

# Deep Exclusive $\pi^-$ Production using a Transversely Polarized $^3\text{He}$ Target and the SoLID Spectrometer

Garth Huber



University  
of Regina

APS April Meeting  
April 17, 2021

Supported by:



SAPIN-2016-00031

- GPDs interrelate the longitudinal momentum and transverse spatial structure of partons within a fast moving hadron.
- GPDs are universal quantities and reflect nucleon structure independently of the probing reaction.

- At leading twist–2, four quark chirality conserving GPDs for each quark, gluon type.
- Because quark helicity is conserved in the hard scattering regime, the produced meson acts as helicity filter.
  - Pseudoscalar mesons  $\rightarrow \tilde{H} \tilde{E}$
  - Vector mesons  $\rightarrow H E$ .

$H^{q,g}(x, \xi, t)$   
spin avg  
no hel. flip

$E^{q,g}(x, \xi, t)$   
spin avg  
helicity flip

$\tilde{H}^{q,g}(x, \xi, t)$   
spin diff  
no hel. flip

$\tilde{E}^{q,g}(x, \xi, t)$   
spin diff  
helicity flip

- Additional chiral–odd GPDs ( $H_T E_T \tilde{H}_T \tilde{E}_T$ ) offer a new way to access the transversity–dependent quark–content of the nucleon.

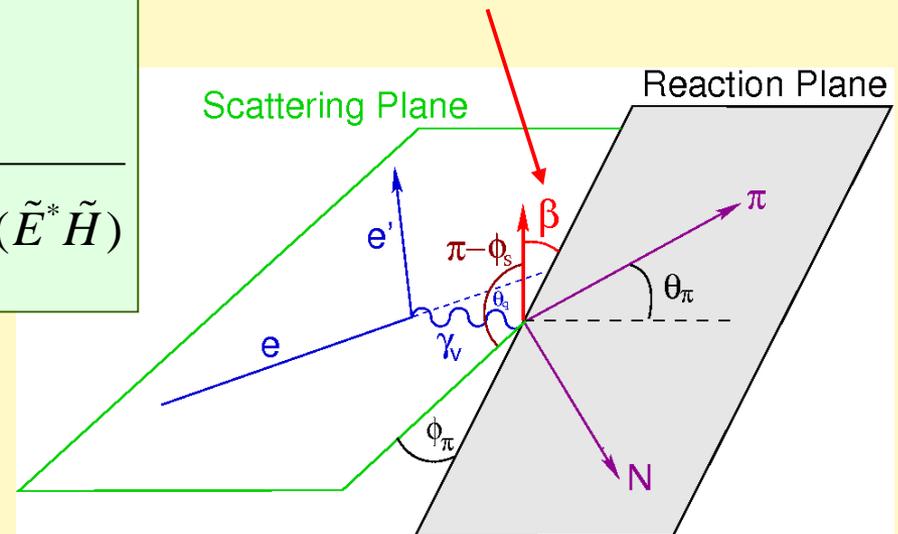
# How to determine $\tilde{E}$

- **GPD  $\tilde{E}$  not related to an already known parton distribution.**
- **Experimental information on  $\tilde{E}$  can provide new nucleon structure info unlikely to be available from any other source.**
- The most sensitive observable to probe  $\tilde{E}$  is the transverse single-spin asymmetry in exclusive  $\pi$  production:

$$A_L^\perp = \frac{\left( \int_0^\pi d\beta \frac{d\sigma_L^\pi}{d\beta} - \int_\pi^{2\pi} d\beta \frac{d\sigma_L^\pi}{d\beta} \right)}{\left( \int_0^{2\pi} d\beta \frac{d\sigma_L^\pi}{d\beta} \right)}$$

$$= \frac{\sqrt{-t'}}{2m_p} \frac{\pi\xi\sqrt{1-\xi^2} \text{Im}(\tilde{E}^* \tilde{H})}{(1-\xi^2)\tilde{H}^2 - \frac{t\xi^2}{4m_p} \tilde{E}^2 - 2\xi^2 \text{Re}(\tilde{E}^* \tilde{H})}$$

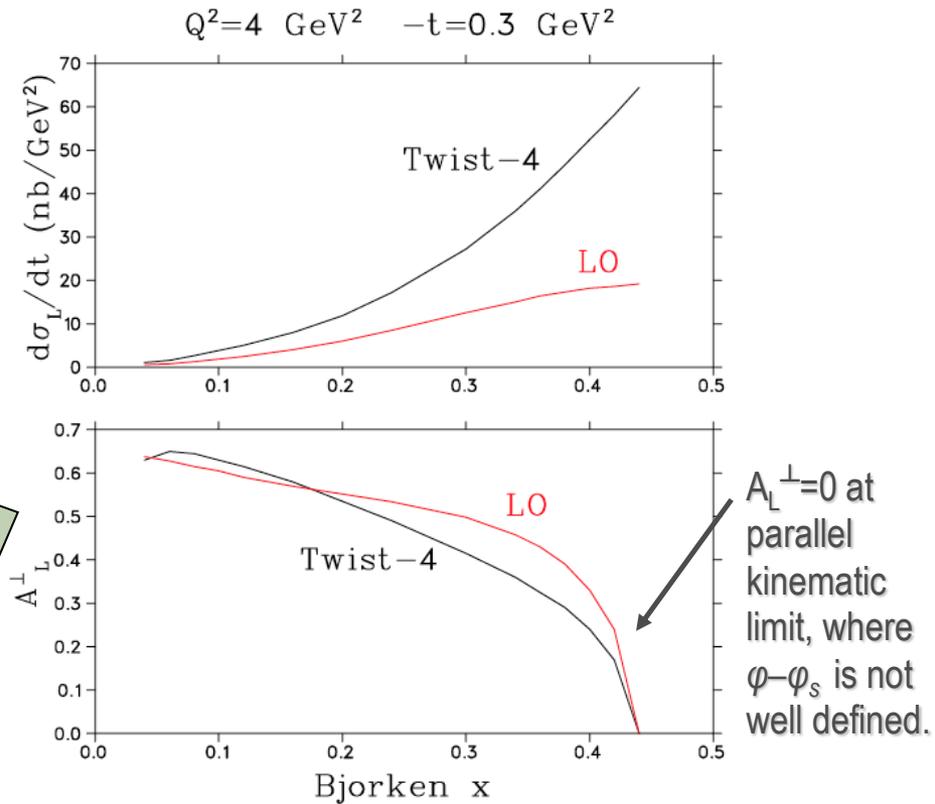
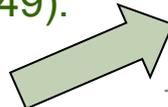
$d\sigma_\pi^L$  = exclusive  $\pi$  cross section for longitudinal  $\gamma^*$   
 $\beta$  = angle between transversely polarized target vector and the reaction plane.



Refs: A.V. Belitsky, D. Mueller, PLB513 (2001) 349  
 L.L. Frankfurt, et al., PRD 60(1999) 014101

- $A_L^\perp$  is expected to display precocious factorization at only  $Q^2 \sim 2-4 \text{ GeV}^2$ :**

- At  $Q^2=10 \text{ GeV}^2$ , Twist-4 effects can be large, but cancel in  $A_L^\perp$  (Belitsky & Müller PLB 513(2001)349).
- At  $Q^2=4 \text{ GeV}^2$ , higher twist effects even larger in  $\sigma_L$ , but still cancel in the asymmetry (CIPANP 2003).



**This relatively low value of  $Q^2$  for the expected onset of precocious scaling is important, because it is experimentally accessible at JLab 12 GeV.**

# Transverse Target Single Spin Asymmetry in DEMP

Unpolarized Cross section  $2\pi \frac{d^2 \sigma_{UU}}{dtd\phi} = \varepsilon \frac{d\sigma_L}{dt} + \frac{d\sigma_T}{dt} + \sqrt{2\varepsilon(\varepsilon+1)} \frac{d\sigma_{LT}}{dt} \cos \phi + \varepsilon \frac{d\sigma_{TT}}{dt} \cos 2\phi$

Transversely polarized cross section has additional components

$$\frac{d^3 \sigma_{UT}}{dtd\phi d\phi_s} = - \frac{P_{\perp} \cos \theta_q}{\sqrt{1 - \sin^2 \theta_q \sin^2 \phi_s}}$$

Gives rise to Asymmetry Moments

$$A(\phi, \phi_s) = \frac{d^3 \sigma_{UT}(\phi, \phi_s)}{d^2 \sigma_{UU}(\phi)} = - \sum_k A_{UT}^{\sin(\mu\phi + \lambda\phi_s)_k} \sin(\mu\phi + \lambda\phi_s)_k$$

$$\left( \begin{aligned} & \sin \beta \text{Im}(d\sigma_{++}^{+-} + \varepsilon d\sigma_{00}^{+-}) \\ & + \sin \phi \sqrt{\varepsilon(1+\varepsilon)} \text{Im}(d\sigma_{+0}^{+-}) \\ & + \sin(\phi + \phi_s) \frac{\varepsilon}{2} \text{Im}(d\sigma_{+-}^{+-}) \\ & + \sin(2\phi - \phi_s) \sqrt{\varepsilon(1+\varepsilon)} \text{Im}(d\sigma_{+0}^{-+}) \\ & + \sin(3\phi - \phi_s) \frac{\varepsilon}{2} \text{Im}(d\sigma_{+-}^{-+}) \end{aligned} \right)$$

$\sigma_{mn}^{ij} \rightarrow$  nucleon polarizations  $ij = (+1/2, -1/2)$   
 photon polarizations  $mn = (-1, 0, +1)$

Unseparated  $\sin\beta = \sin(\phi - \phi_s)$  Asymmetry Moment

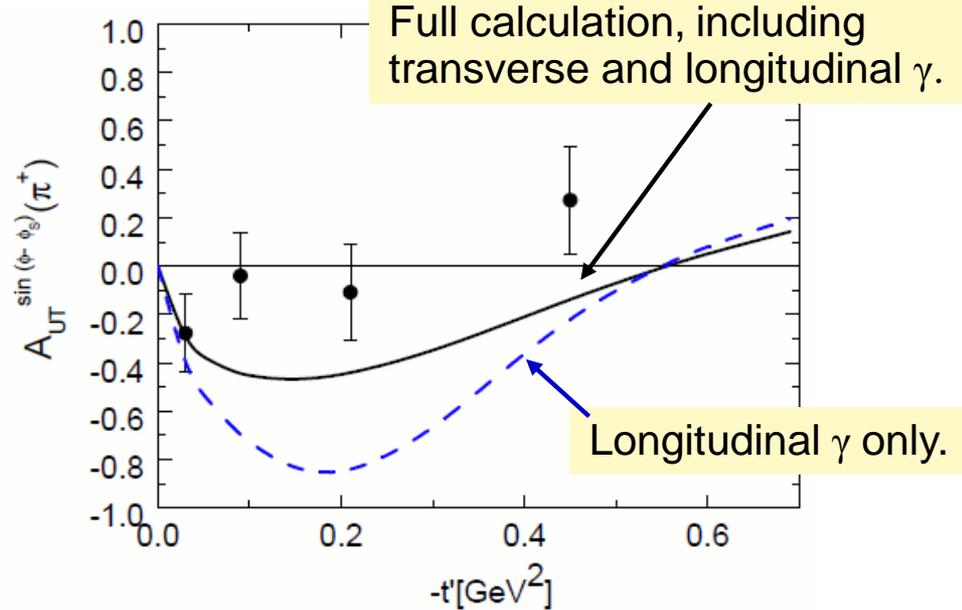
$$A_{UT}^{\sin(\phi - \phi_s)} \sim \frac{d\sigma_{00}^{+-}}{d\sigma_L \binom{++}{00}} \sim \frac{\text{Im}(\tilde{E}^* \tilde{H})}{|\tilde{E}|^2} \text{ where } \tilde{E} \gg \tilde{H}$$

Ref: M. Diehl, S. Sapeta, Eur.Phys.J. C41(2005)515.

Note: Trento convention used for rest of talk

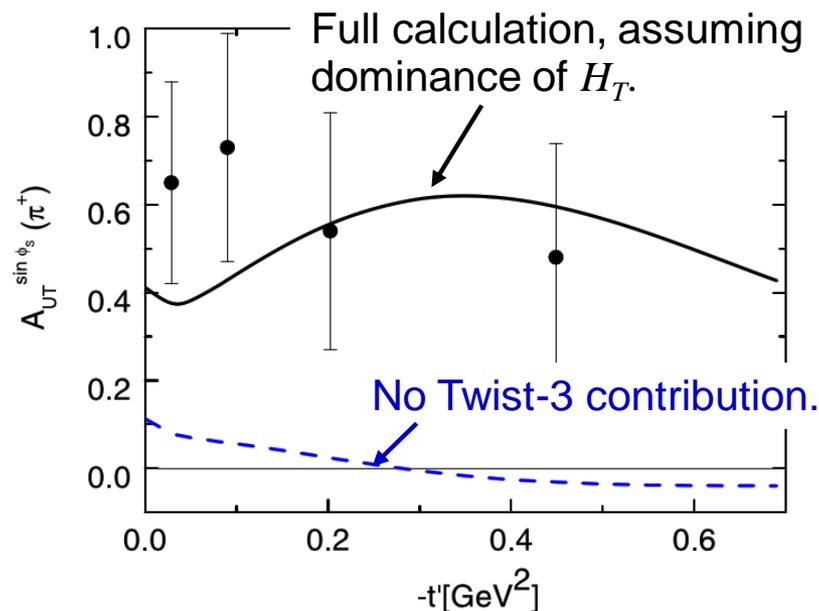
# HERMES $\sin(\varphi-\varphi_S)$ Asymmetry Moment

- Exclusive  $\pi^+$  production by scattering 27.6 GeV positrons or electrons from transverse polarized  $^1\text{H}$  [PL **B682**(2010)345].
- Analyzed in terms of 6 Fourier amplitudes for  $\varphi_\pi, \varphi_S$ .
- $\langle x_B \rangle = 0.13$ ,  $\langle Q^2 \rangle = 2.38 \text{ GeV}^2$ ,  $\langle -t \rangle = 0.46 \text{ GeV}^2$ .



- **Since there is no L/T separation,  $A_{UT}^{\sin(\varphi-\varphi_S)}$  is diluted by the ratio of the longitudinal cross section to the unseparated cross section.**
- Goloskokov and Kroll indicate the HERMES results have significant contributions from transverse photons, as well as from L and T interferences [Eur Phys.J. **C65**(2010)137].
- **Because no factorization theorems exist for exclusive  $\pi$  production by transverse photons, these data cannot be trivially interpreted in terms of GPDs.**

- While most of the theoretical interest and the primary motivation of our experiment is the  $\sin(\varphi-\varphi_s)$  asymmetry moment, there is growing interest in the  $\sin(\varphi_s)$  moment, which may be interpretable in terms of the transversity GPDs.



- In contrast to the  $\sin(\varphi-\varphi_s)$  modulation, which has contributions from LL and TT interferences, the  $\sin(\varphi_s)$  modulation measures only the LT interference.
- **The HERMES  $\sin(\varphi_s)$  modulation is large and nonzero at  $-t'=0$ , giving the first clear signal for strong contributions from transversely polarized photons at rather large values of  $W$  and  $Q^2$ .**
- **Goloskokov and Kroll calculation [Eur.Phys.J. C65(2010)137] assumes the transversity GPD  $H_T$  dominates and that the other three can be neglected.**

# Measure DEMP with SoLID – Polarized $^3\text{He}$

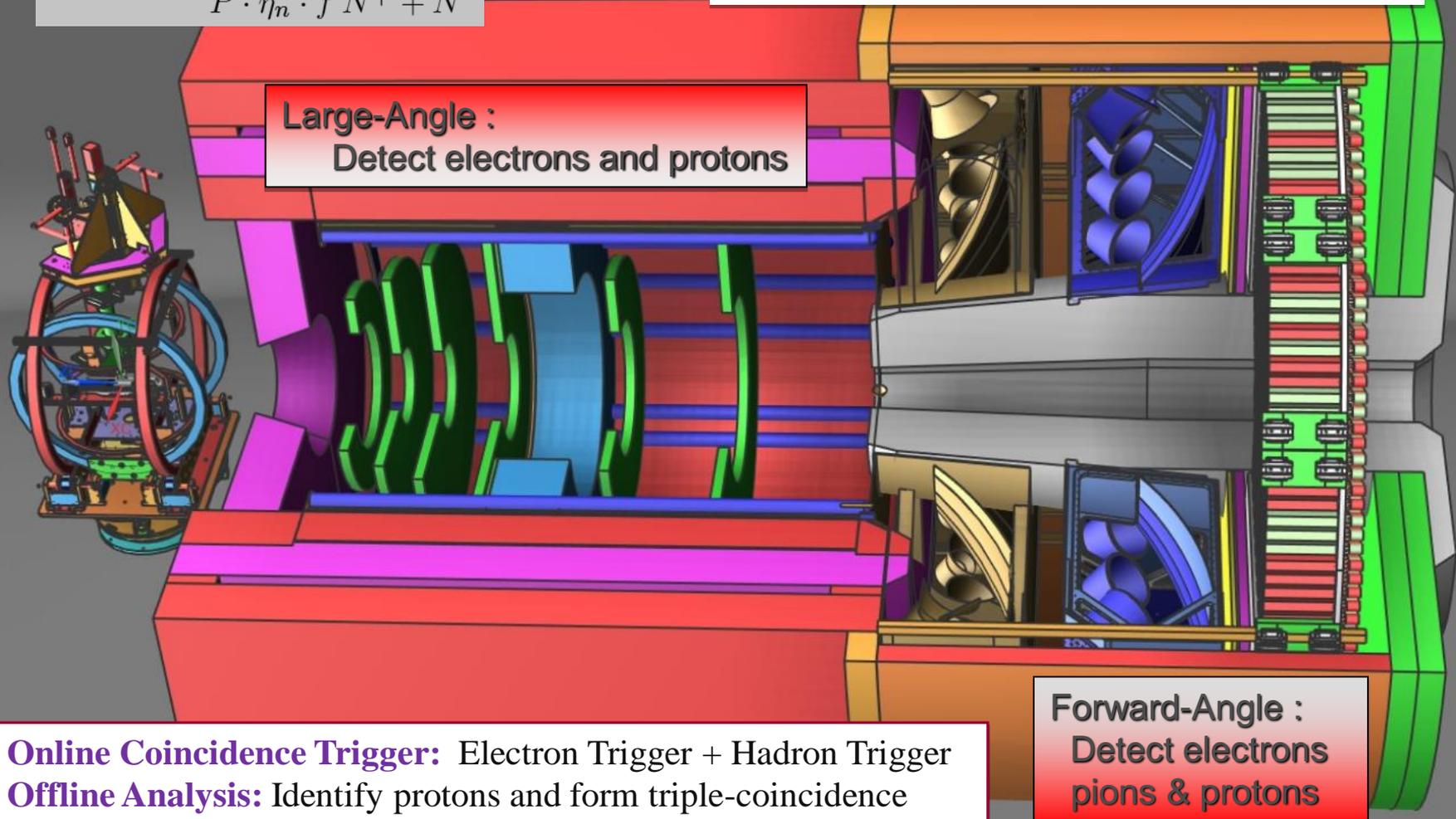
$\vec{n}(e, e' \pi^-)p$ : with transversely polarized  $^3\text{He}$

$$\langle A_{UT} \rangle = \frac{1}{P \cdot \eta_n \cdot f} \frac{N^+ - N^-}{N^+ + N^-}$$

Run in parallel with E12-10-006:

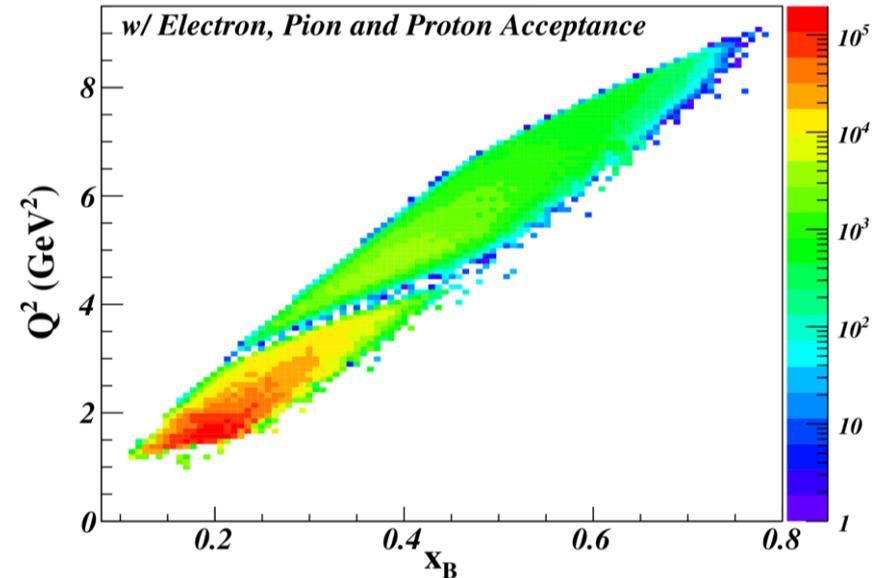
$E_0 = 11.0 \text{ GeV}$  (48 days)

Luminosity =  $10^{36} \text{ cm}^{-2} \text{ s}^{-1}$  (per nucleon)

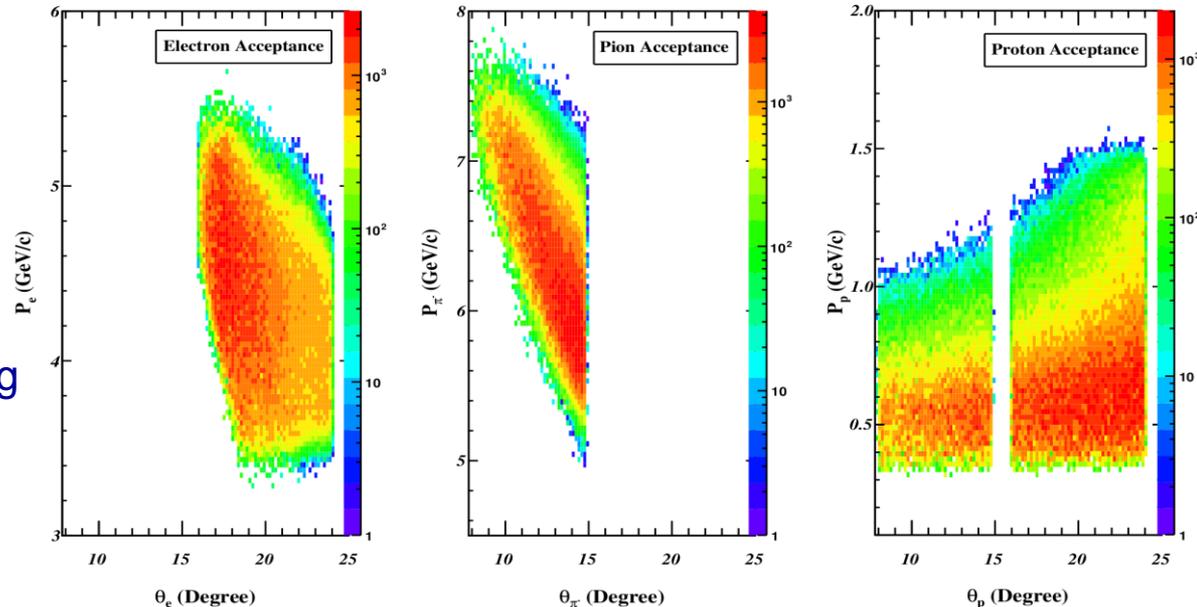


# SoLID Acceptance and Projected Rates

$Q^2 > 1 \text{ GeV}^2$ $W > 2 \text{ GeV}$	$Q^2 > 4 \text{ GeV}^2$ $W > 2 \text{ GeV}$
DEMP: $n(e, e' \pi^- p)$ Triple Coin (Hz)	
4.95	0.40
SIDIS: $n(e, e' \pi^-) X$ Double Coin (Hz)	
1425	35.8



$Q^2 > 4 \text{ GeV}^2$ ,  $W > 2 \text{ GeV}$ ,  $0.55 < \epsilon < 0.75$  cuts applied.

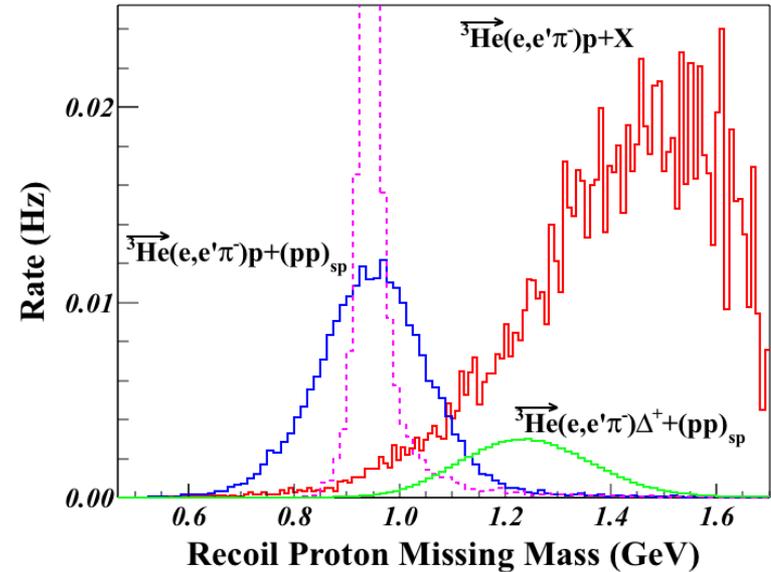


- Event generator is based on data from HERMES, Halls B,C with VR Regge+DIS model used as a constraint in unmeasured regions.
- Generator includes electron radiation, multiple scattering and ionization energy loss.
- Every detected particle is smeared in  $(P, \theta, \phi)$  with resolution from SoLID tracking studies, and acceptance profiles from SoLID-SIDIS GEMC study applied.

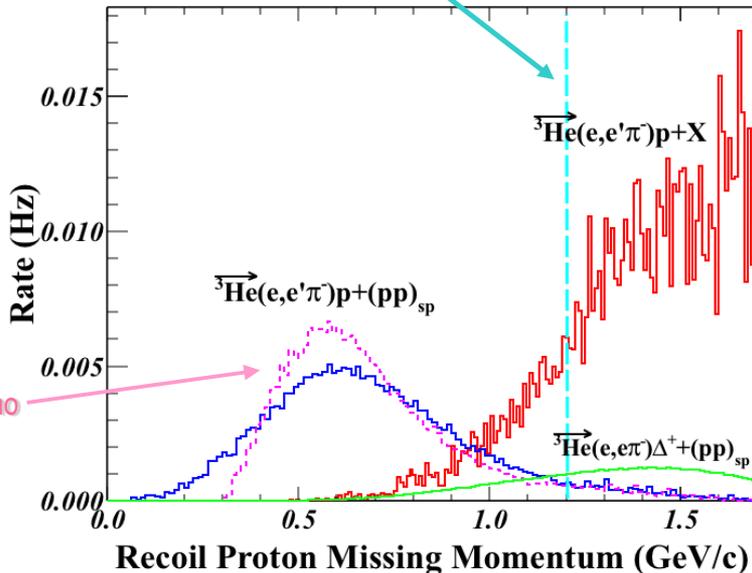
# Example Cuts to Reduce Background

## Two different background channels were simulated:

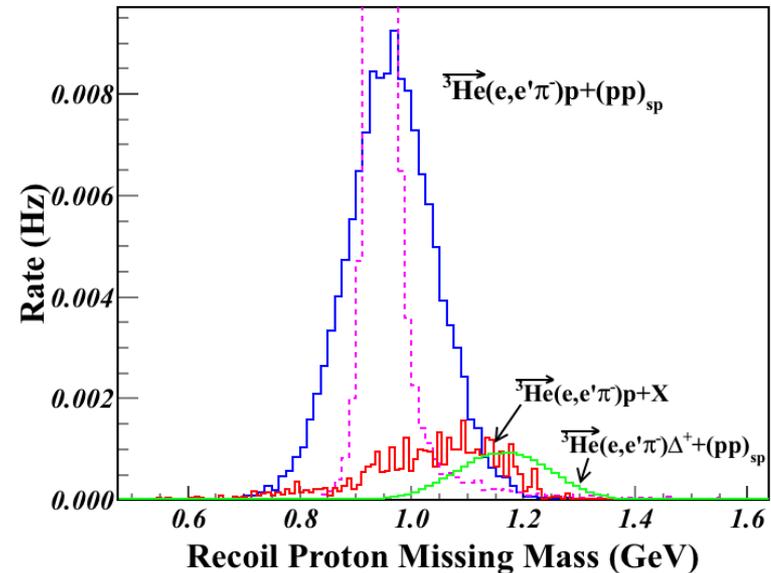
- SoLID–SIDIS generator  $p(e, e' \pi^-)X$  and  $n(e, e' \pi^-)X$ , where we assume all  $X$  fragments contain a proton (over-estimate).
- $en \rightarrow \pi^- \Delta^+ \rightarrow \pi^- \pi^0 p$  where the  $\Delta^+$  (polarized) decays with  $l=1, m=0$  angular distribution (more realistic).



Apply  $P_{miss} > 1.2$  GeV/c cut

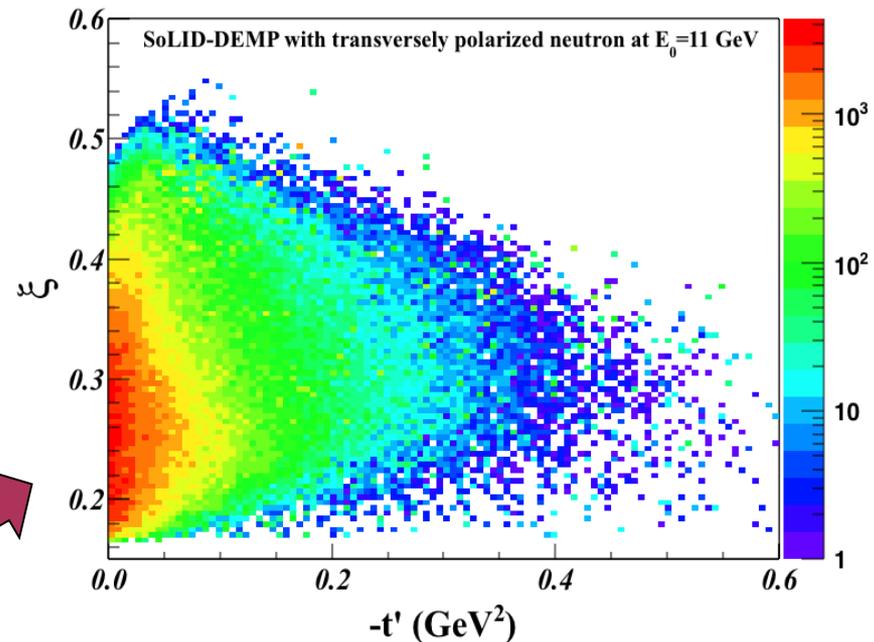
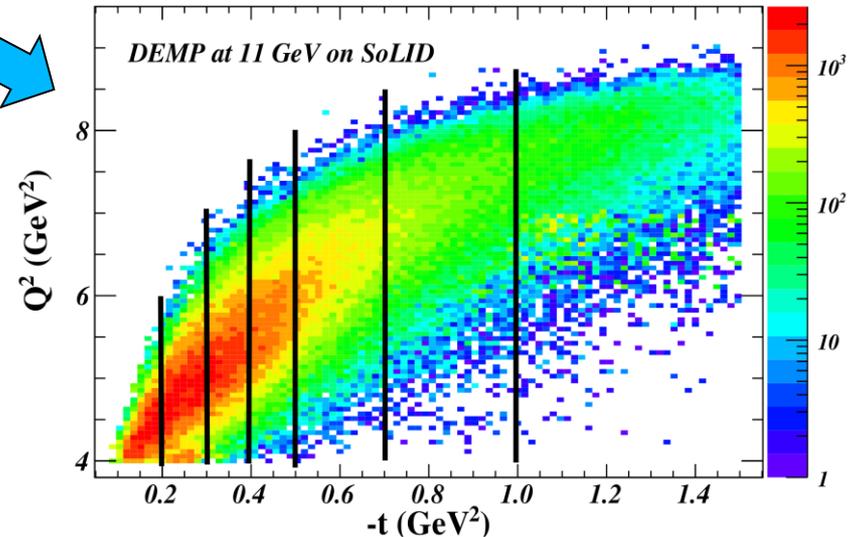


Background remaining after  $P_{miss}$  cut



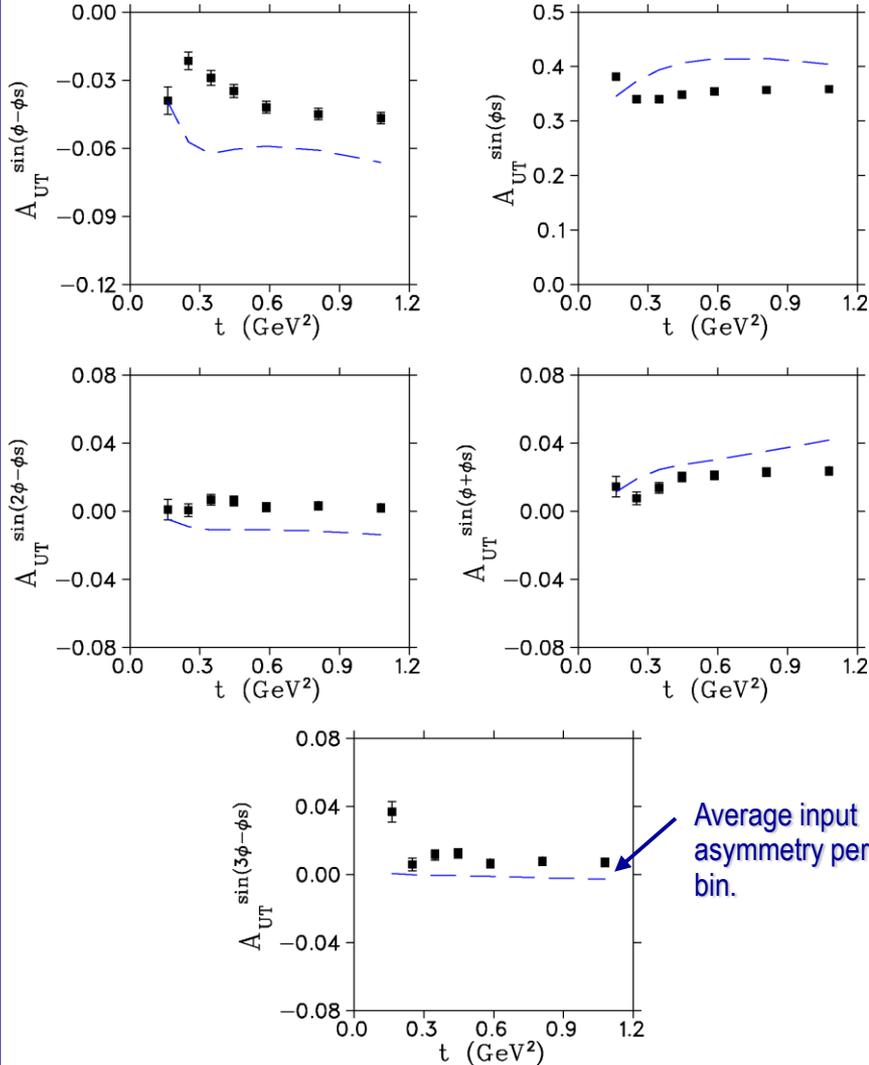
# Kinematic Coverage and Binning

- We binned the simulated data in 7  $t$ -bins.
- In actual data analysis, we will consider alternate binning.
- All JLab data cover a range of  $Q^2$ ,  $x_{Bj}$  values.
  - $x_{Bj}$  fixes the skewness ( $\xi$ ).
  - $Q^2$  and  $x_{Bj}$  are correlated. In fact, we have an almost linear dependence of  $Q^2$  on  $x_{Bj}$ .
- HERMES and COMPASS experiments are restricted kinematically to very small skewness ( $\xi < 0.1$ ).
- With SoLID, we can measure the skewness dependence of the relevant GPDs over a fairly large range of  $\xi$ .



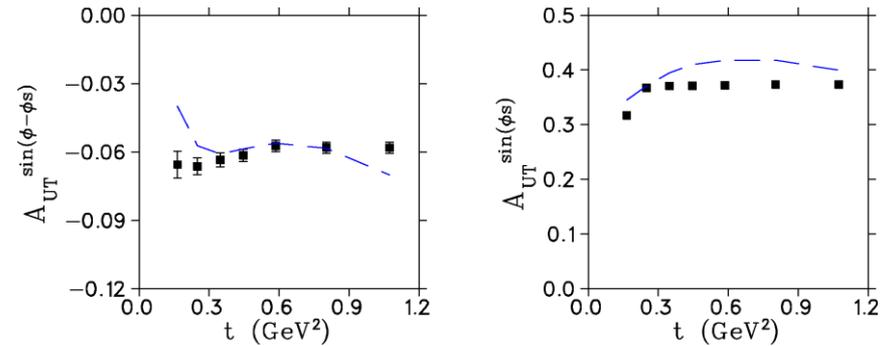
## All effects on.

Includes all scattering, energy loss, resolution and Fermi momentum effects.



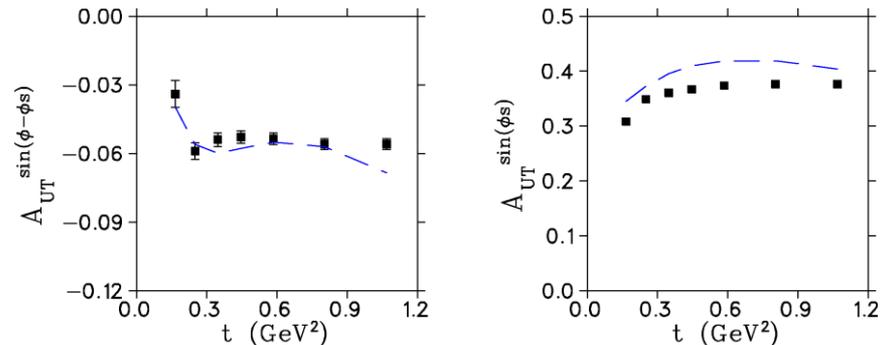
## Only Fermi momentum off.

Includes all scattering, energy loss, resolution effects. Similar to where proton resolution is good enough to correct for Fermi momentum effects.



## All effects off.

- Agreement between input and output fit values is very good. Validates the Unbinned Maximum Likelihood analysis procedure.



- $A_{UT}^{\sin(\varphi-\varphi_S)}$  transverse single–spin asymmetry in exclusive  $\pi$  production is particularly sensitive to the spin–flip GPD  $\tilde{E}$ . Factorization studies indicate precocious scaling to set in at moderate  $Q^2 \sim 2\text{--}4 \text{ GeV}^2$ , while scaling is not expected until  $Q^2 > 10 \text{ GeV}^2$  for absolute cross section.
- $A_{UT}^{\sin(\varphi_S)}$  asymmetry can also be extracted from same data, providing powerful additional GPD–model constraints and insight into the role of transverse photon contributions at small  $-t$ , and over wide range of  $\xi$ .
- **High luminosity and good acceptance capabilities of SoLID make it well–suited for this measurement. It is the only feasible manner to access the wide  $-t$  range needed to fully understand the asymmetries.**
- **SoLID measurement is also important preparatory work for EIC.**