** , indicating the need for improved models to fully describe the separated cross-section data.
 - **CKY Regge Model** — Choi, Kong, Yu, JKPS 67, L1089 (2015)
 - **This motivated the development of the CKY model**, aimed at improving the description of σ_T without degrading the successful description of σ_L .
 - **PKT Model** — Perry, Kizilersu, Thomas, PLB 807, 135581 (2020)
 - A **modern pion electroproduction model** was developed to improve theoretical consistency and study model dependence in the extracted **$F_\pi(Q^2)$** .
 - The PKT model provides an alternative framework for comparing how different reaction models affect the form-factor extraction.
- ❑ **More theoretical models are welcome**, especially those capable of describing both σ_L and σ_T consistently.

VGL Model to extract F_π from σ_L

- ❑ Feynman Propagator is given as; $\left[\frac{1}{t-m_\pi^2}\right]$ replaced by π and ρ Regge propagators.
- ❑ It represents the exchange of a series of particles, compared to a single particle.
- ❑ Model parameters fixed from pion photoproduction.
- ❑ Free parameters: Λ_π , Λ_ρ (trajectory cutoff).
- ❑ ρ exchange does not significantly influence σ_L at small $-t$.
- ❑ Pionform factor is a free parameter in the model, parameterized as:

$$F_\pi = \frac{1}{1 + Q^2/\Lambda_\pi}$$

Fit σ_L to the model which gives F_π at each Q^2 .



[Vanderhaeghen, Guidal, Laget, PRC 57(1998)1454]

Comparison with VGL Regge Model

VGL Regge Model

Phys Rev C 57(1998)1454

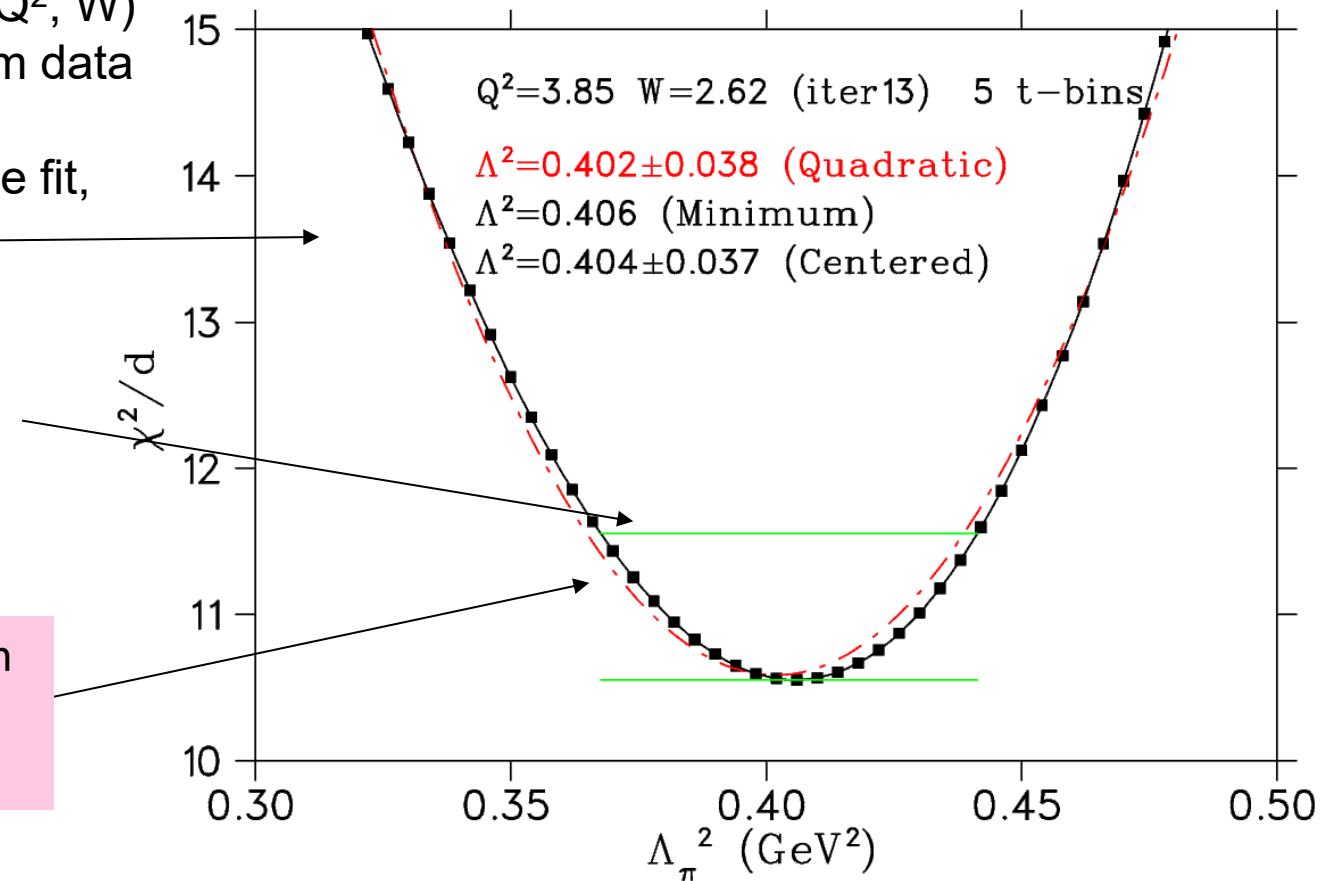
- The first model written for our data
- As the t-channel π exchange diagram is by itself not gauge invariant, the s-channel (for π^+ production) or u-channel (for π^- production) nucleon exchange diagram was also Reggeized, to ensure gauge invariance of their sum
- Two model parameters: σ_L sensitive to Λ_π^2 , much less sensitive to Λ_ρ^2

Determine best fit Λ_π^2 value

- Calculate VGL model for many Λ_π^2 using precise (Q^2 , W) values for each t-bin and compute χ^2 deviation from data
- Results depend on how many t-bins are used in the fit, #bins used shown here

Green line indicates increase in χ^2 by 1 unit, used to estimate in the uncertainty in Λ_π^2

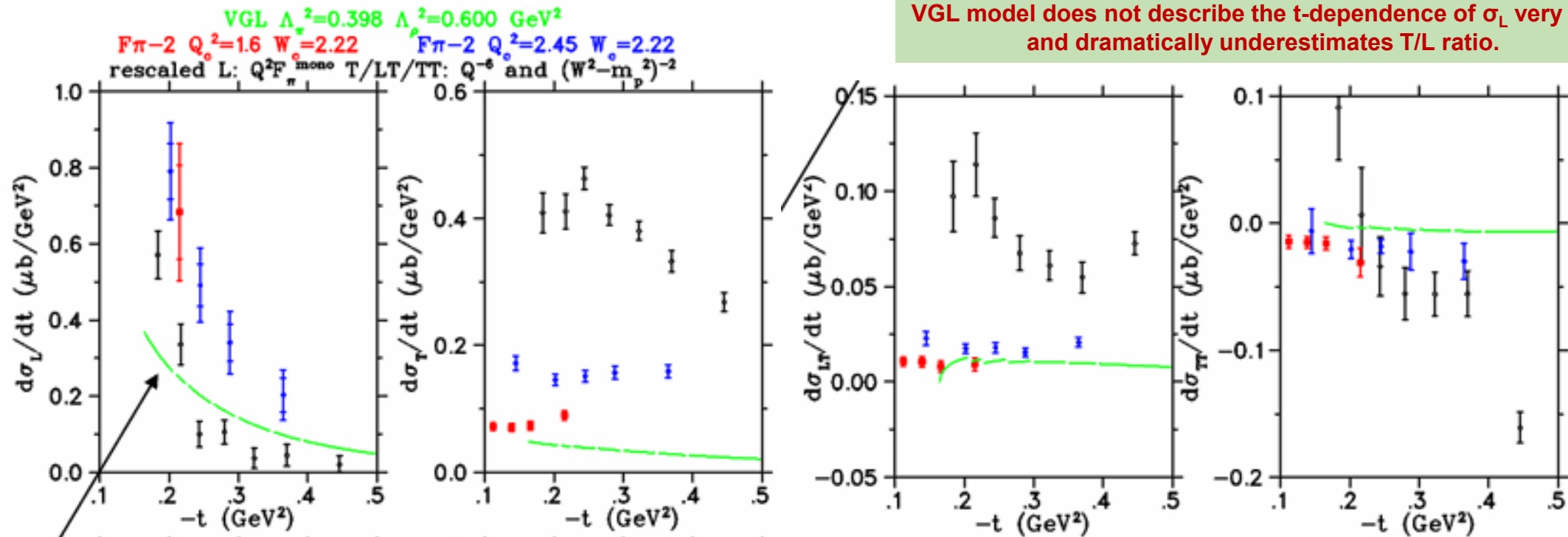
Red fit is a symmetric parabola, the difference from the black line indicates whether the uncertainty is asymmetric



$\Lambda_p^2=0.600$ for all calcs

Comparison with VGL Model

- Jefferson Lab F_π experiments use the Vanderhaeghen–Guidal–Laget (VGL) model as it has proven to give a reliable description of σ_L across a wide kinematic domain [Vanderhaeghen, Guidal, Laget, PRC 57(1998)1454].
- σ_L : results a little low compared to the scaled expectation from Fpi-2. If this is confirmed in our final analysis, it could indicate F_π is dropping more rapidly.
- σ_T : surprisingly large, an interesting rising trend from $Q^2=1.6$ -2.45-3.85.
- σ_{LT} & σ_{TT} : also much larger than Fpi-2.



Plots from Prof. Garth Huber

Model is evaluated at precise kinematics of data. Discontinuities indicate change in (Q^2, W) for each t -bin.

Comparison with CKY Regge Model

CKY Regge Model

J Korea PhysSoc 67(2015)L1089

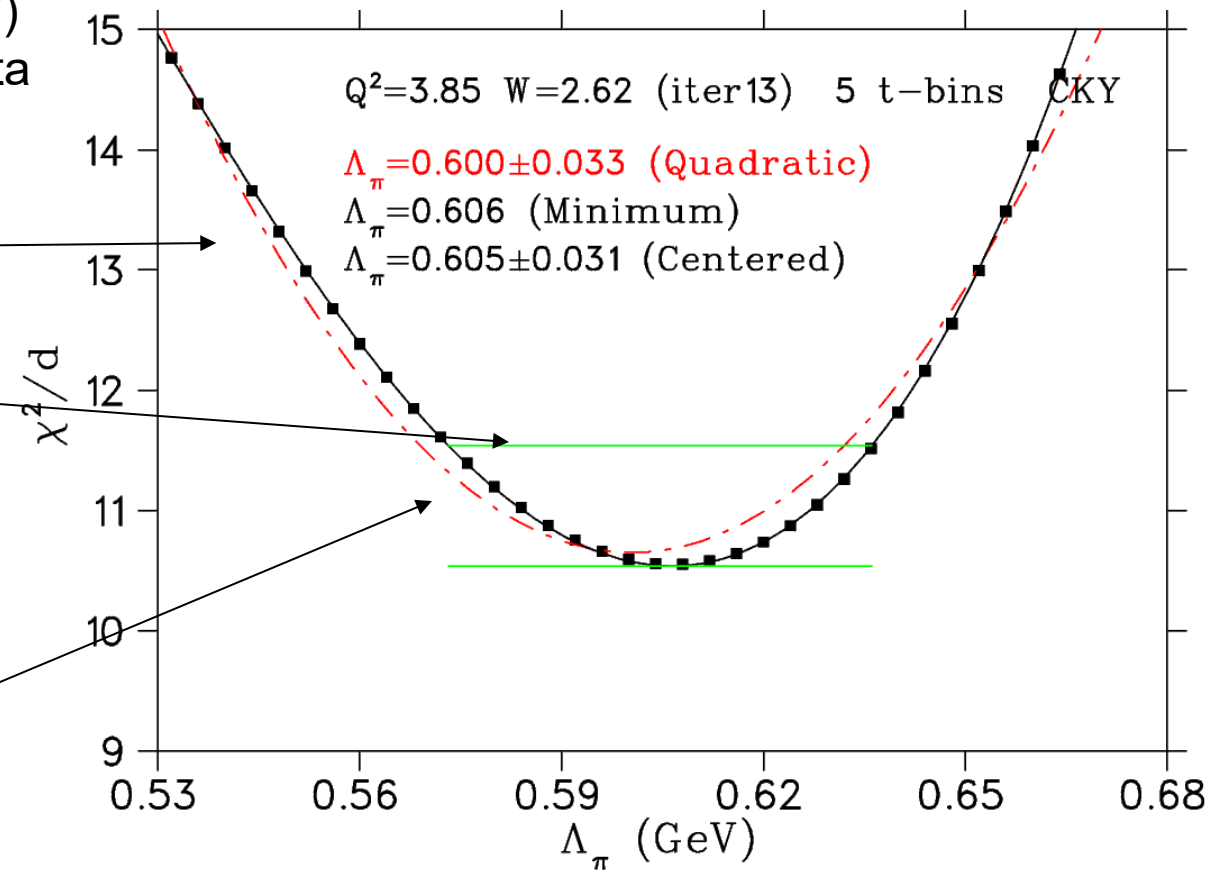
- The fact that VGL dramatically underestimates σ_T while doing a much better job on σ_L led to new model development, with the goal of increasing σ_T without degrading σ_L description
- The Kaskulov-Mosel and Vrancx-Ryckebusch models include intermediate (DIS-like) processes on the recoil nucleon, leading to the same final state. The CKY model is an extension of this work
- Third model parameter, Λ_N , for proton exchange in s-channel, is added to restore gauge invariance and plays a role in σ_T

Determine best fit Λ_π value

- Calculate CKY model for many Λ_π using precise (Q^2 , W) values for each t-bin and compute χ^2 deviation from data
- Results depend on how many t-bins are used in the fit, #bins used shown here

Green line indicates increase in χ^2 by 1 unit, used to estimate in the uncertainty in Λ_π

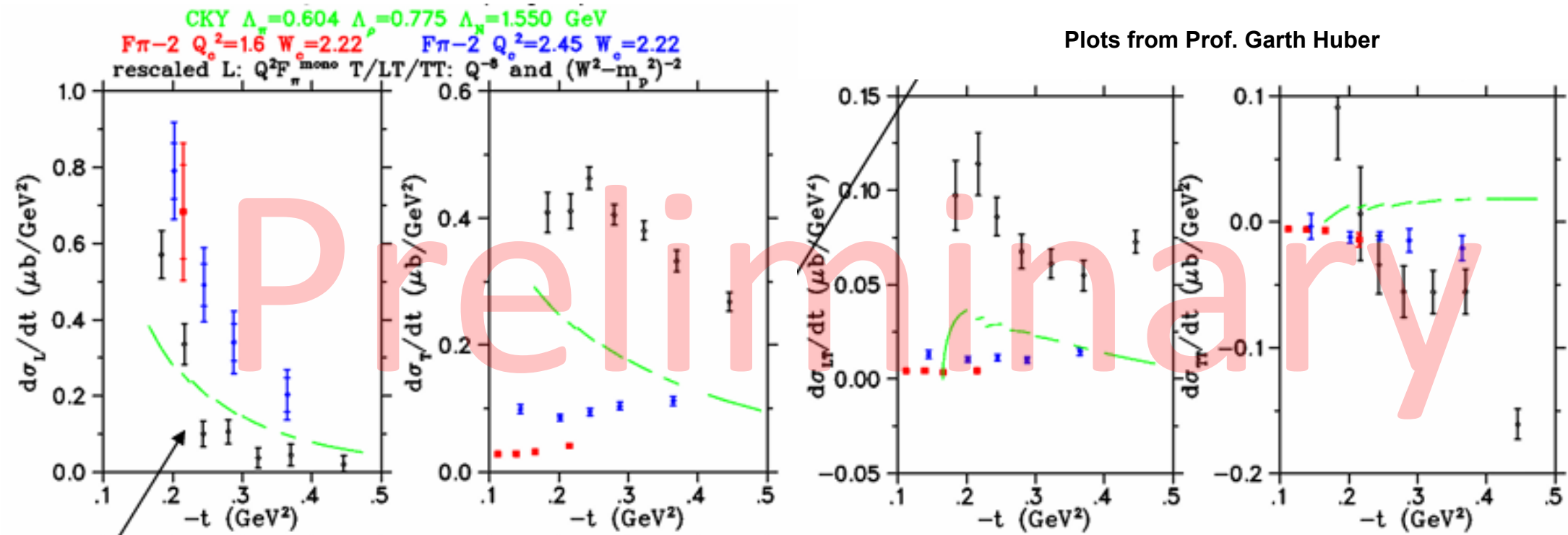
Red fit is a symmetric parabola, the difference from the black line indicates whether the uncertainty is asymmetric



$\Lambda_p=0.600$, $\Lambda_N=1.550$ for all calcs

Comparison with CKY Model

- The VGL dramatically underestimates σ_T while doing a much better job on σ_L led to new model development CKY, with the goal of increasing σ_T without degrading σ_L description [T.K. Choi, K.J. Kong, B.G. Yu, J.Kor.Phys.Soc. 67(2015) L1089; arXiv: 1508.00969].
- **CKY model also does not describe t-dependence of σ_L very well.**
- Much better T/L ratio, as expected, but still low.
- σ_{LT} somewhat higher than VGL, but still lower.



Model is evaluated at precise kinematics of data. Discontinuities indicate change in (Q^2, W) for each t -bin.