

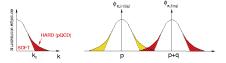
Outline

- Charged Meson Form Factors
- JLab 12 GeV data
- EIC Measurement and Simulation

Charged Meson Form Factors

- $\,\circ\,$ Simple $q\bar{q}$ valence structure of mesons makes them an excellent testing ground
- ullet Pion form factor , F_π , is the overlap integral -

$$F_{\pi}(Q^2) = \int \phi_{\pi}^*(p)\phi_{\pi}(p+q)dp$$

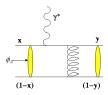


- Meson wave function can be split into $\phi_{\pi}^{\rm soft}$ $(k < k_0)$ and $\phi_{\pi}^{\rm hard}$, the hard tail
 - ullet Can treat $\phi_\pi^{
 m hard}$ in pQCD, cannot with $\phi_\pi^{
 m soft}$
- Study of Q^2 dependence of form factor focuses on finding description of hard and soft contributions

The Pion in pQCD 1/2

• At very large Q^2 , F_{π} can be calculated using pQCD via -

$$F_{\pi}(Q^2) = \frac{4_F \alpha_s(Q^2)}{Q^2} \Big| \sum_{n=0}^{\infty} a_n \left(\log \left(\frac{Q^2}{\Lambda^2} \right) \right)^{-\gamma_n} \Big|^2 \left[1 + O\left(\alpha_s(Q^2), \frac{m}{Q} \right) \right]$$



The Pion in pQCD 2/2

• At asymptotically high Q^2 ($Q^2 \to \infty$), the pion distribution amplitude becomes - $3f_{\pi}$

$$\phi_{\pi}(x)
ightarrow rac{3f_{\pi}}{\sqrt{n_c}} x(1-x)$$

- $\, \bullet \,$ With $f_\pi = 93$ MeV , the $\pi^+ \to \mu^+ \nu$ decay constant
- F_{π} takes the form -

$$Q^2 F_\pi \to 16\pi\alpha_s(Q^2) f_\pi^2$$

- This only relies on asymptotitc freedom in QCD, i.e. $(\partial \alpha_s/\partial \mu) < 0$ as $\mu \to \infty$
- Q^2F_{π} should behave as $\alpha_s(Q^2)$, even for moderately large Q^2
- Pion form factor seems to be the best tool for experimental study of the nature of the quark-gluon coupling constant renormalisation

Eqns - G.P. Lepage, S.J. Brodsky, PLB 87, p359, 1979 | Closing Statement - A.V. Efremov, A.V. Radyushkin PLB 94, p245, 1980

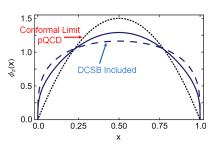
Implications for Pion Structure 1/2

• Previous pQCD derivation used normalisation of F_π based on the conformal limit of the pion's twist 2-PDA -

$$\phi_{\pi}^{cl}(x) = 6x(1-x)$$

- Gives F_{π} that are "too small"
- Incorporating the DCSB effects yields Pion PDA -

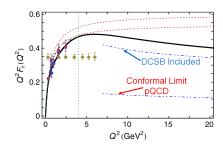
$$\phi_{\pi}(x) = \frac{8}{\pi} \sqrt{x(1-x)}$$



L. Chang, et al., PRL110(2013) 132001

Implications for Pion Structure 2/2

- Using this $\phi_{\pi}(x)$ in the pQCD expression brings the F_{π} calculation much closer to the data
- Underestimates the full computation by $\sim 15\%$ for $Q^2 \geqslant 8~GeV^2$



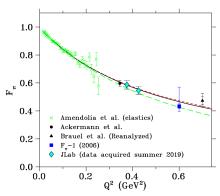
L. Chang, et al., PRL111(2013) 141802

Measurement of F_{π} - Low Q^2

- At low Q^2 , F_{π} can be measured model independently
 - \circ High energy elastic π^- scattering from atomic electrons in H
- CERN SPS used 300 GeV pions to measure F_{π} up to

 $Q^2 = 0.25 \ GeV^2$

- Used data to extract pion charge radius $r_{\pi} = 0.657 \pm 0.012$ fm
- Maximum accessible Q² approximately proportional to pion beam energy
 - $Q^2 = 1 \text{ GeV}^2$ requires 1 TeVpion beam (!)



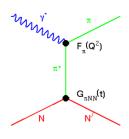
Amendolia, et al., NPB 277(1986) p168, P. Brauel, et al., ZPhysC (1979), p101, H. Ackerman, et al., NPB137 (1978), p294

Measurement of F_{π} - Larger Q^2

- To access higher Q^2 , must measure F_{π} indirectly
 - Use the "pion cloud" of the proton via pion electroproduction $p(e,e'\pi^+)n$
- At small -t, the pion pole process dominates the longitudinal cross section, σ_L
- In the Born term model, F_{π}^2 appears as -

$$rac{d\sigma_L}{dt} \propto rac{-tQ^2}{(t-m_\pi^2)} g^2(t) F_\pi^2(Q^2,t)$$

- Drawbacks of this technique -
 - Isolating σ_L experimentally challenging
 - Theoretical uncertainty in F_{π} extraction \rightarrow Model dependent

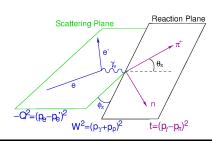


Measurement of F_{π} at JLab

 The physical cross section for the electroproduction process is given by -

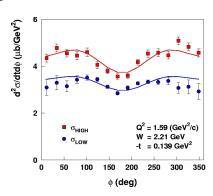
$$2\pi \frac{d^2\sigma}{dtd\phi} = \epsilon \frac{d\sigma_L}{dt} + \frac{d\sigma_T}{dt} + \sqrt{2\epsilon(\epsilon+1)} \frac{d\sigma_{LT}}{dt} \cos\phi + \epsilon \frac{d\sigma_{TT}}{dt} \cos 2\phi,$$
$$\epsilon = \left(1 + 2\frac{(E_e - E_{e'})^2 + Q^2}{Q^2} \tan^2\frac{\theta_{e'}}{2}\right)^{-1}$$

- $\circ \epsilon \rightarrow Virtual photon polarisation$
- L-T separation required to isolate σ_L from σ_T
- Need data at lowest -t possible, σ_L has maximum pole contribution here



Measuring $\frac{d\sigma_L}{dt}$ at JLab

- Rosenbluth separation required to isolate σ_L
 - Fix W, Q^2 and -t, measure cross section at two beam energies
 - \circ Carry out simultaneous fit at two different ϵ values to determine interference terms
- Careful control of point-to-point systematics crucial, $1/\Delta\epsilon$ error amplification in σ_L
- Spectrometer acceptance, kinematics and efficiencies must all be carefully studied and understood



T. Horn, et al., PRL 97(2006) 192001

Extracting F_{π} at JLab

- Only reliable approach for extracting F_{π} from σ_L is to use a model that incorporates the π^+ production mechanism and the spectator nucleon
- ullet JLab F_π experiments use the VGL Regge model
 - Reliably describes σ_L across a wide kinematic domaon
- Ideally, want a better understanding of the model dependence of the result
- There has been considerable recent interest
 - T.K. Choi, K.J. Kong, B.G. Yu, arXiv 1508.00969
 - T. Vrancx, J. Ryckebusch, PRC 89(2014)025203
 - M.M. Kaskulov, U. Mosel, PRC 81(2010)045202
 - S.V. Goloskokov, P.Kroll, EPJC 65(2010)137
- We aim to publish our experimentally measured cross section data so that updated values of F_{π} can be extracted as the models improve

VGL - Vanderhaeghen-Guidal-Laget Model - Vanderhaeghen, Guidal, Laget, PRC 57(1998) 1454

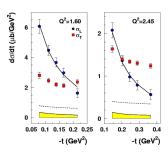
$F_{\pi}(Q^2)$ from JLab Data

VGL model incorporates π^+ production mechanism and spectator neutron effects

- Feynman propagator $\frac{1}{t-m^2}$ replaced by π and ρ Regge propagators
- Represents the exchange of a series of particles, compared to a single particle
- Free parameters Λ_{π} , Λ_{o} -Trajectory cutoff parameters
- At small -t, σ_I only

sensitive to
$$F_{\pi}$$

$$F_{\pi} = \frac{1}{1 + Q^2/\Lambda_{\pi}^2}$$

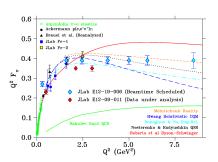


Error bars indicate statistical and random (pt-pt) systematic uncertainties in quadrature. Yellow band indicates the correlated (scale) and partly correlated (t-corr) systematic uncertainties. $\Lambda_{\pi}^{2} = 0.513, 0.491 \text{ GeV}^{2}, \Lambda_{\rho}^{2} = 1.7 \text{ GeV}^{2}$

T. Horn, et al., PRL 97(2006) 192001

Current and Projected JLab F_{π} Data

- JLab 12 GeV program includes measurements of F_{π} to higher Q^2
- No other facility worldwide can perform this measurement
- New overlap points at $Q^2=1.6, 2.45$ will be closer to pole to constrain $-t_{min}$ dependence
- Check π^+/π^- ratios at modest Q^2 to test t-channel dominance



• New low Q^2 point will provide best comparison of the electroproduction extraction of F_{π} vs elastic $\pi + e$ data

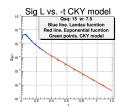
DEMP Studies at the EIC

- Measurements of the $p(e,e'\pi^+n)$ reaction at the EIC have the potential to extend the Q^2 reach of F_π measurements even further
- A challenging measurement however
 - Need good identification of $p(e, e'\pi^+n)$ triple coincidences
 - \circ Conventional L-T separation not possible \to would need lower than feasible proton energies to access low ϵ
- Utilise new EIC software framework to assess the feasibility of the study with updated design parameters
 - Feed in events generated from a DEMP event generator

DEMP Event Generator

- Want to examine exclusive reactions
 - $p(e, e'\pi^+n)$ exclusive reaction is reaction of interest $\rightarrow p(e, e'\pi^+)X$ SIDIS events are background
- Generator uses Regge-based $p(e, e'\pi^+)n$ model from T.K. Choi, K.J. Kong and B.G. Yu (CKY) arXiv 1508.00969
 - MC event generator created by parametrising CKY σ_L , σ_T for $5 < Q^2 < 35$, 2 < W < 10, 0 < -t < 1.2

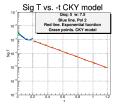


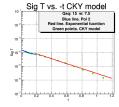


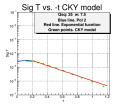


DEMP Event Generator

- Want to examine exclusive reactions
 - $p(e, e'\pi^+n)$ exclusive reaction is reaction of interest $\rightarrow p(e, e'\pi^+)X$ SIDIS events are background
- Generator uses Regge-based $p(e, e'\pi^+)n$ model from T.K. Choi, K.J. Kong and B.G. Yu (CKY) arXiv 1508.00969
 - MC event generator created by parametrising CKY σ_L , σ_T for $5 < Q^2 < 35$, 2 < W < 10, 0 < -t < 1.2

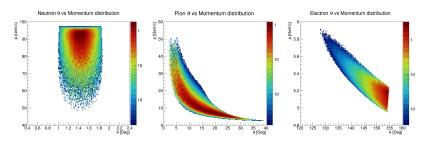






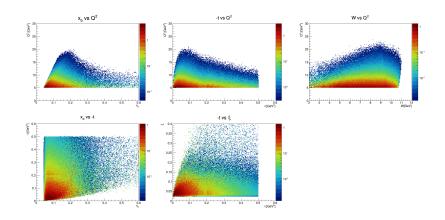
DEMP Acceptance for $-t < 0.5 \ GeV^2$

- $5(e^-)$ on 100(p) GeV collisions, 25 mrad crossing angle
- Events weighted by cross section
- No smearing



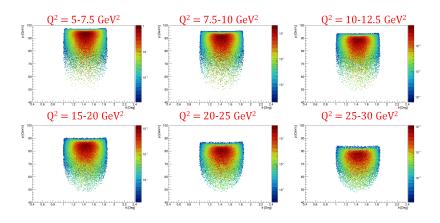
• Neutrons within 0.2° of outgoing proton beam, offset is due to the crossing angle (25 mrad $\approx 1.4^\circ$)

DEMP Kinematic Coverage - 5 on 100

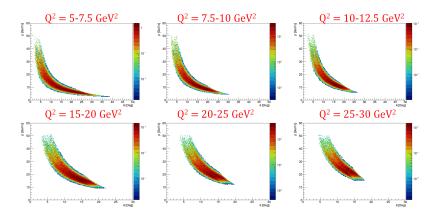


 $\xi = \text{skewnees}$

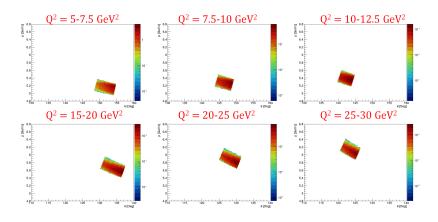
Neutron Acceptance Across Q^2 - 5 on 100



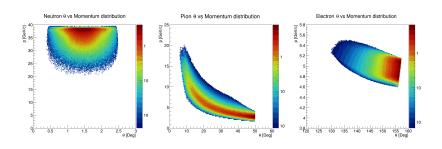
Pion Acceptance Across Q^2 - 5 on 100



Electron Acceptance Across Q^2 - 5 on 100

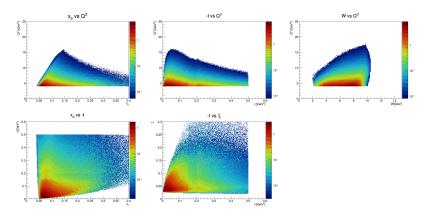


DEMP Acceptance - 5 on 41



- $Q^2 > 4$ GeV^2 cut applied, low Q^2 events dominate otherwise • High weight on low Q^2 events
- Neutron distribution broader in θ
 - May miss ZDC? Need to run full simulation and see

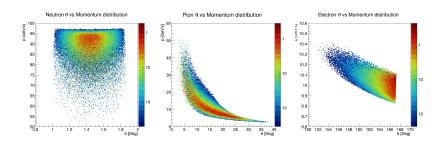
DEMP Kinematic Coverage - 5 on 41



 $\xi = \text{skewnees}$

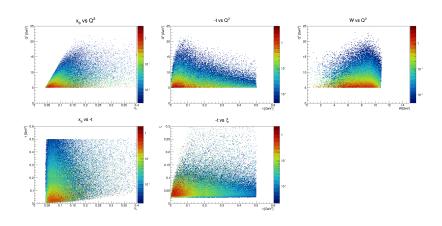
- $Q^2 > 4 \text{ GeV}^2$ cut applied
- Similar kinematic coverage to 5 on 100

DEMP Acceptance - 10 on 100



- Distributions broadly similar to 5 on 100
- Fewer events
- Electrons at higher momentum and wider angle

DEMP Kinematic Coverage - 10 on 100



 $\xi = \text{skewnees}$

ullet Similar to other energies, events shifted to higher W

Background Events

- Want to isolate a clean sample of $p(e, e'\pi^+n)$ events by detecting the neutron
- SIDIS $p(e, e'\pi^+)X$ events a large source of background
 - Utilised the EIC SIDIS event generator by Duke University to generate SIDIS background events
 - /work/eic/evgen_DUKE/e5p100 on the JLab farm
- Both the DEMP and SIDIS generators produce LUND format files that can be interpreted within the EIC software container

DEMP vs SIDIS Kinematics

- DEMP events are $e'\pi^+n$ triple coincidence
- SIDIS events are $e'\pi^+$ double coincidence, p_{miss} reconstructed

$$p_{miss} = |\underline{p}_e + \underline{p}_p - \underline{p}_{e'} - \underline{p}_{\pi^+}|$$
 Mis mom. SIDIS results in the side of the s

- SIDIS events overwhelm foreground exclusive events, but distributed over wider momentum range and at larger -t
- Note Plots from earlier study with smearing included

Isolating σ_L from σ_T in an e-p Collider

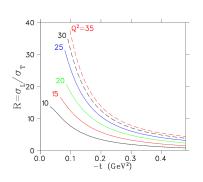
For a collider -

$$\epsilon = \frac{2(1-y)}{1+(1-y)^2}$$
 with $y = \frac{Q^2}{x(s_{tot} - M_N^2)}$

- y is the fractional energy loss
- ullet Systematic uncertainties in σ_L magnified by $1/\Delta\epsilon$
 - \bullet Ideally, $\Delta\epsilon > 0.2$
- To access $\epsilon < 0.8$ with a collider, need y > 0.5
 - \circ Only accessible at small s_{tot}
 - \circ Requires low proton energies ($\sim 10~GeV$), luminosity too low
- Conventional L-T separation not practical, need another way to determine σ_I

σ_I Isolation with a Model

- QCD scaling predicts $\sigma_L \propto Q^{-6}$ and $\sigma_T \propto Q^{-8}$
- At the high Q^2 and W accessible at the EIC, phenomenological models predict $\sigma_L\gg\sigma_T$ at small -t
- Can attempt to extract σ_L by using a model to isolate dominant $d\sigma_L/dt$ from measured $d\sigma_{UNS}/dt$
- Critical to confirm the validity of the model used!



Predictions are assuming $\epsilon > 0.9995$ with the kinematic ranges seen earlier

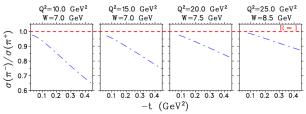
T.Vrancx, J. Ryckebusch, PRC 89(2014)025203

Model Validation via π^-/π^+ ratios

- Measure exclusive ${}^2H(e,e'\pi^+n)n$ and ${}^2H(e,e'\pi^-p)p$ in same kinematics as $p(e,e'\pi^+n)$
- π t-channel diagram is purely isovector \rightarrow G-Parity conserved

$$R = \frac{\sigma [n(e, e'\pi^{-}p)]}{\sigma [p(e, e'\pi^{+}n)]} = \frac{|A_V - A_S|^2}{|A_V - A_S|^2}$$

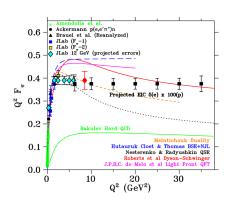
- R will be diluted if σ_T not small or if there are significant non-pole contributions to σ_L
- Compare R to model expectations



EIC Kinematic Reach

Assumptions

- $5(e^-)$ on 100(p)
- $\int \mathcal{L} = 20 \text{ fb}^{-1} \text{yr}^{-1}$
- Clean identification of $p(e, e'\pi^+n)$
- Syst.Unc:2.5% pt-pt, 12% scale
- $R = \sigma_L/\sigma_T = 0.013 0.14$ at lowest -t from VR model
- $\delta R = R$ Syst.Unc in model subtraction to isolate σ_L
- π pole dominance at small -t confirmed in 2H π^+/π^- ratios



 Results look promising, but need further studies and further energy combinations

Outlook and Future Plans

- Higher Q^2 data on F_π vital for our understanding of hadronic physics
 - Pion properties connected to DCSB
 - F_{π} is our best hope of observing QCD's transition from confinement-dominated physics to perturbative QCD
- Measurement of F_{π} at the EIC will be challenging
 - Conventional L-T separation not possible
 - Should be possible to use a model to separate σ_L from the unseparated cross section
 - Can use π^-/π^+ ratio in e+d collisions to validate model
 - Process files through full geant simulation, process other beam energy combinations, contribute to yellow report
- Building on our current event generator, may examine possibility of a Kaon event generator based on VR model
 - Could attempt to measure F_K in a similar manner
 - Further challenges to address for such a study!

Thanks for listening, any questions?



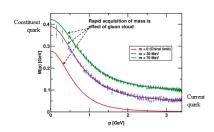


S.J.D. Kay, G.M. Huber, Z. Ahmed, Daniele Binosi, Huey-Wen Lin, Timothy Hobbs, Arun Tadepalli, Rachel Montgomery, Paul Reimer, David Richards, Rik Yoshida, Craig Roberts, Thia Keppel, John Arrington, Lei Chang, Ian L. Pegg, Jorge Segovia, Carlos Ayerbe Gayoso, Wenliang Li, Yulia Furletova, Dmitry Romanov, Markus Diefenthaler, Richard Trotta, Tanja Horn, Rolf Ent, Tobias Frederico

This research was supported by the Natural Sciences and Engineering Research Council of Canada (NSERC), FRN: SAPIN-2016-00031

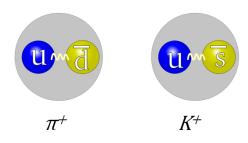
Recent Theoretical Advances

- Have a much better understanding of how Dynamical Chiral Symmetry Breaking (DCSB) generates hadron mass
- Evolution of the current-quark of pQCD into constituent quark was observed as its momentum becomes smaller
- The constituent quark mass arises from a cloud of low momentum gluons attaching themselves to the current quark
- Non-perturbative effect that generates a quark mass from nothing, occurs in even in the chiral (m = 0) limit



M.S. Bhagwat, et al., PRC 68(2003) 015203, L. Chang, et al., Chin.J.Phys. 49(2011)955

A $2^{\rm nd}$ Test Case - The Charged Kaon



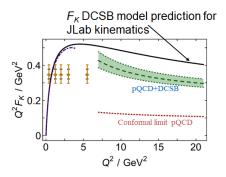
• In the hard scattering limit, pCQD predicts F_{π} and F_{K} will behave similarly -

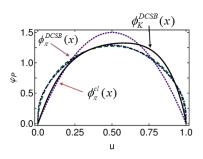
$$\frac{F_K(Q^2)}{F_\pi(Q^2)} \to \frac{f_K^2}{f_\pi^2}$$

• Should compare the magnitude and Q^2 dependences of both form factors

Effects of DCSB on K^+ Properties

- \circ K^+ PDA is also broad, concave and asymmetric
- Heavier s quark carries more bound state momentum than the u quark, shift is less then one might expect based on the difference in current quark masses.





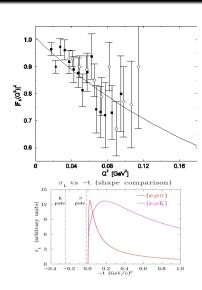
C. Shi, et al., PRD 92 (2015) 014035, F. Guo, et al., PRD 96(2017) 034024 (Full calculation)

F_K Measurement at JLab

- Similar to F_{π} , elastic K^+ scattering from electrons used to determine F_K at low Q^2
- Can "kaon cloud" of the proton be used in the same way as the pion to extract F_k from electroproduction?
- Kaon pole further from kinematically allowed region

$$\frac{d\sigma_L}{dt} \propto \frac{-tQ^2}{(t-m_K^2)}g_K^2(T)F_K^2(Q^2,t)$$

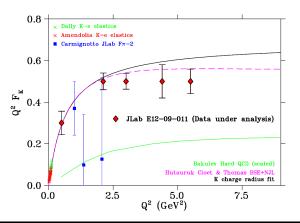
• Issues are being explored and tested in JLab E12-09-011



19/06/2020

F_K Measurement at JLab - Projections

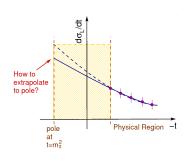
- Points with projected errors shown below
- Data has all been acquired and analysis is in progress
- y positioning of points arbitrary



Chew-Low Method to determine F_{π}

- $p(e,e'\pi^+)n$ data obtained away from $t=m_\pi^2$ pole
- "Chew Low" extrapolation method must know analytical dependence of $d\sigma_L/dt$ in unphysical region
- Extrapolation method last used in 1972 by Devenish and Lyth
- Very large systematic uncertainties
- Failed to produce a reliable result
- Different polynomial fits equally likely in physical region
 - Form factor values divergent when extrapolated

We do not use the Chew-Low method

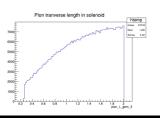


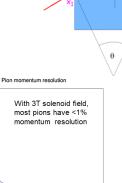
Old Momentum Resolution Estimate

Intrinsic momentum resolution from n equidistant measurements

$$\frac{\delta p}{p} = \frac{p}{0.3B} \frac{\sigma_{r\phi}}{L'^2} \sqrt{\frac{720}{n+4}}$$

- R. L. Glcukstern, NIM24(1963), p381
- B= central field (T), $\sigma_{r\phi}=$ position resolution (m), L= length of transverse path through field (m), N= number of measurements
- Assumed n = 5, B = 3 T





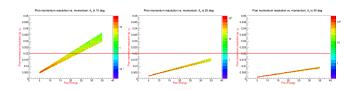
⊗B

40000

35000

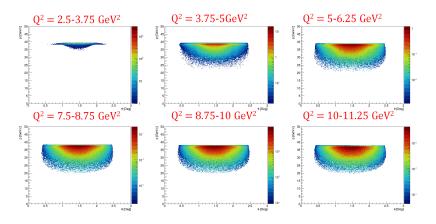
0.01 0.02 0.03 0.04 0.05 0.06 0.07 0.08 0.09

Old π Momentum Resolution with 3 T Solenoid

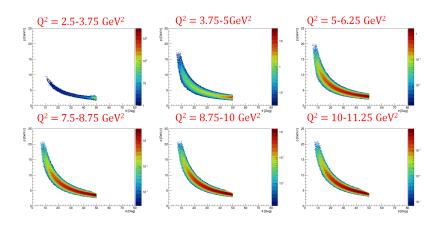


- Pion momentum resolution suffers when the pion is emitted at a shallow angle to the solenoidal field
- \bullet To simplify the MC study, assumed $\delta p/p=2\%$ for all angles, for both pion and electron
 - Typical π^+ angles: $7-30^\circ$
 - \circ Typical e^- angles: $25-45^\circ$

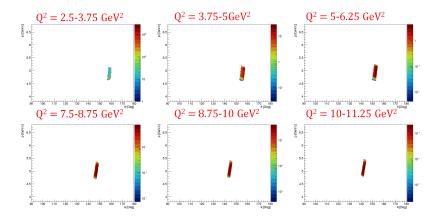
Neutron Acceptance Across Q^2 - 5 on 41



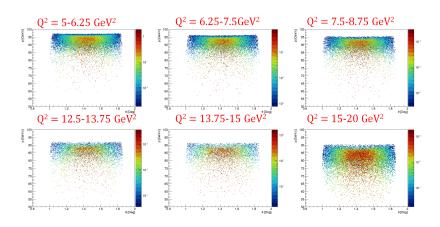
Pion Acceptance Across Q^2 - 5 on 41



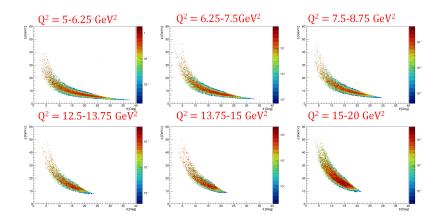
Electron Acceptance Across Q^2 - 5 on 41



Neutron Acceptance Across Q^2 - 10 on 100



Pion Acceptance Across Q^2 - 10 on 100



Electron Acceptance Across Q^2 - 10 on 100

