

Meson and Nucleon Form Factors

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University
of Regina

Science at the Luminosity Frontier: JLab @ 22 GeV
December 9, 2024

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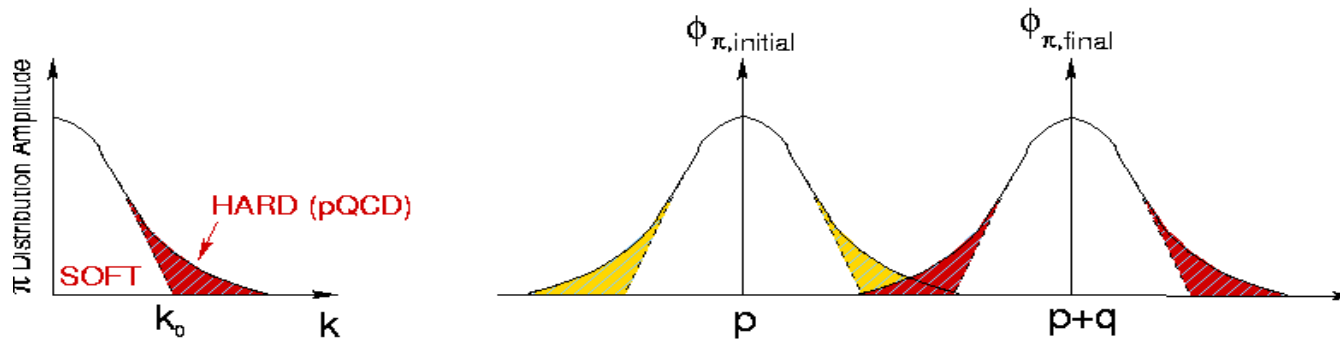


SAPIN-2021-00026

Simple $q\bar{q}$ valence structure of mesons presents the ideal testing ground for our understanding of bound quark systems.

In quantum field theory, the form factor is the overlap integral:

$$F_\pi(Q^2) = \int \phi_\pi^*(p) \phi_\pi(p+q) dp$$



The meson wave function can be separated into ϕ_π^{soft} with only low momentum contributions ($k < k_0$) and a hard tail ϕ_π^{hard} .

While ϕ_π^{hard} can be treated in pQCD, ϕ_π^{soft} cannot.

From a theoretical standpoint, the study of the Q^2 -dependence of the form factor focuses on finding a description for the hard and soft contributions of the meson wave-function.

A program of study unique to Hall C (until completion of EIC)

At large Q^2 , perturbative QCD (pQCD) can be used

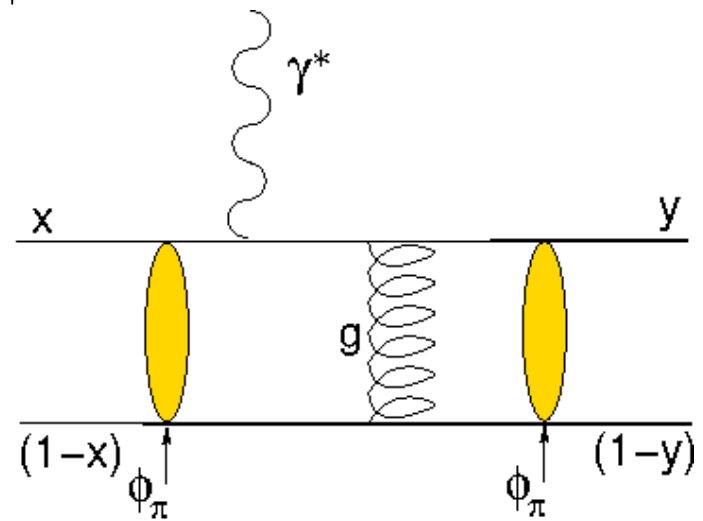
$$F_\pi(Q^2) = \frac{4\pi C_F \alpha_s(Q^2)}{Q^2} \left| \sum_{n=0}^{\infty} a_n \left(\log \left(\frac{Q^2}{\Lambda^2} \right) \right)^{-\gamma_n} \right|^2 \left[1 + O \left(\alpha_s(Q^2), \frac{m}{Q} \right) \right]$$

at asymptotically high Q^2 , only the hardest portion of the wave function remains

$$\phi_\pi(x) \xrightarrow{Q^2 \rightarrow \infty} \frac{3f_\pi}{\sqrt{n_c}} x(1-x)$$

and F_π takes the very simple form

$$F_\pi(Q^2) \xrightarrow{Q^2 \rightarrow \infty} \frac{16\pi\alpha_s(Q^2)f_\pi^2}{Q^2}$$



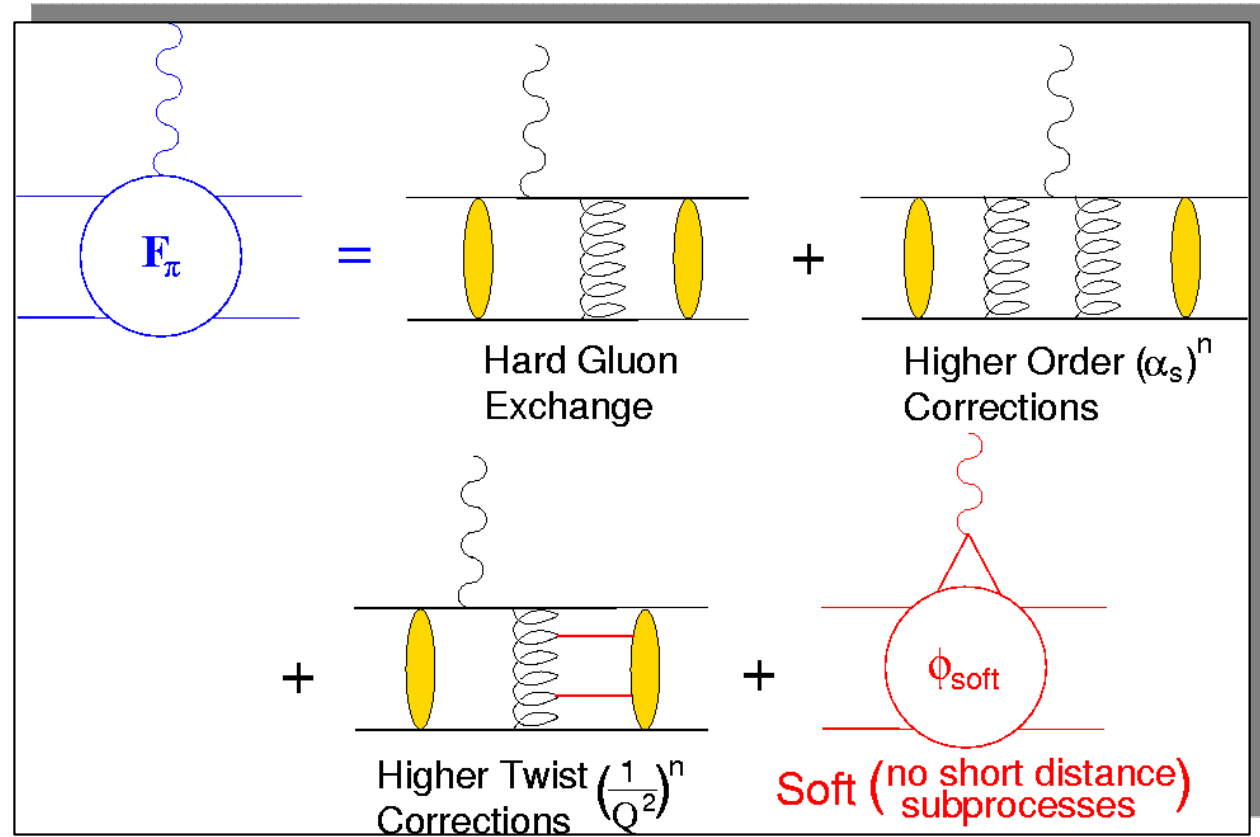
where $f_\pi=92.4$ MeV is the $\pi^+ \rightarrow \mu^+ \nu$ decay constant.

G.P. Lepage, S.J. Brodsky, Phys.Lett. **87B**(1979)359

This prediction only relies on asymptotic freedom in QCD, *i.e.* $(\partial\alpha_s/\partial\mu) < 0$ as $\mu \rightarrow \infty$

Pion Form Factor at Finite Q^2

- At finite momentum transfer, higher order terms contribute.
- Calculation of higher order, “hard” (short distance) processes difficult, but tractable.

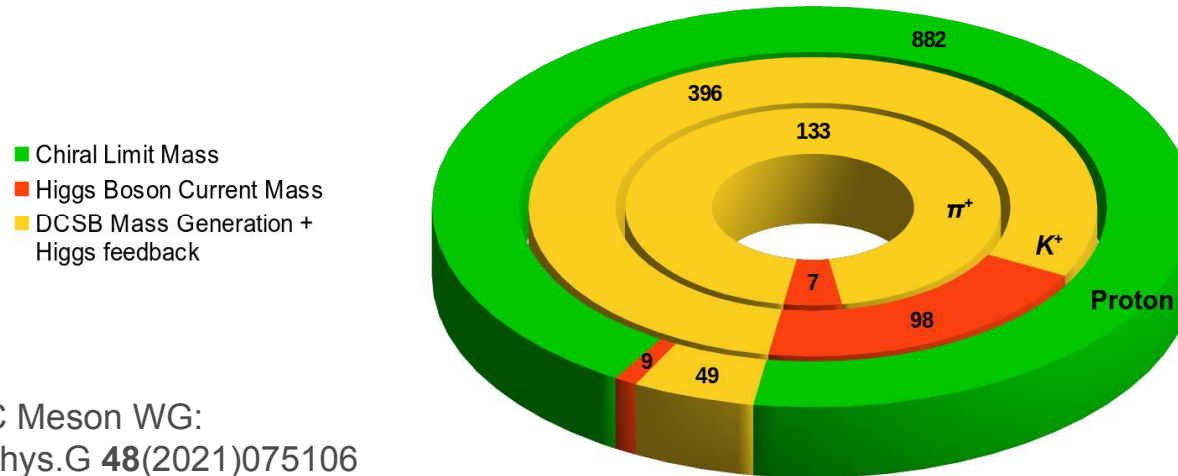


$Q^2 F_\pi$ should behave like $\alpha_s(Q^2)$ even for moderately large Q^2 .

→ Pion form factor seems to be best tool for experimental study of nature of the quark-gluon coupling constant renormalization.

[A.V. Radyushkin, JINR 1977, arXiv:hep-ph/0410276]

Hadron Mass Budget



EIC Meson WG:
J.Phys.G 48(2021)075106

Stark Differences between proton, K^+ , π^+ mass budgets

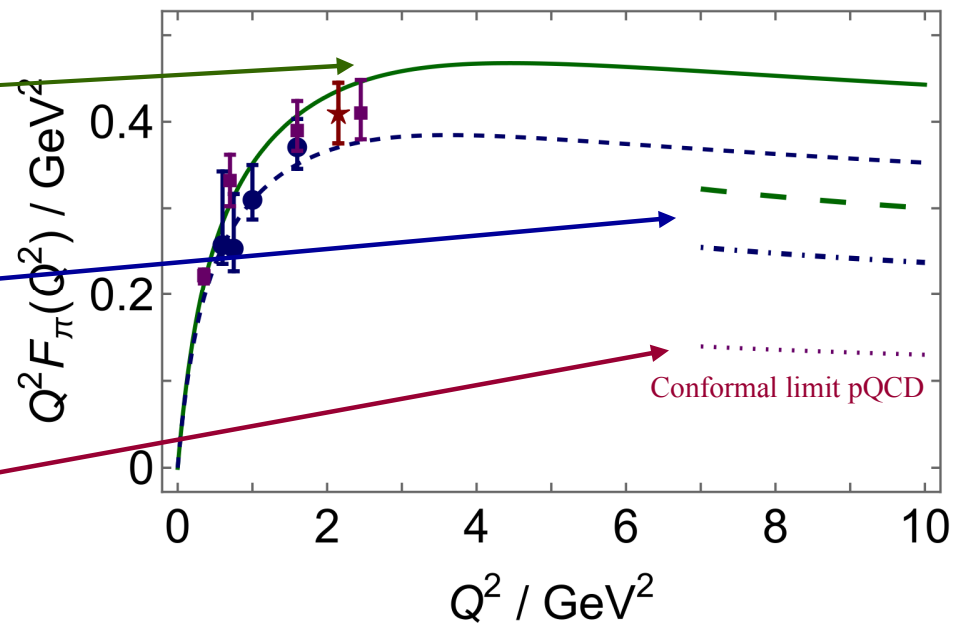
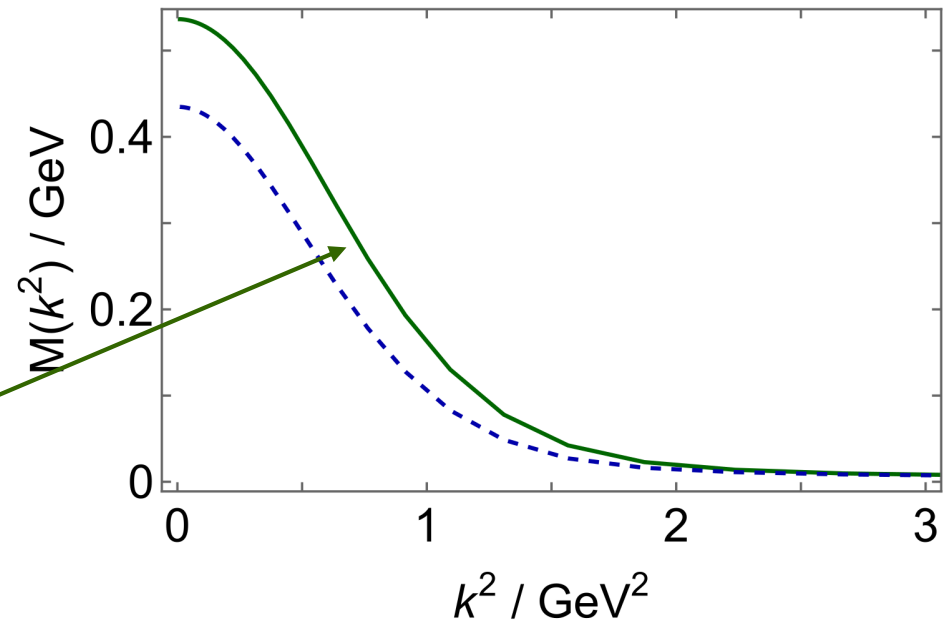
- Due to Emergent Hadronic Mass (EHM), Proton mass large in absence of quark couplings to Higgs boson (chiral limit).
- Conversely, and yet still due to EHM and DCSB, K and π are massless in chiral limit (i.e. they are Goldstone bosons of QCD).
- The mass budgets of these crucially important particles demand interpretation.
- Equations of QCD stress that any explanation of the proton's mass is incomplete, unless it simultaneously explains the light masses of QCD's Goldstone bosons, the π and K .

Synergy: Emergent Mass and π^+ Form Factor

At empirically accessible energy scales, π^+ form factor is sensitive to emergent mass scale in QCD

- Two dressed-quark mass functions distinguished by amount of DCSB
 - DCSB emergent mass generation is 20% stronger in system characterized by solid green curve, which is more realistic case
- $F_\pi(Q^2)$ obtained with these mass functions
 - $r_\pi=0.66$ fm with solid green curve
 - $r_\pi=0.73$ fm with solid dashed blue curve
- $F_\pi(Q^2)$ predictions from QCD hard scattering formula, obtained with related, computed pion PDAs
- QCD hard scattering formula, using conformal limit of pion's twist-2 PDA

$$\phi_\pi^{cl}(x) = 6x(1-x)$$



Why Meson Form Factors?

- **The π^+ form factor is our best hope of observing experimentally QCD's transition from soft QCD to hard QCD**
 - This transition is expected to occur at a much lower Q^2 than for the proton
- **K^+ form factor:**
 - How does meson structure change when s quark is substituted for d quark?
 - At what Q^2 will the K^+ to π^+ form factor ratio converge to the value predicted by QCD?
- **The normalization of π^+ and K^+ form factors at high Q^2 is sensitive to quark and gluon energy contributions to emergent hadronic mass**
 - A comparison of π^+ and K^+ form factors over a wide range of Q^2 will provide unique information relevant to our understanding of hadronic mass generation

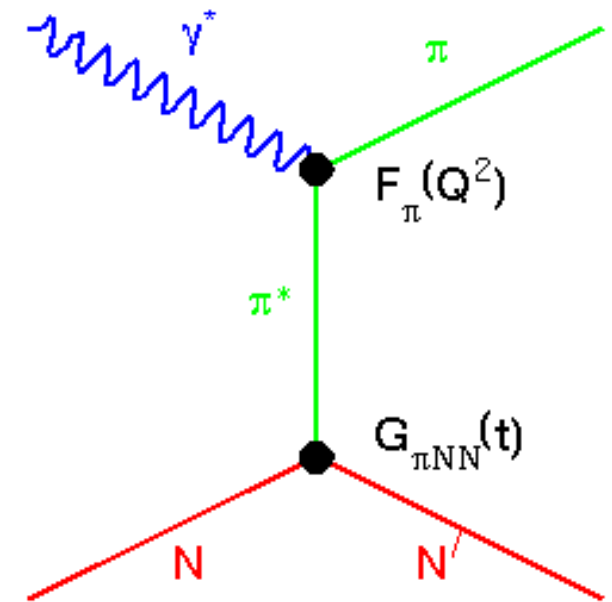
Measurement of π^+ Form Factor – Larger Q^2

At larger Q^2 , F_π must be measured indirectly using the “pion cloud” of the proton via pion electroproduction $p(e, e'\pi^+)n$

$$|p\rangle = |p\rangle_0 + |n\pi^+\rangle + \dots$$

- At small $-t$, the pion pole process dominates the longitudinal cross section, σ_L
- In Born term model, F_π^2 appears as,

$$\frac{d\sigma_L}{dt} \propto \frac{-tQ^2}{(t - m_\pi^2)} g_{\pi NN}^2(t) F_\pi^2(Q^2, t)$$



Drawbacks of this technique

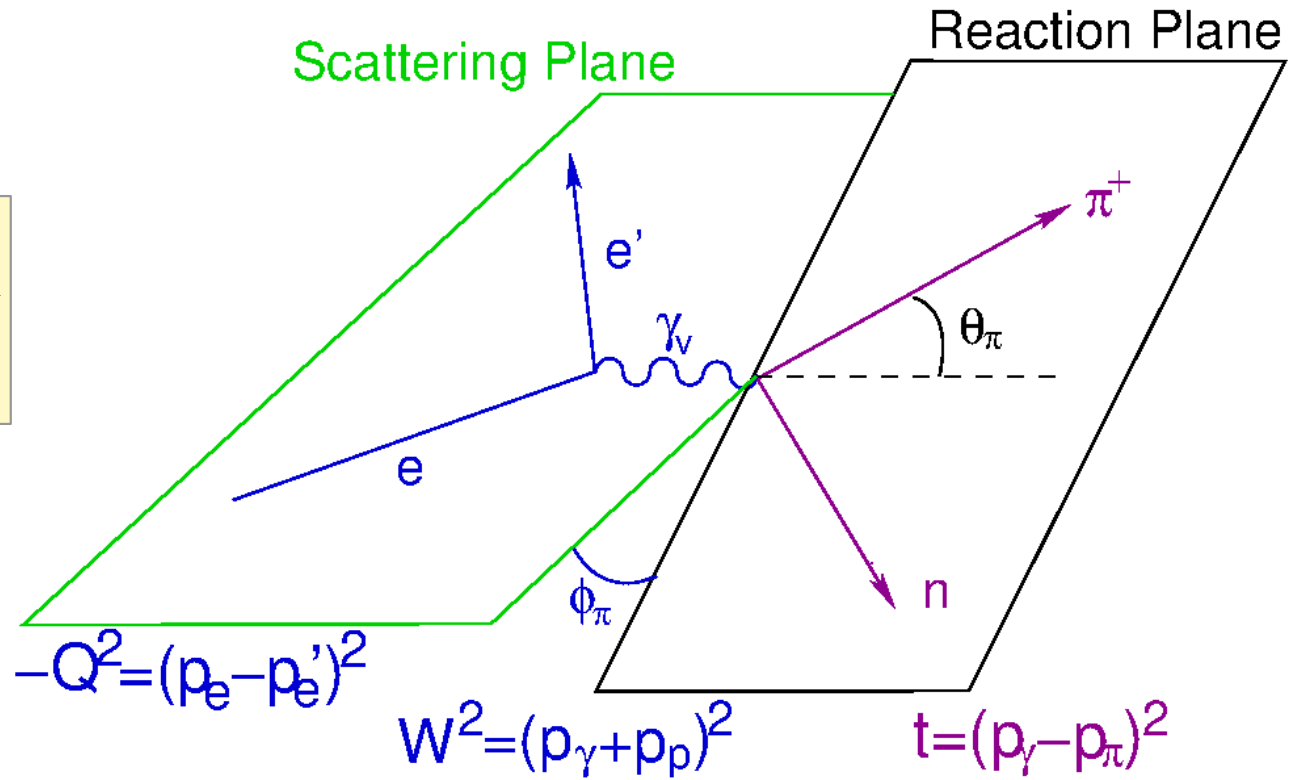
1. Isolating σ_L experimentally challenging
2. Theoretical uncertainty in form factor extraction.

K^+ pole is further in the unphysical region, uncertainties will be larger

$$2\pi \frac{d^2\sigma}{dt d\phi} = \varepsilon \frac{d\sigma_L}{dt} + \frac{d\sigma_T}{dt} + \sqrt{2\varepsilon(\varepsilon+1)} \frac{d\sigma_{LT}}{dt} \cos\phi + \varepsilon \frac{d\sigma_{TT}}{dt} \cos 2\phi$$

Virtual-photon polarization:

$$\varepsilon = \left(1 + 2 \frac{(E_e - E_{e'})^2 + Q^2 \tan^2 \frac{\theta_{e'}}{2}}{Q^2} \right)^{-1}$$



- L-T separation required to separate σ_L from σ_T
- Need to take data at smallest available $-t$, so σ_L has maximum contribution from the π^+ pole
- Need to measure t -dependence of σ_L at fixed Q^2, W

Error in $d\sigma_L/dt$ is magnified by $1/\Delta\varepsilon$, where $\Delta\varepsilon=(\varepsilon_{\text{Hi}}-\varepsilon_{\text{Low}})$

→ To keep magnification factor $<5\times$, need $\Delta\varepsilon>0.2$, preferably more!

$$\frac{d^2\sigma}{dt d\phi} = \varepsilon \frac{d\sigma_L}{dt} + \frac{d\sigma_T}{dt} + \sqrt{2\varepsilon(\varepsilon+1)} \frac{d\sigma_{LT}}{dt} \cos\phi_\pi + \varepsilon \frac{d\sigma_{TT}}{dt} \cos 2\phi_\pi$$

$$\frac{\Delta\sigma_L}{\sigma_L} = \frac{1}{(\varepsilon_1 - \varepsilon_2)} \left(\frac{\Delta\sigma}{\sigma} \right) \sqrt{(R + \varepsilon_1)^2 + (R + \varepsilon_2)^2} \quad \text{where } R = \frac{\sigma_T}{\sigma_L}$$

$$\frac{\Delta\sigma_T}{\sigma_T} = \frac{1}{(\varepsilon_1 - \varepsilon_2)} \left(\frac{\Delta\sigma}{\sigma} \right) \sqrt{\varepsilon_1^2 \left(1 + \frac{\varepsilon_2}{R} \right)^2 + \varepsilon_2^2 \left(1 + \frac{\varepsilon_1}{R} \right)^2}$$

The relevant quantities for F_π extraction are R and $\Delta\varepsilon$

$$\frac{d\sigma_L}{dt} \propto \frac{-tQ^2}{(t - m_\pi^2)} g_{\pi NN}^2(t) F_\pi^2(Q^2, t)$$

What is being measured?

- Scattered electron and π^+/K^+ in coincidence with the two high performance spectrometers in Hall C
 - High momentum, forward angle (5.5°) meson detection is required, with good Particle ID to separate π^+ , K^+ , p
 - Good momentum resolution required to reconstruct crucial kinematics, such as M_{miss} , Q^2 , W , t
 - Need to measure the longitudinal cross section $d\sigma_L/dt$ needed for form factor extraction

The role of 22 GeV electrons?

- Allows access to higher Q^2
- Expanded range of virtual photon polarization $\Delta\varepsilon=(\varepsilon_{\text{HI}}-\varepsilon_{\text{LO}})$, leading to reduced errors in the extraction of $d\sigma_L/dt$
 - Uncertainty in $\sigma_L \sim 1/\Delta\varepsilon$, desire $\Delta\varepsilon > 0.2$, preferably larger

Upgrade Scenarios Considered

Phase 1: higher energy beam, keep HMS+SHMS largely as is, with relatively small DAQ and PID upgrades

- See what can be accomplished in “cost effective approach”
- Goal: to extend kinematic range of L/T–separated measurements beyond what is possible with JLab 11 GeV beam

Phase 2: Replace HMS with a new Very High Momentum Spectrometer (VHMS) to enable measurements utilizing full 22 GeV beam energy

- See what extra physics can be obtained for significantly larger investment

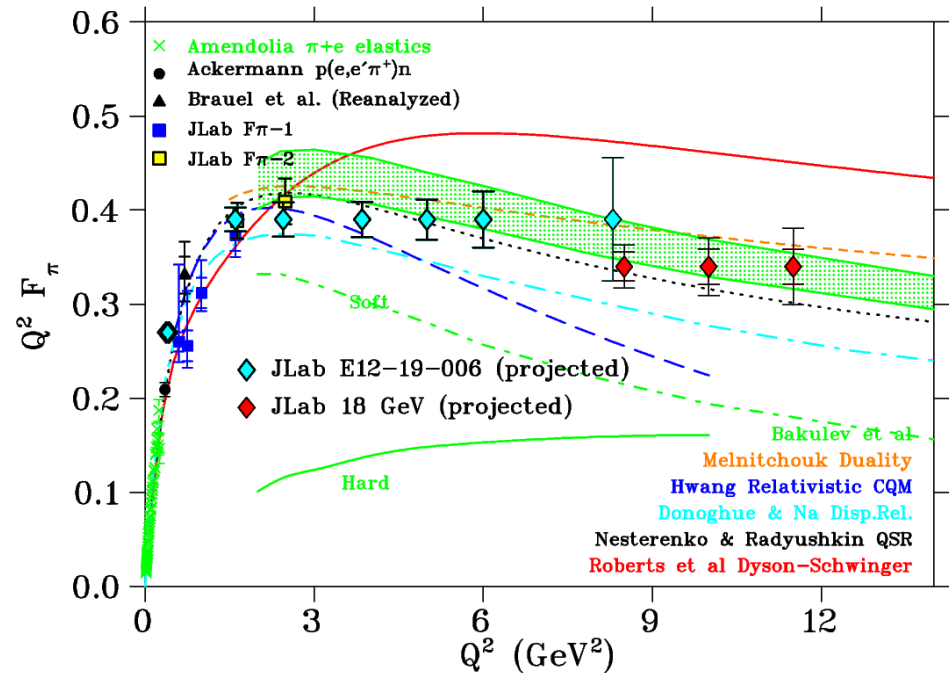


Hall C
instrumentation
has been
optimized for
specifically such
studies

Phase 1 Scenario: π^+ Form Factor

- 7.2 GeV/c HMS & 11.0 GeV/c SHMS allow a lot of kinematic flexibility, **with no major upgrades**

- Success depends on good K^+/π^+ separation in SHMS at high momenta, likely requires a modest aerogel detector upgrade
- Experiment could be done as soon as beam energy is available!
- Maximum beam energy and higher Q^2 reach constrained by sum of HMS+SHMS maximum momenta



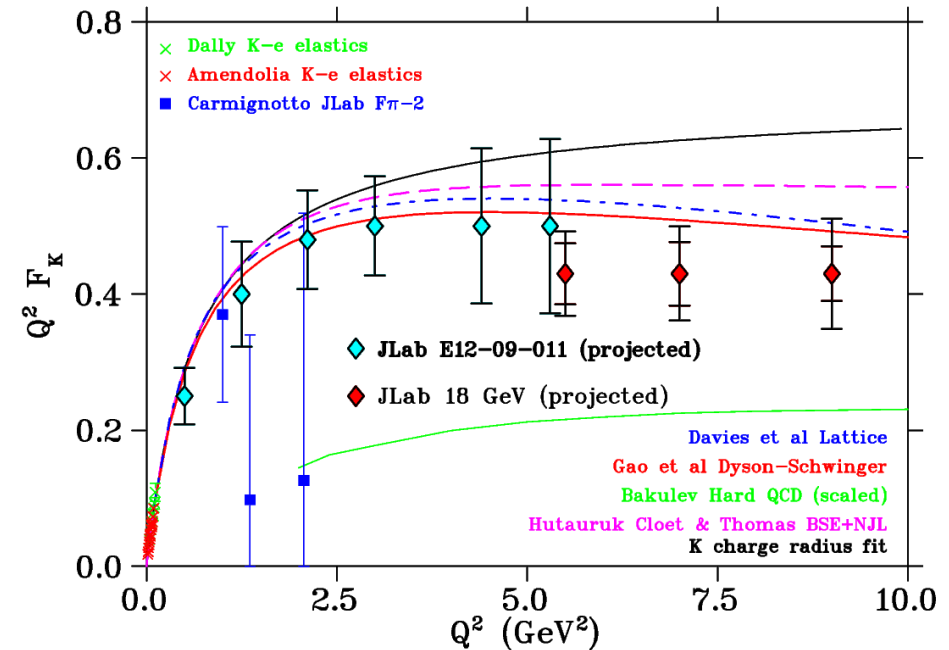
- F_π assumes same statistics as acquired in PionLT experiment
- Inner error bar is projected statistical and systematic error
- Outer error bar also includes a model uncertainty in the form factor extraction, added in quadrature

	10.6 GeV	18.0 GeV	Improvement in $\delta F_\pi / F_\pi$
$Q^2=8.5$	$\Delta\varepsilon=0.22$	$\Delta\varepsilon=0.40$	17.9% \rightarrow 4.6%
$Q^2=10.0$	New high quality F_π data		
$Q^2=11.5$	Larger F_π extraction uncertainty due to higher $-t_{\min}$		

- Since quality L–T separations are impossible at EIC (can't access $\varepsilon < 0.95$) this extension of L–T separated data considerably increases F_π data set overlap between JLab and EIC

Phase 1 Scenario: K^+ Form Factor

- 7.2 GeV/c HMS & 11.0 GeV/c SHMS allow a lot of kinematic flexibility
- Maximum beam energy and higher Q^2 reach constrained by sum of HMS+SHMS maximum momenta
- Success depends on good K^+/π^+ separation in SHMS at high momenta, likely requires a modest aerogel detector upgrade
- Counting rates are roughly 10x lower than pion form factor measurement



- Projected running times extremely long
- F_K errors uncertain, as E12-09-011 analysis not yet completed
- F_K feasibility studies at EIC are ongoing, but we already know that such measurements there are exceptionally complex.
- JLab measurements likely a complement to those at EicC.

	10.6 GeV	16.0 GeV	Improvement in $\delta F_K/F_K$
$Q^2=5.5$	$\Delta\varepsilon=0.33$	$\Delta\varepsilon=0.40$	17.9% \rightarrow 10.4%
$Q^2=7.0$	New high quality F_K data		
$Q^2=9.0$	Larger F_K extraction uncertainty due to higher $-t_{\min}$		

Phase 2 Scenario: π^+ Form Factor

■ Replace HMS with VHMS for π^+ , use SHMS for e'

- Assume $\theta_{\min}=5.5^\circ$, $\theta_{\text{open}}=15.0^\circ$
- VHMS: $\Delta\Omega$, $\Delta P/P$ similar SHMS

■ $P_{\text{VHMS}}=15.0$ GeV/c is sufficient, constrained by max beam energy

■ $\theta_{\text{VHMS}}\sim 5.5^\circ$ allows improved $\Delta\varepsilon$, but does not affect maximum Q^2 reach

■ $\theta_{\text{SHMS}}<12.0^\circ$, $P_{\text{SHMS}}>9.0$ not used

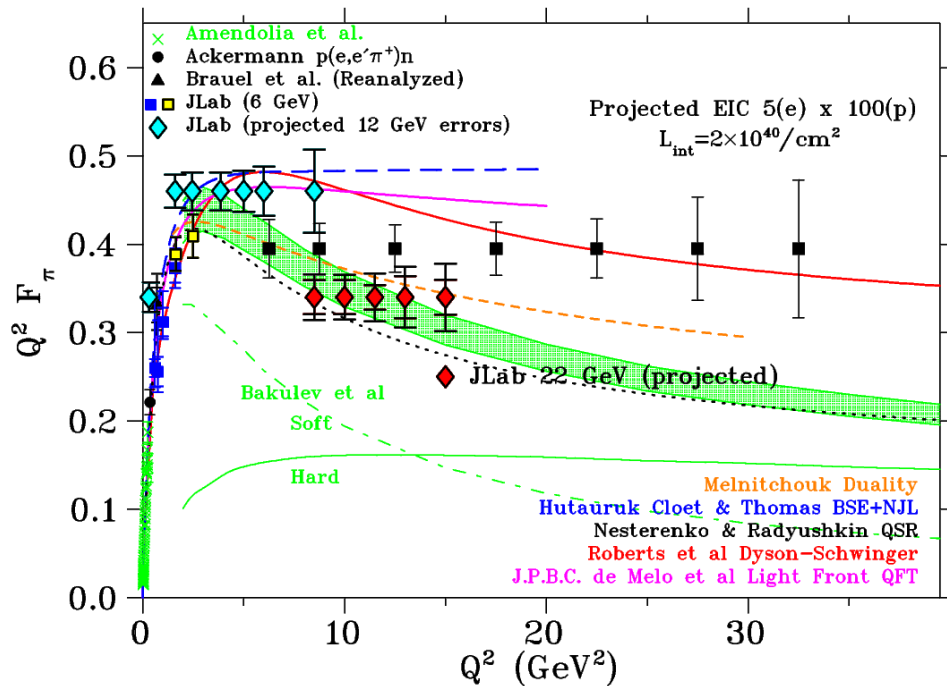
■ Dramatic increase in upper Q^2 11.5 \rightarrow 15.0 GeV²

■ Error bars for $Q^2=8.5\text{--}11.5$ GeV² substantially decrease due to smaller $-t_{\min}$ (better $R=\sigma_T/\sigma_L$) and shorter running times

■ $Q^2=15.0$ GeV² point would be “expensive” in terms of running time, but it would likely have very high scientific priority

■ Feasible scenario for Phase 2 Upgrade

$p(e,e'\pi^+)n$ Kinematics					
E_{beam}	$\theta_{\text{SHMS}}(e')$	$P_{\text{SHMS}}(e')$	$\theta_{q(\text{VHMS})}(\pi^+)$	$P_{\text{VHMS}}(\pi^+)$	Time FOM
$Q^2=8.5$ $W=4.18$ $-t_{\min}=0.15$ $\Delta\varepsilon=0.28$					
17.0	21.39	3.63	5.55	13.29	20.5
22.0	12.15	8.63	7.62	13.29	1.8
$Q^2=10.0$ $W=4.08$ $-t_{\min}=0.21$ $\Delta\varepsilon=0.30$					
17.0	24.49	3.27	5.52	13.62	53.3
22.0	13.46	8.27	7.85	13.62	4.3
$Q^2=11.5$ $W=3.95$ $-t_{\min}=0.29$ $\Delta\varepsilon=0.31$					
17.0	27.34	3.03	5.55	13.82	124.8
22.0	14.66	8.03	8.12	13.82	9.3
$Q^2=13.0$ $W=3.96$ $-t_{\min}=0.35$ $\Delta\varepsilon=0.25$					
18.0	27.55	3.18	5.54	14.63	209.5
22.0	16.49	7.18	7.69	14.63	24.4
$Q^2=15.0$ $W=3.73$ $-t_{\min}=0.52$ $\Delta\varepsilon=0.26$					
18.0	30.24	3.06	5.73	14.66	560
22.0	17.88	7.06	8.07	14.66	65.7



JLab 22 GeV			EIC 5x100		
Q^2	W	$-t_{\min}$	Q^2	W	$-t_{\min}$
8.5	4.18	0.15	5.9	7.71	0.02
10.0	4.08	0.21	8.5	8.06	0.02
11.5	3.95	0.29	11.7	8.53	0.02
13.0	3.96	0.35	16.9	8.88	0.03
15.0	3.73	0.52	22.5	9.03	0.05

Quality L/T-separations impossible at EIC (can't access $\epsilon < 0.95$)

- High W can be accessed, so $-t_{\min}$ is low
- Projected T/L ratio: ~ 0.05 at $-t_{\min}$, to 0.5 at $-t = 0.3 \text{ GeV}^2$
- Model must be used to correct for σ_T contribution
- Model must be validated from other data, adds systematic uncertainty

JLab will remain ONLY source of quality L–T separated data!

- $-t_{\min}$ is higher, but true L/T separation is performed
- Overlap of F_π data set between JLab and EIC needed to constrain EIC model uncertainty

Nucleon charge/magnetic FF ratios

- Proton electric form factor possesses a zero
 - $Q^2 = 8.86^{+1.93}_{-0.86} \text{ GeV}^2$
- Neutron electric form factor is positive definite
 - $G_E^n(Q^2) > G_E^p(Q^2)$ on $Q^2 > 4.7 \text{ GeV}^2$
 - On this domain, electric form factor of charge-neutral neutron is larger than that of charge-one proton
 - A remarkable, non-intuitive result!
- Verification of this is within JLab reach
 - perhaps already in Gen-II data?
- Curves are Continuum Schwinger Model (CSM)

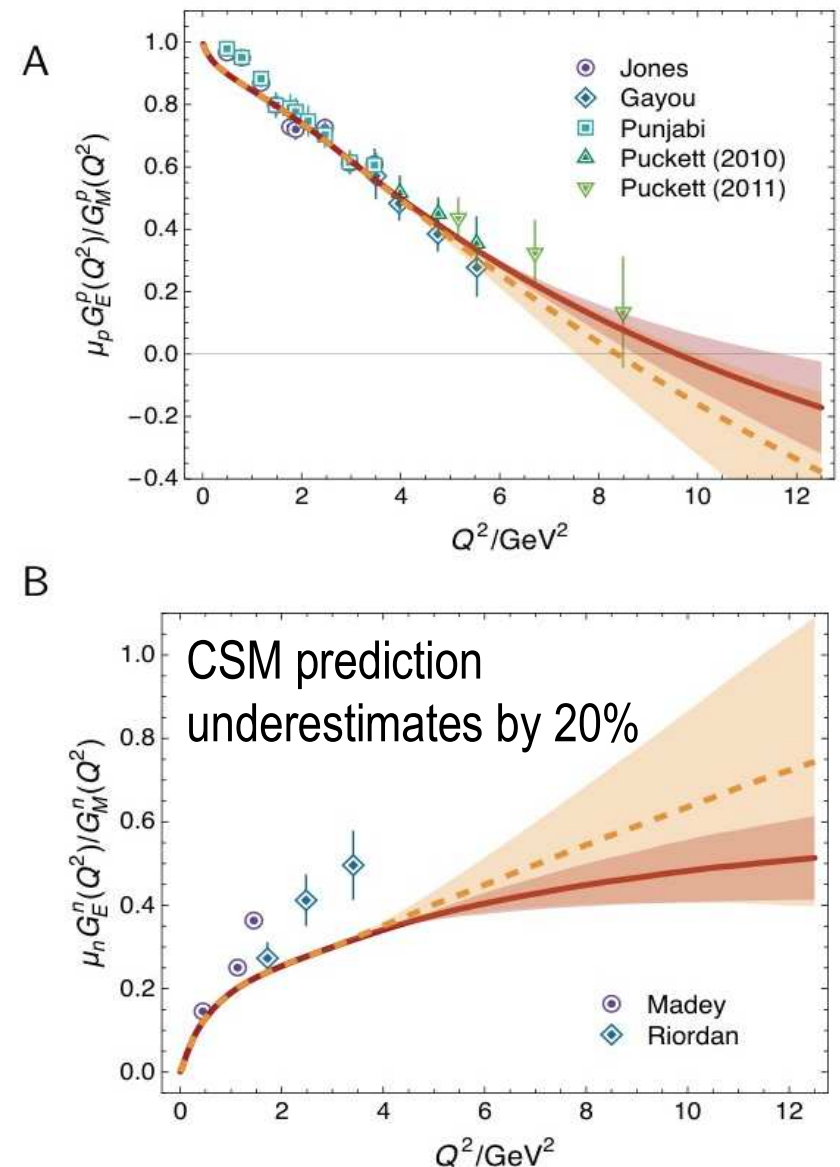


FIG. 6. Panel A: $\mu_p G_E^p/G_M^p$. Panel B: $\mu_n G_E^n/G_M^n$. SPM I – dashed orange curve within like-coloured band; and SPM II – solid red curve within like-coloured band. Data: proton – Refs. [20–24]; and neutron – Refs. [87, 97].

■ Isospin symmetry limit:

- Behaviors of $\mu_p G_E^p/G_M^p$ and $\mu_p G_E^n/G_M^n$ are correlated ($e_u=2/3$, $e_d=-1/3$)

- $G_E^p = e_u G_E^u + e_d G_E^d$

- $G_E^n = e_u G_E^d + e_d G_E^u$

■ G_E^p possesses a zero because

- although remaining positive, G_E^u/G_M^p falls steadily with increasing Q^2
- while $G_E^d/G_M^p > 0$ and approximately constant

■ G_E^n predicted to NOT exhibit a zero at high Q^2 because

- $e_u > 0$, G_E^d/G_M^p large & positive
- $|e_d G_E^u|$ always $< e_u G_E^d$

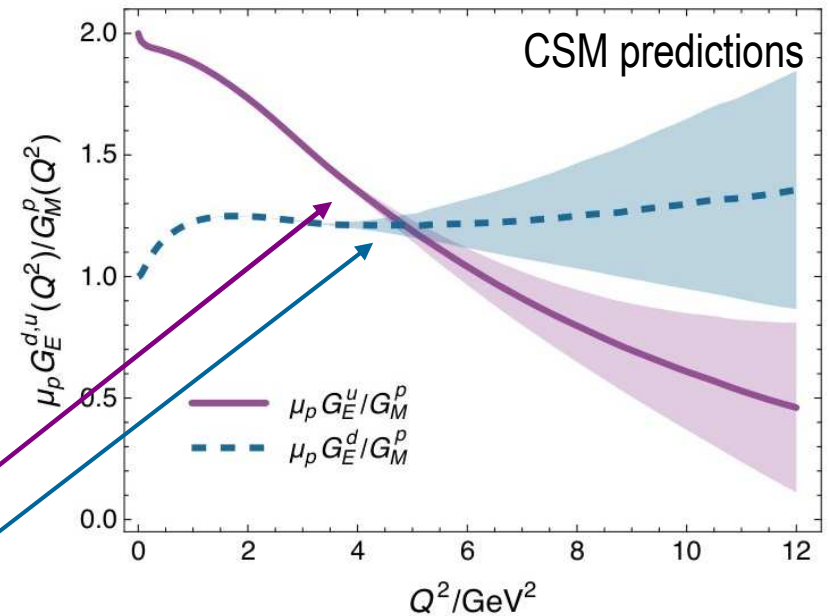
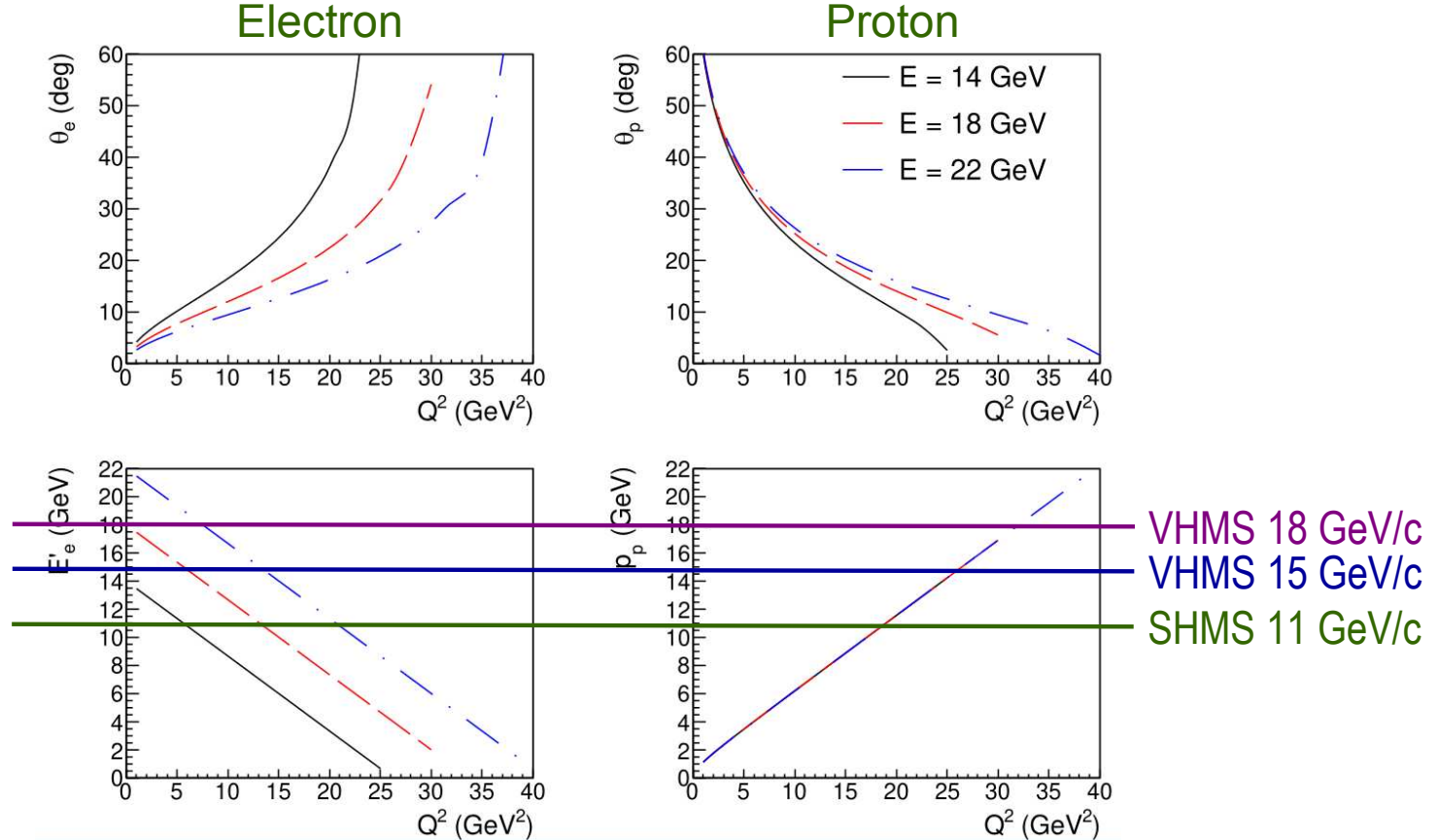


Figure 8: Flavour separation of the charge and magnetisation form factors, with each function normalised by G_M^p in order to highlight their differing Q^2 -dependence.

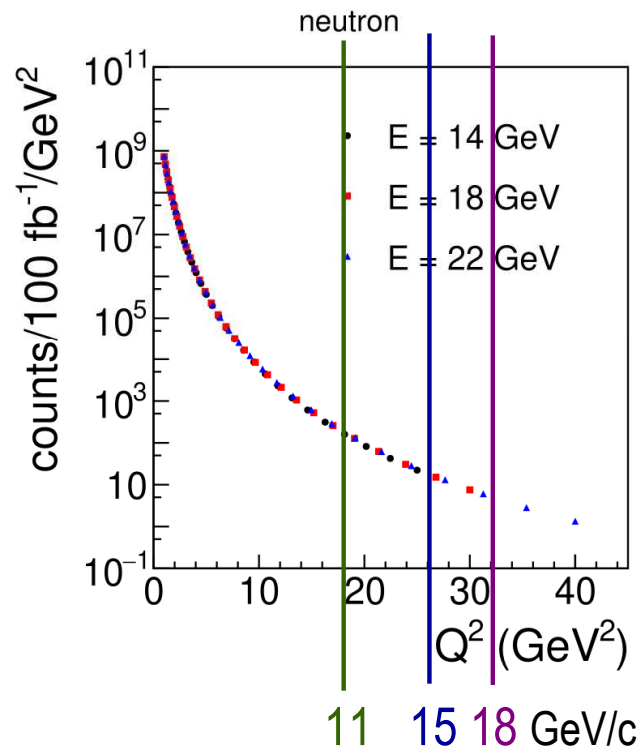
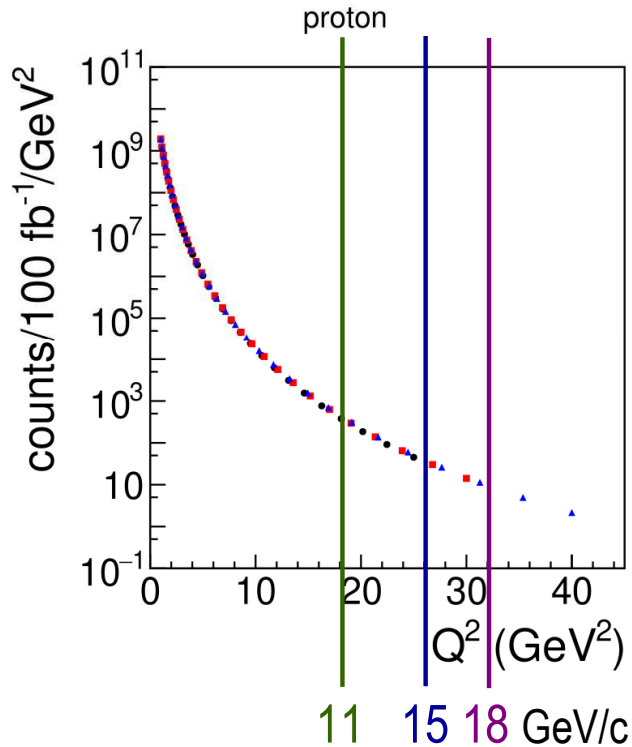
- For nucleons, discovering QCD scaling and scaling violations requires $Q^2 > 18 \text{ GeV}^2$
- Each feature is a sensitive expression of emergent phenomena in QCD

Elastic eN Kinematics @ 14, 18, 22 GeV



- **Scattering angles well matched** to acceptances of a variety of JLab detectors (CLAS12, SoLID, HMS+SHMS, SBS+BigBite)
- **High proton momentum is a challenge:**
 - SHMS could measure up to $Q^2=18$ GeV²
 - VHMS could allow $Q^2=26-32$ GeV², depending on capability
 - Large acceptance spectrometers would have resolution challenges at high Q^2

Elastic eN Count Rate Projections



	Luminosity (cm ⁻² s ⁻¹ , e-N)	fb ⁻¹ / day
NH ₃ /ND ₃ polarized p/n	10 ³⁵	8.64
³ He SEOP polarized n	10 ³⁷	8.64
LH ₂ /LD ₂ unpolarized	10 ³⁸⁻³⁹	8600 – 86000

- **Event rate per unit integrated luminosity for 14, 18, 22 GeV beam**
 - Assumes Q² bin width ≈ Q² spacing between points
 - Assumes 2π azimuthal acceptance
- **JLab provides lots of count rate up to maximum accessible Q²**
 - EIC best case: 100 fb⁻¹/year

- **Meson and nucleon form factors are fundamental structure observables that are intimately linked to many open questions in QCD**
 - QCD's transition from soft to hard degrees of freedom
 - Unraveling emergent phenomena in QCD
- **22 GeV upgrade enables a significant expansion of Q^2 reach**
 - π^+ form factor up to $Q^2=11.5(15)$ GeV^2 with SHMS(VHMS)
 - K^+ form factor up to $Q^2=9$ GeV^2 with SHMS
 - Unfortunately, running times are very long
 - Proton/Neutron up to $Q^2=18(26-32)$ GeV^2
- **Understanding how QCD explains the emergence of hadron mass and structure requires investment in facilities that can deliver precision data on mesons, nucleons, and beyond**

F_π PionLT Projected Uncertainties

$p(e,e'\pi^+)n$ Kinematics										
E_{beam} (GeV)	ϵ	θ_{HMS} (e')	P_{HMS} (e')	$\theta_{q(\text{SHMS})}$ (π^+)	P_{SHMS} (π^+)	LH ₂ Run hrs	Online #Events /t-bin	$\delta\sigma_{\text{UNS}}$ (stat & est uncorrel syst unc)	$\delta\sigma_{\text{L}}$ (stat & est uncorrel+ correlated syst unc)	δF_π (stat & est syst unc / incl est model unc)
$Q^2=5.0$ $W=2.95$ $-t_{\text{min}}=0.21$ $R(\text{VR})=\text{T/L}=0.79$										
8.0	0.22	44.37	1.10	6.17	6.72	64	3.9k	2.3%	10.8%	5.4% / 12.2%
10.6	0.60	20.57	3.72	10.47	6.72	36	6.4k	2.1%		
$Q^2=6.0$ $W=3.19$ $-t_{\text{min}}=0.37$ $R(\text{VR})=\text{T/L}=0.70$										
9.2	0.18	47.04	1.03	5.06	8.04	166	4.1k	2.3%	10.7%	5.4% / 12.2%
10.6	0.40	28.18	2.40	7.65	8.04	97	6.4k	2.1%		
$Q^2=8.5$ $W=2.79$ $-t_{\text{min}}=0.55$ $R(\text{VR})=\text{T/L}=1.71$										
9.2	0.15	58.53	0.97	5.44	7.91	615	2.5k	2.6%	35.8%	17.9% / 33.5%
10.6	0.38	34.11	2.34	8.67	7.91	101	1.9k	2.4%		

- σ_{UNS} errors include 0.6% pt-pt and 1.6% t-correlated syst unc from E12-06-101 proposal
- σ_{L} errors include 3.3% scale systematic uncertainty from E12-06-101 proposal

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 - Experiment could be done as soon as beam energy is available!
 - Maximum beam energy and higher Q^2 reach constrained by sum of HMS+SHMS maximum momenta
 - $Q^2=8.5$ and 11.5 Time FOM similar to PionLT $Q^2=6.0$ and 8.5 points

p(e,e' π^+)n Kinematics					
E_{beam}	$\theta_{\text{HMS}} (e')$	$P_{\text{HMS}} (e')$	$\theta_{\text{q(SHMS)}} (\pi^+)$	$P_{\text{SHMS}} (\pi^+)$	Time FOM
$Q^2=8.5$ $W=3.64$ $-t_{\text{min}}=0.24$ $\Delta\varepsilon=0.40$					
13.0	34.30	1.88	5.29	10.99	64.7
18.0	15.05	6.88	8.94	10.99	2.2
$Q^2=10.0$ $W=3.44$ $-t_{\text{min}}=0.37$ $\Delta\varepsilon=0.40$					
13.0	37.78	1.83	5.56	10.97	122.7
18.0	16.39	6.83	9.57	10.97	4.5
$Q^2=11.5$ $W=3.24$ $-t_{\text{min}}=0.54$ $\Delta\varepsilon=0.29$					
14.0	31.73	2.75	7.06	10.96	82.4
18.0	17.70	6.75	10.05	10.96	8.8

- **Since quality L–T separations are impossible at EIC (can't access $\varepsilon < 0.95$) this extension of L–T separated data considerably increases F_π data set overlap between JLab and EIC**

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$p(e,e'\pi^+)n$ Kinematics										
E_{beam} (GeV)	ϵ	θ_{HMS} (e')	P_{HMS} (e')	$\theta_{q(\text{SHMS})}$ (π^+)	P_{SHMS} (π^+)	Time FOM	#Events /t-bin	$\delta\sigma_{\text{UNS}}$ (stat & est uncorrel syst unc)	$\delta\sigma_{\text{L}}$ (stat & est uncorrel+ correlated syst unc)	δF_π (stat & est syst unc / incl est model unc)
$Q^2=8.5$ $W=3.64$ $-t_{\text{min}}=0.24$ $R(\text{VR})=\text{T/L}=0.56$										
13.0	0.25	34.30	1.88	5.29	10.99	64.7	3.2k	2.5%	9.1%	4.6% / 9.6%
18.0	0.65	15.05	6.88	8.94	10.99	2.2	4.8k	2.2%		
$Q^2=10.0$ $W=3.44$ $-t_{\text{min}}=0.37$ $R(\text{VR})=\text{T/L}=0.80$										
13.0	0.24	37.78	1.83	5.56	10.97	122.7	2.8k	2.56%	11.1%	5.6% / 12.7%
18.0	0.64	16.39	6.83	9.57	10.97	4.5	4.2k	2.3%		
$Q^2=11.5$ $W=3.24$ $-t_{\text{min}}=0.54$ $R(\text{VR})=\text{T/L}=1.20$										
14.0	0.22	31.73	2.75	7.06	10.96	82.4	2.0k	2.8%	16.3%	5.4% / 18.8%
18.0	0.63	17.70	6.75	10.05	10.96	8.8	2.0k	2.8%		

- σ_{UNS} errors include 0.6% pt-pt and 1.6% t-correlated syst unc from E12-06-101 proposal
- σ_{L} errors include 3.3% scale systematic uncertainty from E12-06-101 proposal

Phase 1 Scenario: K^+ Form Factor

- 7.2 GeV/c HMS & 11.0 GeV/c SHMS allow a lot of kinematic flexibility
- Maximum beam energy and higher Q^2 reach constrained by sum of HMS+SHMS maximum momenta
- Success depends on good K^+/π^+ separation in SHMS at high momenta, likely requires a modest aerogel detector upgrade
- Counting rates are roughly 10x lower than pion form factor measurement

	10.6 GeV	16.0 GeV	Improvement in $\delta F_K/F_K$
$Q^2=5.5$	$\Delta\varepsilon=0.33$	$\Delta\varepsilon=0.40$	17.9%→10.4%
$Q^2=7.0$	New high quality F_K data		
$Q^2=9.0$	Larger F_K extraction uncertainty due to higher $-t_{min}$		

p(e,e' K^+) Λ Kinematics					
E_{beam}	θ_{HMS} (e')	P_{HMS} (e')	$\theta_{q(SHMS)}$ (K^+)	P_{SHMS} (K^+)	Time FOM
$Q^2=5.5$ $W=3.56$ $-t_{min}=0.32$ $\Delta\varepsilon=0.40$					
11.0	30.69	1.79	5.50	8.84	746
16.0	12.92	6.79	9.18	8.84	150
$Q^2=7.0$ $W=3.90$ $-t_{min}=0.33$ $\Delta\varepsilon=0.29$					
14.0	25.16	2.64	5.51	10.98	620
18.0	13.91	6.64	7.85	10.98	192
$Q^2=9.0$ $W=3.66$ $-t_{min}=0.54$ $\Delta\varepsilon=0.30$					
14.0	29.17	2.54	5.98	10.97	964
18.0	15.90	6.54	8.69	10.97	350

- F_K feasibility studies at EIC are ongoing, but we already know that such measurements there are exceptionally complex.
- JLab measurements likely a complement to those at EicC.

p(e,e'K⁺)Λ Kinematics										
E_{beam} (GeV)	ϵ	θ_{HMS} (e')	P_{HMS} (e')	$\theta_{\text{q(SHMS)}}$ (K ⁺)	P_{SHMS} (K ⁺)	Time FOM	#Events /t-bin	$\delta\sigma_{\text{UNS}}$ (stat & est uncorrel syst unc)	$\delta\sigma_{\text{L}}$ (stat & est uncorrel+ correlated syst unc)	δF_{π} (stat & est syst unc / incl est model unc)
$Q^2=5.5$ $W=3.56$ $-t_{\text{min}}=0.22$ $R(\text{VR})=\text{T/L}=1.17$										
11.0	0.36	30.69	1.79	5.50	8.84	746	3.0k	2.5%	20.8%	10.4% / 20.4%
16.0	0.64	12.92	6.79	9.18	8.84	150	3.0k	2.6%		
$Q^2=7.0$ $W=3.90$ $-t_{\text{min}}=0.33$ $R(\text{VR})=\text{T/L}=1.19$										
14.0	0.34	25.16	2.64	5.51	10.98	620	2.5k	2.6%	21.6%	10.8% / 21.0%
18.0	0.63	13.91	6.64	7.85	10.98	192	2.5k	2.6%		
$Q^2=9.0$ $W=3.66$ $-t_{\text{min}}=0.54$ $R(\text{VR})=\text{T/L}=1.63$										
14.0	0.32	29.17	2.54	5.98	10.97	964	2.0k	2.8%	28.1%	9.4% / 28.0%
18.0	0.62	15.60	6.54	8.69	10.97	350	2.0k	2.8%		

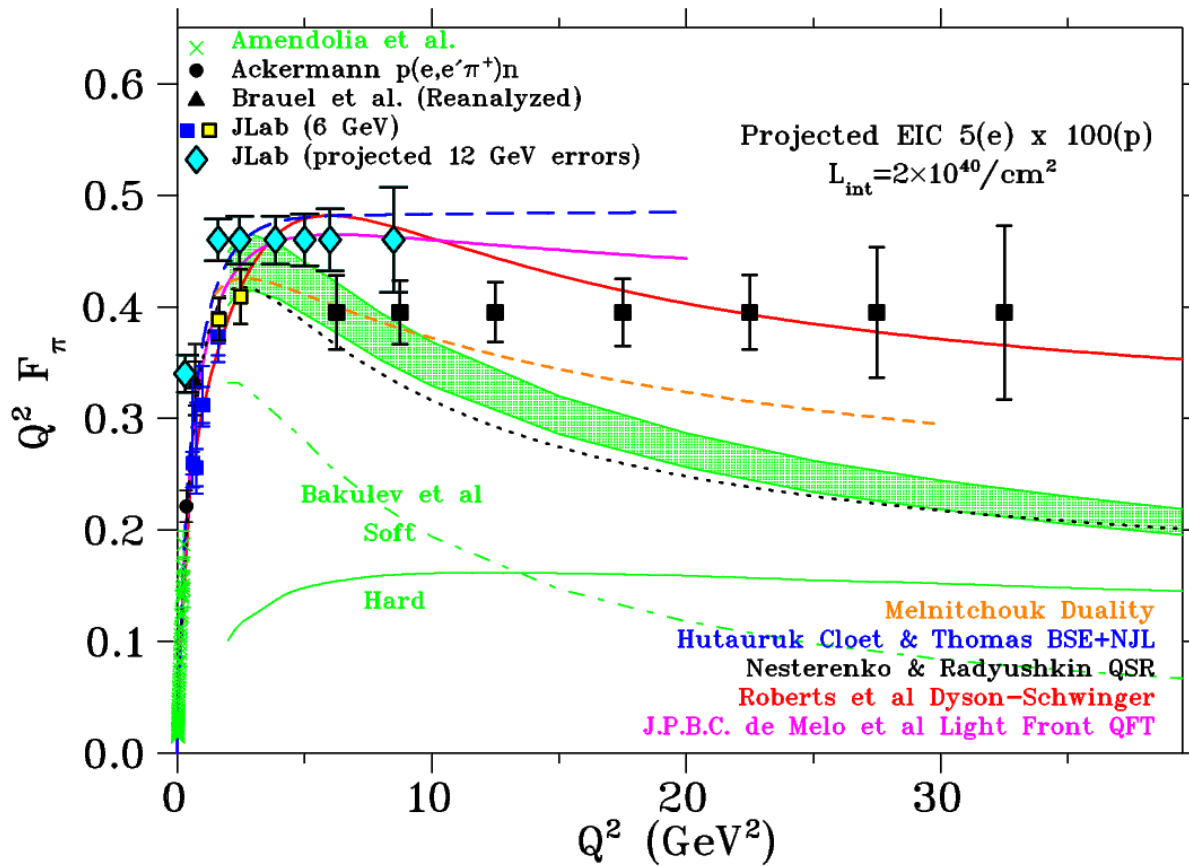
- σ_{UNS} errors include 0.6% pt-pt and 1.6% t-correlated syst unc from E12-06-101 proposal
- σ_{L} errors include 3.3% scale systematic uncertainty from E12-06-101 proposal

F_π Hall C Phase 2 (new VHMS for π^+ , SHMS for e')



E_{beam} (GeV)	ϵ	θ_{SHMS} (e')	P_{SHMS} (e')	$\theta_{q(\text{VHMS})}$ (π^+)	P_{VHMS} (π^+)	Time FOM	#Events /t-bin	$\delta\sigma_{\text{UNS}}$ (stat & est uncorrel syst unc)	$\delta\sigma_{\text{L}}$ (stat & est uncorrel+ correlated syst unc)	δF_π (stat & est syst unc / incl est model unc)
$Q^2=8.5$ $W=4.18$ $-t_{\text{min}}=0.15$ $R(\text{VR})=T/L=0.37$										
17.0	0.39	21.39	3.63	5.55	13.29	20.5	3.2k	2.5%	11.2%	5.6% / 10.6%
22.0	0.67	12.15	8.63	7.62	13.29	1.8	4.8k	2.2%		
$Q^2=10.0$ $W=4.08$ $-t_{\text{min}}=0.21$ $R(\text{VR})=T/L=0.45$										
17.0	0.21	24.49	3.27	5.56	13.62	53.3	3.2k	2.5%	11.2%	5.6% / 10.6%
22.0	0.64	13.46	8.27	9.57	13.62	4.3	4.8k	2.2%		
$Q^2=11.5$ $W=3.95$ $-t_{\text{min}}=0.29$ $R(\text{VR})=T/L=0.50$										
17.0	0.32	27.34	3.03	5.55	13.82	124.8	2.8k	2.6%	11.3%	5.6% / 12.4%
22.0	0.63	14.66	8.03	8.12	13.82	9.3	4.2k	2.3%		
$Q^2=13.0$ $W=3.96$ $-t_{\text{min}}=0.35$ $R(\text{VR})=T/L=0.80$										
18.0	0.32	27.55	3.18	5.54	14.63	209.5	2.4k	2.7%	14.2%	7.0% / 14.3%
22.0	0.57	16.49	7.18	7.69	14.63	24.4	3.6k	2.4%		
$Q^2=15.0$ $W=3.73$ $-t_{\text{min}}=0.52$ $R(\text{VR})=T/L=1.20$										
18.0	0.30	30.24	3.06	5.73	14.66	560	2.0k	2.8%	17.9%	6.0% / 18.6%
22.0	0.56	17.88	7.06	8.07	14.66	65.7	2.0k	2.8%		

EIC Kinematic Reach (projection)



Assumptions:

- $5(e^-) \times 100(p)$.
- Integrated $L=20 \text{ fb}^{-1}/\text{yr}$.
- Clean identification of exclusive $p(e, e'\pi^+n)$ events.
- Syst. Unc: 2.5% pt-pt and 12% scale.
- $R=\sigma_L/\sigma_T=0.013-0.14$ at lowest $-t$ from VR model, and $\delta R=R$ syst. unc. in model subtraction to isolate σ_L .
- π pole dominance at small $-t$ confirmed in ${}^2\text{H } \pi^-/\pi^+$ ratios.

ECCE 2022 projections shown

Projections to be updated soon using latest ePIC detector simulation

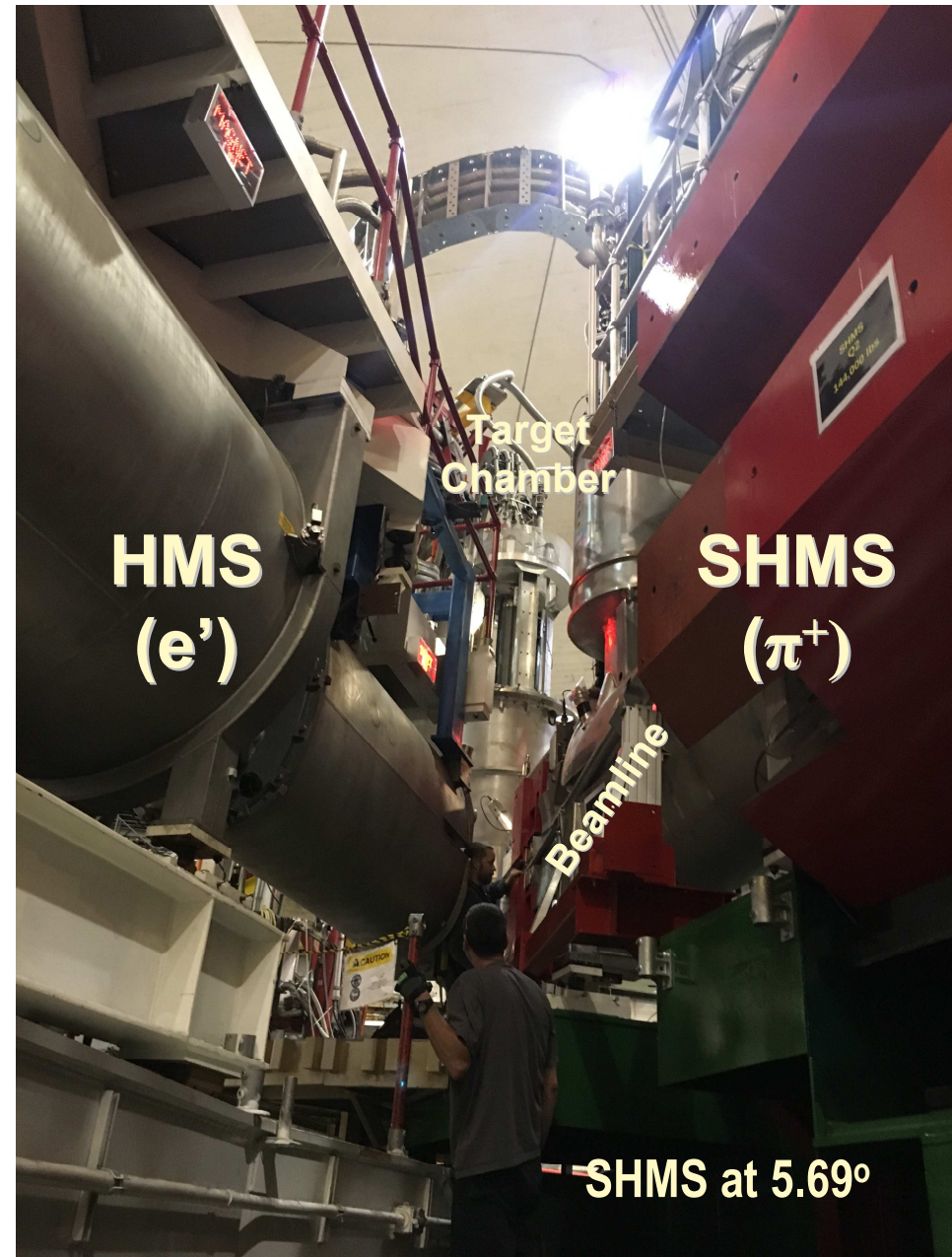
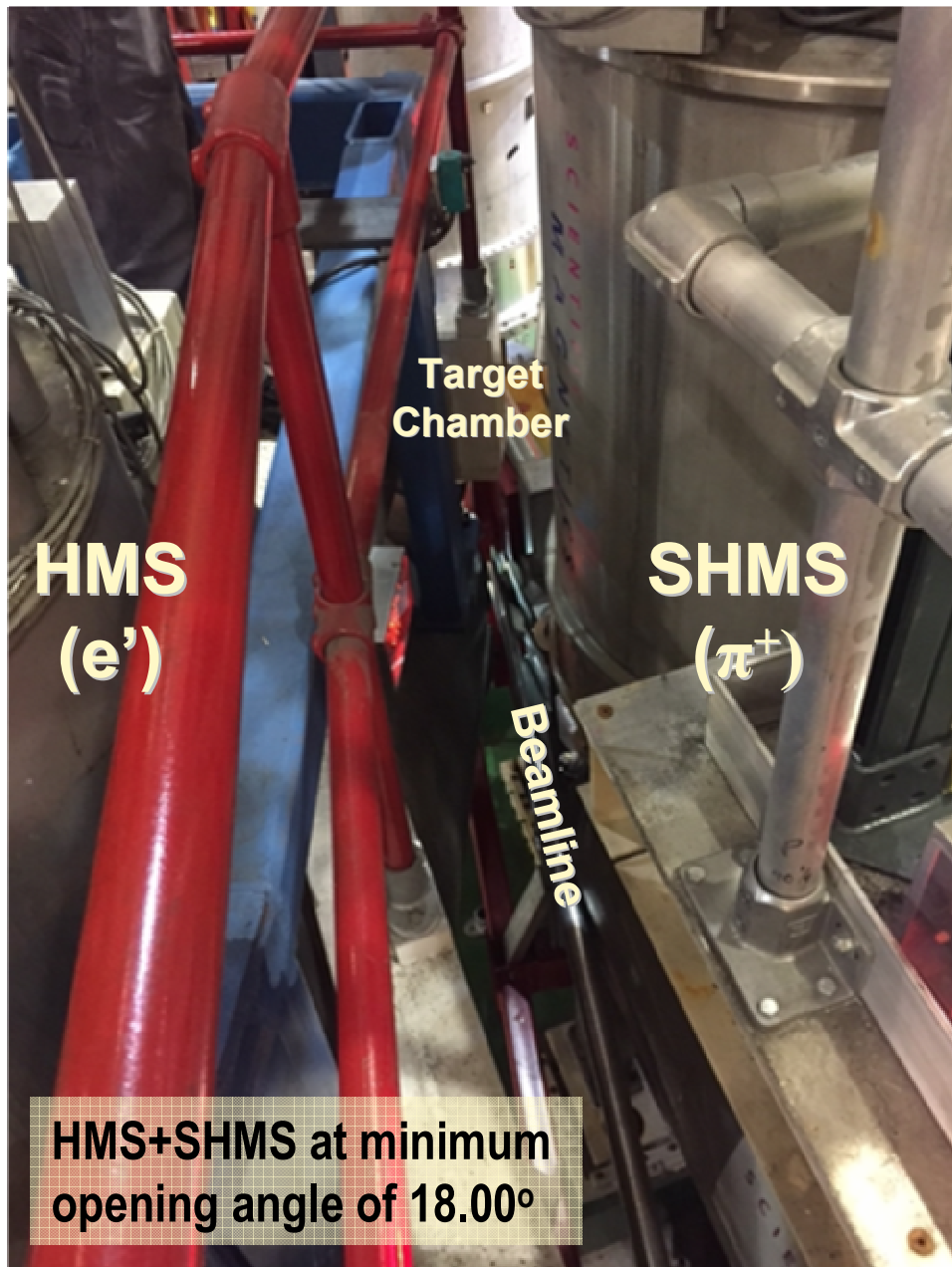
F_π EIC 5x100 $\epsilon=0.999$

Q^2 mean (GeV ²)	W_{mean} (GeV)	$-t_{\text{bin}}$ (GeV ²)	$R(\text{VR})$ =T/L	Rate (Hz) $L=10^{34}\text{cm}^{-2}\text{s}^{-1}$	#Events /t-bin $L_{\text{int}}=2\times 10^{40}\text{cm}^{-2}$	$\delta\sigma_L$ (stat & est uncorrel syst / est correlated syst & σ_T subtraction unc)	δF_π (stat and est correlated+ uncorrel syst unc)
5.87	7.71	0.02	0.12	0.29	290k	3.4% / 15.8%	8.4%
		0.30	1.03	0.12	119k	27.2% / 52.2%	
8.48	8.06	0.02	0.073	0.045	45.1k	3.4% / 13.8%	7.4%
		0.30	0.65	0.019	18.9k	17.8% / 41.1%	
11.71	8.53	0.02	0.047	0.012	12.4k	4.4% / 12.8%	7.1%
		0.30	0.42	8.6E-3	8.6k	12.4% / 32.1%	
16.88	8.88	0.06	0.038	3.2E-4	3.5k	7.1% / 12.5%	8.0%
		0.30	0.27	1.2E-3	1.2k	13.8% / 24.4%	
22.02	9.03	0.06	0.027	5.5E-4	550	16.9% / 12.3%	12.6%
		0.30	0.20	3.4E-4	340	21.0% / 20.3%	
27.03	9.09	0.10	0.035	2.0E-4	200	28.3% / 12.5%	20.0%
		0.30	0.16	1.0E-4	100	36.4% / 18.1%	

- Assume 2.5% pt-pt syst error and 12% scale syst error (similar to HERA-H1 pion struct fcn)
- Assume uncertainty in $R=T/L$ (due to lack of L/T sep) is bounded by R , i.e. $\delta R/R=1$

Hall C during Data Taking

π^+/K^+ FF experiments have challenging forward angle requirements

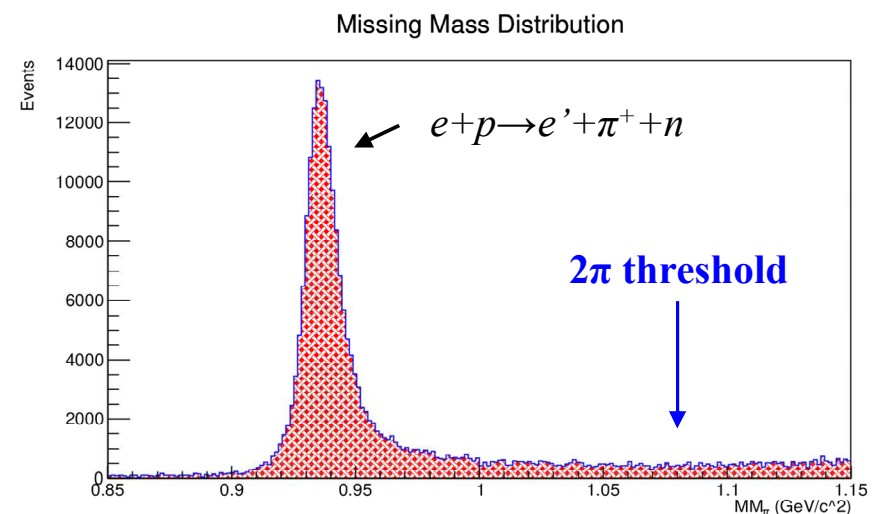
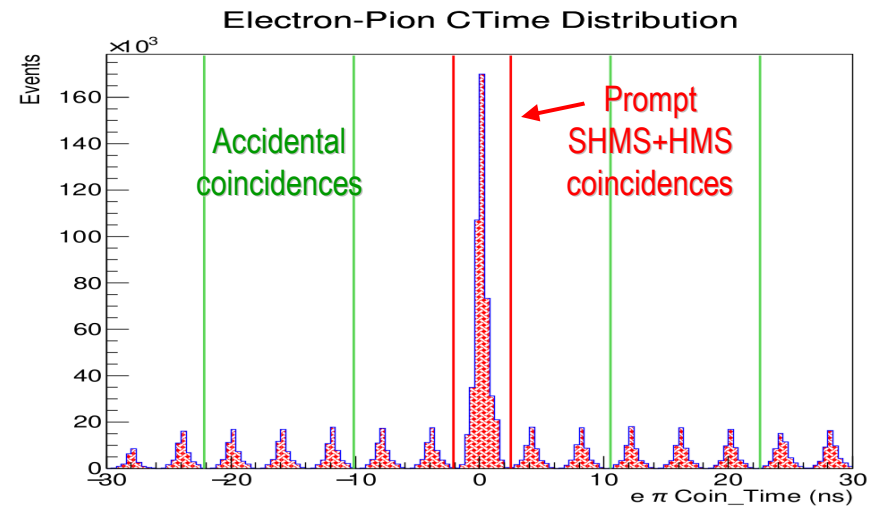
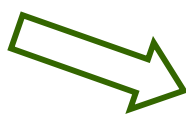
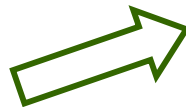


$p(e, e' \pi^+) n$ Event Selection

Coincidence measurement between charged pions in SHMS and electrons in HMS

Easy to isolate
exclusive channel

- Excellent particle identification
- CW beam minimizes “accidental” coincidences
- Missing mass resolution easily excludes 2-pion contributions



PionLT experiment E12-19-006 Data

$Q^2=1.60$, $W=3.08$, $x=0.157$, $\varepsilon=0.685$

$E_{\text{beam}}=9.177$ GeV, $P_{\text{SHMS}}=+5.422$ GeV/c, $\theta_{\text{SHMS}}=10.26^\circ$ (left)

Plots by Muhammad Junaid

Extract $F_\pi(Q^2)$ from JLab σ_L data

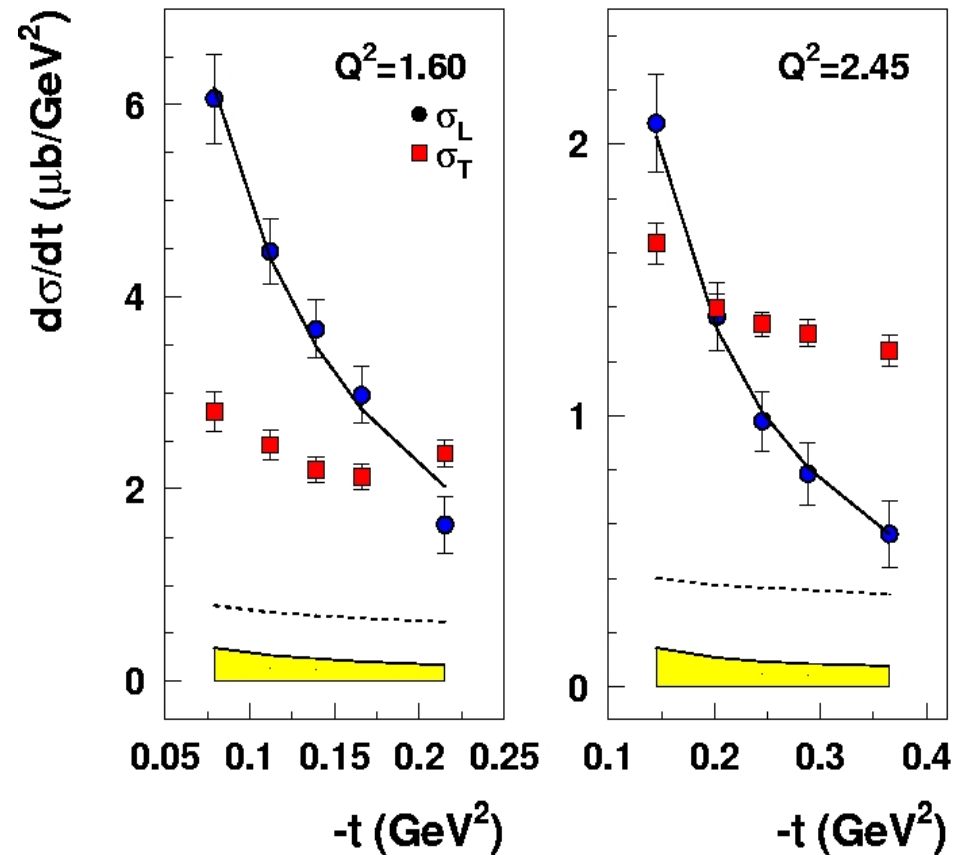
Model incorporates π^+ production mechanism and spectator neutron effects:

VGL Regge Model:

- Feynman propagator $\left(\frac{1}{t - m_\pi^2} \right)$
replaced by π and ρ Regge propagators.
 - Represents the exchange of a series of particles, compared to a single particle.
- Free parameters: $\Lambda_\pi, \Lambda_\rho$ (trajectory cutoff)
[Vanderhaeghen, Guidal, Laget, PRC 57(1998)1454]
- At small $-t$, σ_L only sensitive to F_π

$$F_\pi = \frac{1}{1 + Q^2 / \Lambda_\pi^2}$$

Fit to σ_L to model
gives F_π at each Q^2



Error bars indicate statistical and random (pt-pt) systematic uncertainties in quadrature.

Yellow band indicates the correlated (scale) and partly correlated (t-corr) systematic uncertainties.

$$\Lambda_\pi^2 = 0.513, 0.491 \text{ GeV}^2, \Lambda_\rho^2 = 1.7 \text{ GeV}^2.$$

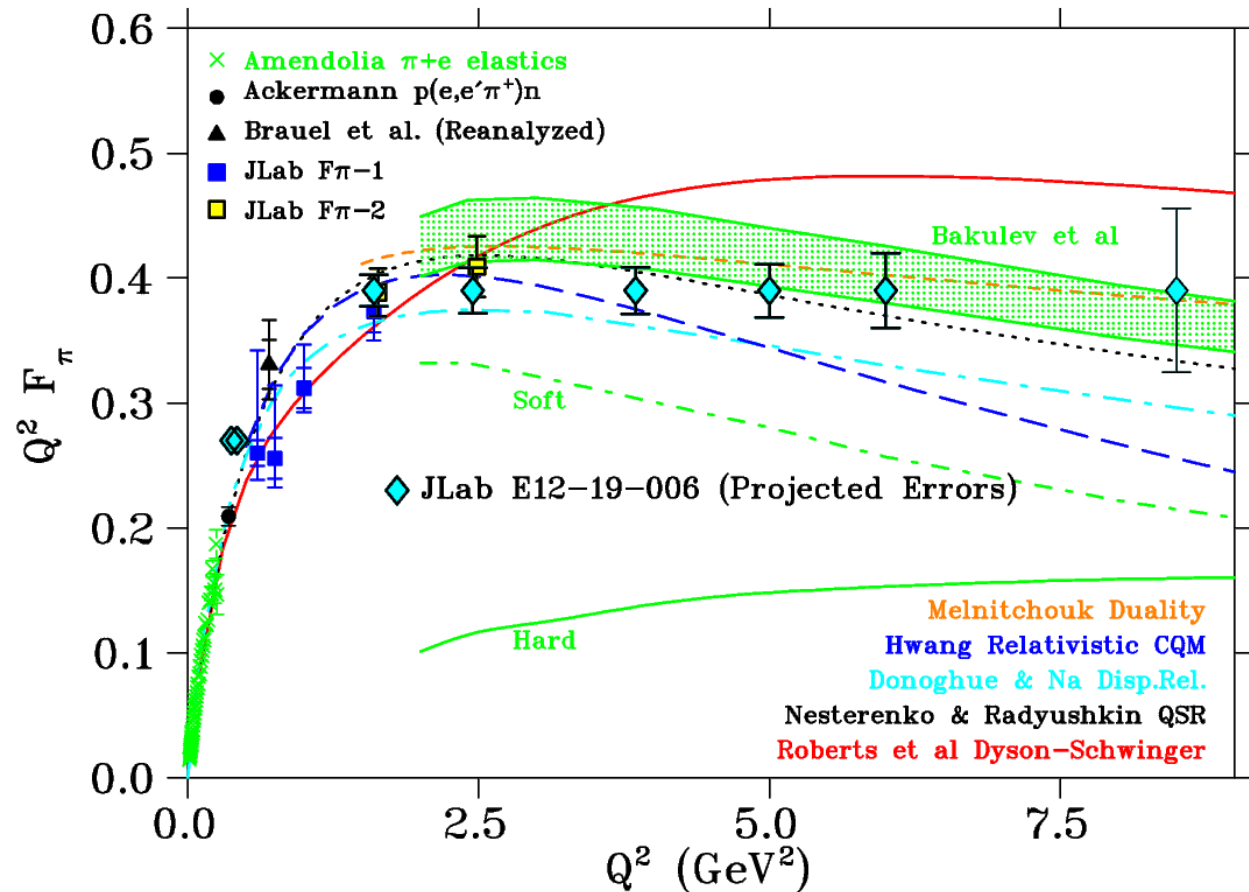
F π -2 data: T. Horn et al., PRL 97(2006)192001.

Current and Projected F_π Data

SHMS+HMS will allow measurement of F_π to much higher Q^2 .

No other facility worldwide can perform this measurement.

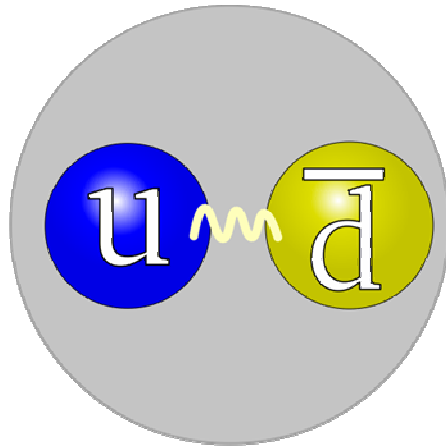
The pion form factor is the clearest test case for studies of QCD's transition from non-perturbative to perturbative regions.



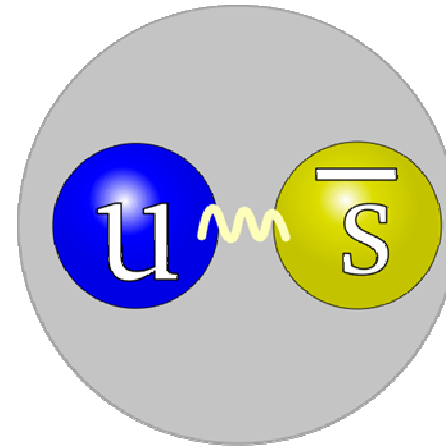
The $\sim 17\%$ measurement of F_π at $Q^2=8.5 \text{ GeV}^2$ is at higher $-t_{min}=0.45 \text{ GeV}^2$

PionLT E12-19-006: D. Gaskell, T. Horn and G. Huber, spokespersons00

The Charged Kaon – a 2nd QCD test case



π^+



K^+

- In the hard scattering limit, pQCD predicts that the π^+ and K^+ form factors will behave similarly

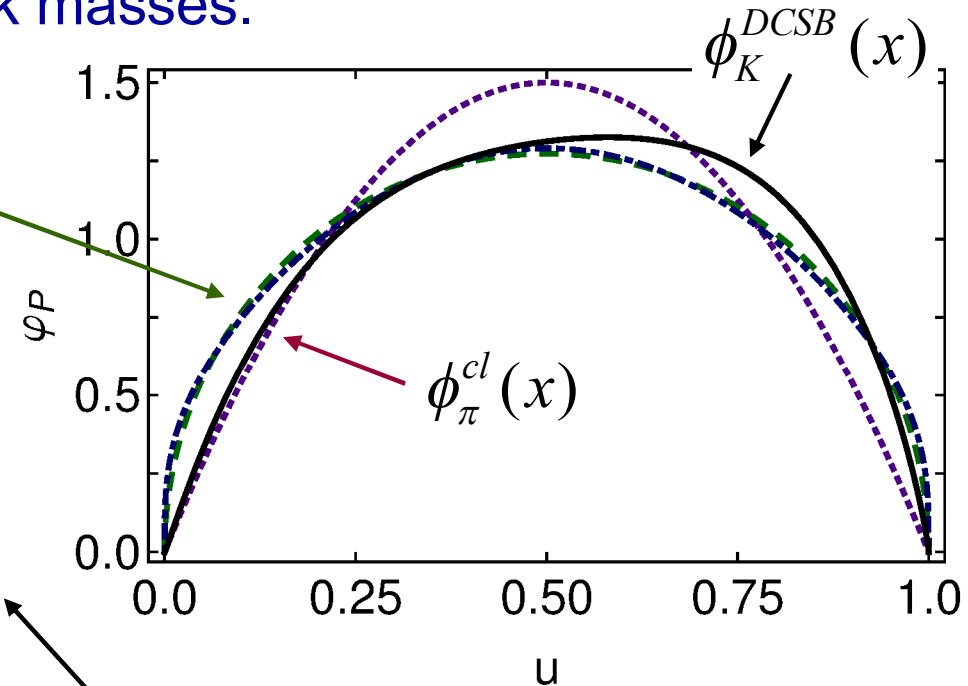
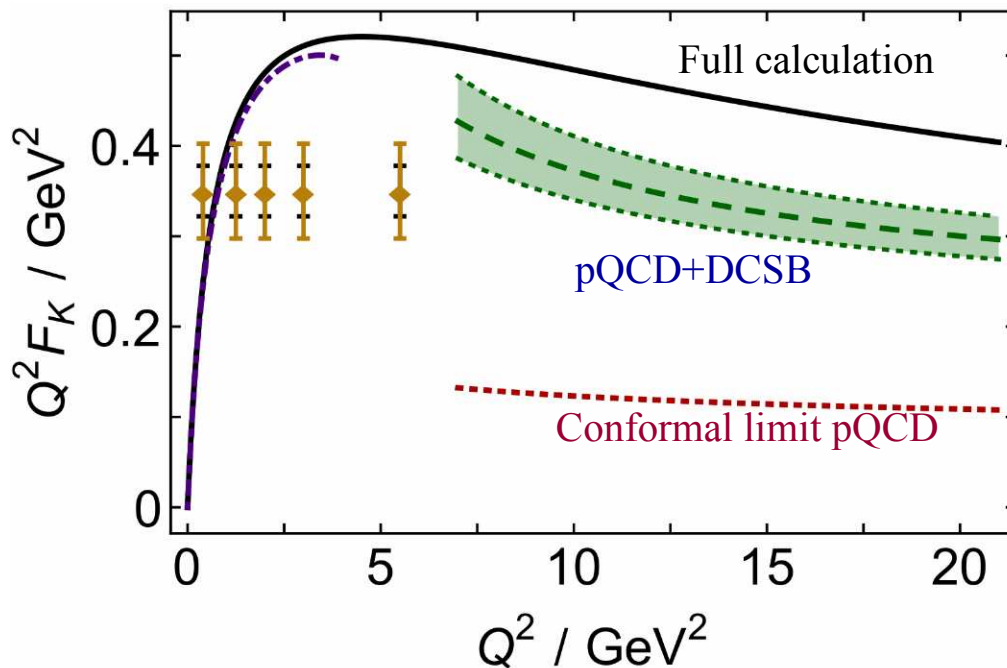
$$\frac{F_K(Q^2)}{F_\pi(Q^2)} \xrightarrow{Q^2 \rightarrow \infty} \frac{f_K^2}{f_\pi^2}$$

- It is important to compare the magnitudes and Q^2 –dependences of both form factors.

K^+ properties also strongly influenced by EHM

- K^+ PDA also is broad, concave and asymmetric.
- While the heavier s quark carries more bound state momentum than the u quark, the shift is markedly less than one might naively expect based on the difference of u, s current quark masses.

[C. Shi, et al., PRD 92 (2015) 014035].

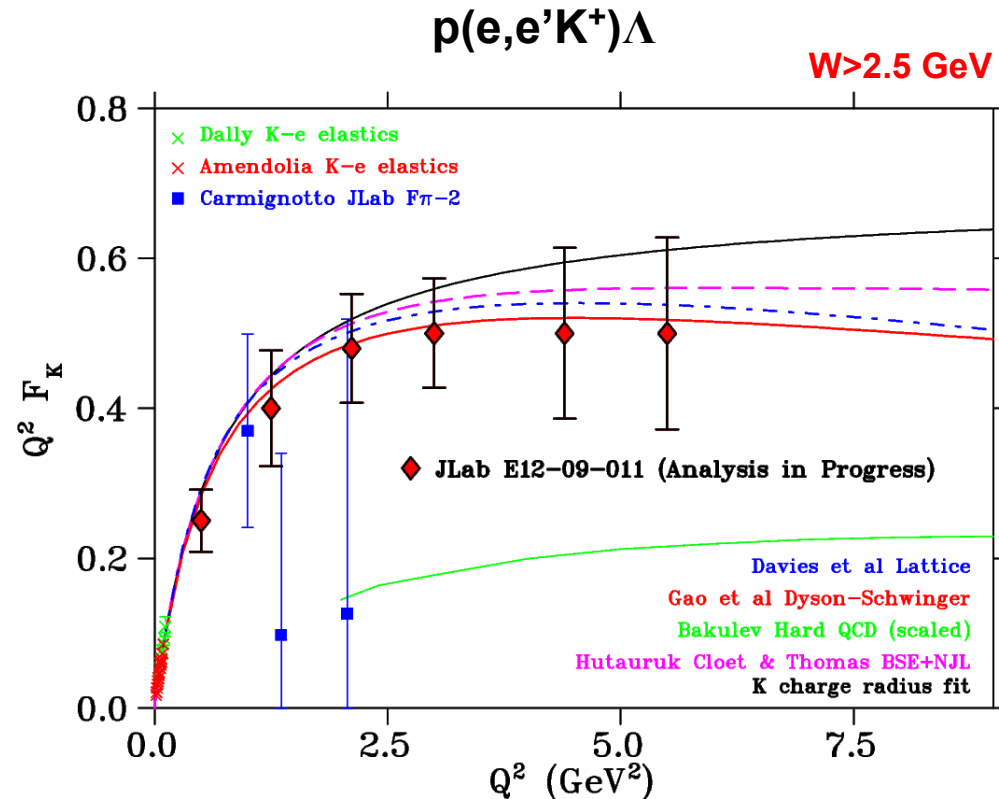


- F_K DCSB model prediction for JLab kinematics

[F. Guo, et al., arXiv: 1703.04875].

Projected Uncertainties for K^+ Form Factor

- First measurement of F_K well above the resonance region.
- Measure form factor to $Q^2=3 \text{ GeV}^2$ with good overlap with elastic scattering data.
 - Limited by $-t < 0.2 \text{ GeV}^2$ requirement to minimize non-pole contributions.
- Data will provide an important second $q\bar{q}$ system for theoretical models, this time involving a strange quark.



KaonLT E12-09-011: T. Horn, G. Huber and P. Markowitz, spokespersons

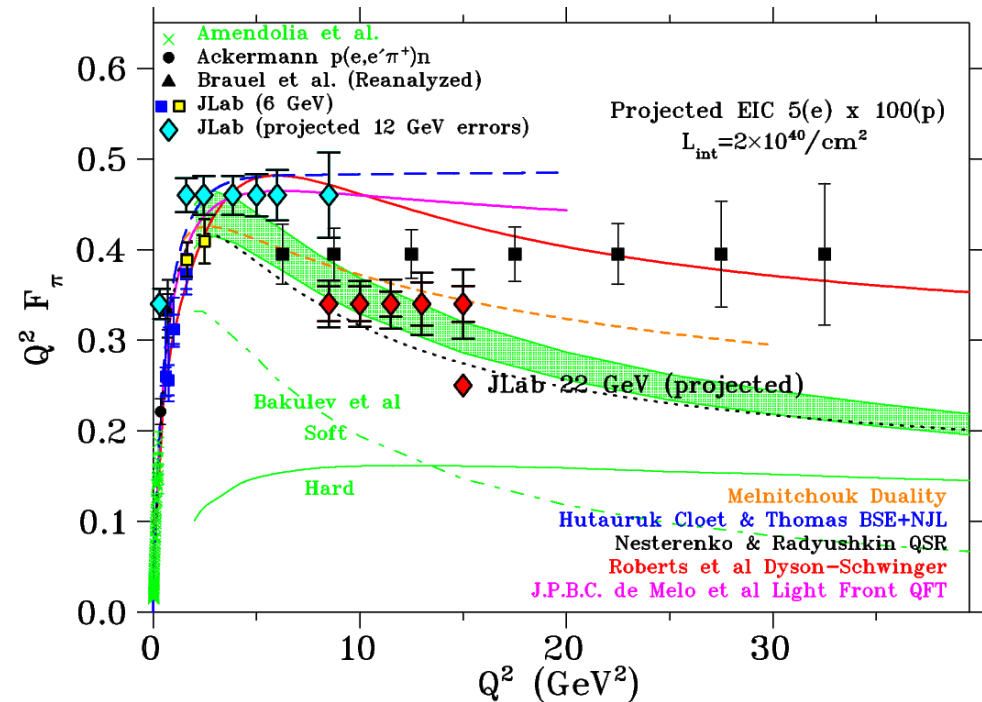
Phase 2 Scenario: π^+ Form Factor

■ Replace HMS with VHMS for π^+ , use SHMS for e'

- Assume $\theta_{\min} = 5.5^\circ$, $\theta_{\text{open}} = 15.0^\circ$
- VHMS: $\Delta\Omega$, $\Delta P/P$ similar SHMS

- $P_{\text{VHMS}} = 15.0$ GeV/c is sufficient, constrained by max beam energy
- $\theta_{\text{VHMS}} \sim 5.5^\circ$ allows improved $\Delta\varepsilon$, but does not affect maximum Q^2 reach

- Dramatic increase in upper Q^2
11.5 \rightarrow 15.0 GeV²
- Error bars for $Q^2 = 8.5\text{--}11.5$ GeV² substantially decrease due to smaller $-t_{\min}$ (better $R = \sigma_T / \sigma_L$) and shorter running times
- Highest Q^2 running time is “expensive” but would have very high scientific priority.



- Extends region of high quality F_π values to $Q^2 = 13$ GeV²
- Somewhat larger errors to $Q^2 = 15$ GeV²
- Provides MUCH improved overlap of F_π data set between JLab and EIC

- **Hall C is world's only facility that can do L–T separations over a wide kinematic range**
- The error magnification in L–T separations depends crucially on the achievable difference in the virtual photon polarization parameter, ε .
 - Errors magnify as $1/\Delta\varepsilon$, where $\Delta\varepsilon = \varepsilon_{\text{High}} - \varepsilon_{\text{Low}}$
 - To keep the magnification <500%, one desires $\Delta\varepsilon > 0.2$
 - This is not feasible at the EIC, as the high ion ring energy constrains $\varepsilon > 0.98$
- **As the interpretation of some EIC data (e.g. GPD extraction) will depend on extrapolation of Hall C L–T separated data, maximizing overlap between Hall C and EIC data sets should be a high priority**
 - An important motivation for extending reach of Hall C data using 22 GeV beam

■ Replace HMS with a higher momentum spectrometer

- For high z reactions, such as DEMP, usable beam energy constrained by sum of HMS+SHMS maximum momenta

- i.e. 22 GeV beam energy is a larger constraint than the maximum HMS momentum

■ New HMS would not extend the Q^2 reach beyond Scenario 1.

However, it would result in smaller errors due to larger $\Delta\varepsilon$ and faster high ε data rates

p(e,e' π^+)n Kinematics					
E_{beam}	$\theta_{\text{HMS}} (e')$	$P_{\text{HMS}} (e')$	$\theta_{\text{q(SHMS)}} (\pi^+)$	$P_{\text{SHMS}} (\pi^+)$	Time FOM
$Q^2=8.5$ $W=3.64$ $-t_{\text{min}}=0.24$ $\Delta\varepsilon=0.53$					
13.0	34.30	1.88	5.29	10.99	64.7
22.0	10.81	10.88	10.23	10.99	0.6
$Q^2=10.0$ $W=3.44$ $-t_{\text{min}}=0.37$ $\Delta\varepsilon=0.54$					
13.0	37.78	1.83	5.56	10.97	122.7
22.0	11.76	10.83	10.97	10.97	1.3
$Q^2=11.5$ $W=3.24$ $-t_{\text{min}}=0.54$ $\Delta\varepsilon=0.29$					
14.0	31.73	2.75	7.06	10.96	82.4
22.0	12.66	10.75	11.56	10.96	2.5

- **This scenario is judged to not be worth it, at least for this reaction channel**

Upgrade HMS Momentum and Angle: F_π

- Upgrade both HMS momentum and forward angle capabilities
 - 7 GeV/c \rightarrow 11 GeV/c
 - $\theta_{\min} = 10.50^\circ \rightarrow 7.5^\circ$
 - $\theta_{\text{open}} = 18.00^\circ \rightarrow 15.00^\circ$
- This upgrade also does not extend the Q^2 reach beyond Scenario 1.
- However, it would result in smaller errors due to larger $\Delta\varepsilon$ and faster high ε data rates

p(e,e' π^+)n Kinematics					
E_{beam}	$\theta_{\text{HMS}} (e')$	$P_{\text{HMS}} (e')$	$\theta_{q(\text{SHMS})} (\pi^+)$	$P_{\text{SHMS}} (\pi^+)$	Time FOM
$Q^2=8.5$ $W=3.64$ $-t_{\min}=0.24$ $\Delta\varepsilon=0.53$					
13.0	34.30	1.88	5.29	10.99	64.7
22.0	10.81	10.88	10.23	10.99	0.6
$Q^2=10.0$ $W=3.44$ $-t_{\min}=0.37$ $\Delta\varepsilon=0.54$					
13.0	37.78	1.83	5.56	10.97	122.7
22.0	11.76	10.83	10.97	10.97	1.3
$Q^2=11.5$ $W=3.24$ $-t_{\min}=0.54$ $\Delta\varepsilon=0.29$					
14.0	31.73	2.75	7.06	10.96	82.4
22.0	12.66	10.75	11.56	10.96	2.5

- Basically the same as Scenario 2. Not worth it, at least for this channel

- **Replace SHMS with higher momentum spectrometer, but keep HMS as is**
- Dramatic increase in upper Q^2 11.5 \rightarrow 15.0 GeV²
- Error bars for $Q^2=8.5-11.5$ GeV² would substantially decrease due to smaller $-t_{min}$ (better $R=\sigma_T/\sigma_L$) and shorter running times
- The $Q^2=15.0$ GeV² point would be “expensive” in terms of running time, but its high scientific priority would make it worthwhile
- **This seems a compelling scenario for a Phase 2 Upgrade**

p(e,e' π^+)n Kinematics					
E_{beam}	θ_{HMS} (e')	P_{HMS} (e')	$\theta_{q(SHMS)}$ (π^+)	P_{SHMS} (π^+)	Time FOM
$Q^2=8.5$ $W=4.06$ $-t_{min}=0.17$ $\Delta\varepsilon=0.26$					
16.0	23.68	3.15	5.52	12.75	17.7
20.0	14.00	7.15	7.55	12.75	1.9
$Q^2=10.0$ $W=3.96$ $-t_{min}=0.23$ $\Delta\varepsilon=0.28$					
16.0	27.41	2.78	5.41	13.09	47.7
20.0	15.60	6.78	7.72	13.09	4.5
$Q^2=11.5$ $W=3.96$ $-t_{min}=0.29$ $\Delta\varepsilon=0.27$					
17.0	27.54	2.98	5.49	13.86	76.3
21.0	16.10	6.98	7.72	13.86	8.1
$Q^2=13.0$ $W=3.96$ $-t_{min}=0.35$ $\Delta\varepsilon=0.25$					
18.0	27.55	3.18	5.54	14.63	123.6
22.0	16.49	7.18	7.69	14.63	14.4
$Q^2=15.0$ $W=3.78$ $-t_{min}=0.50$ $\Delta\varepsilon=0.27$					
18.0	31.30	2.86	5.46	14.87	391
22.0	18.14	6.86	7.86	14.87	41.4