

# Deep Exclusive $p(e, e' \pi^+)n$ and $p(e, e' K^+)\Lambda$ Studies at Jefferson Lab

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of Regina

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On behalf of the PionLT and KaonLT Collaborations

Towards Improved Hadron Femtography with Hard Exclusive Reactions  
ECT\* Workshop, Trento, Italy  
August 9, 2024

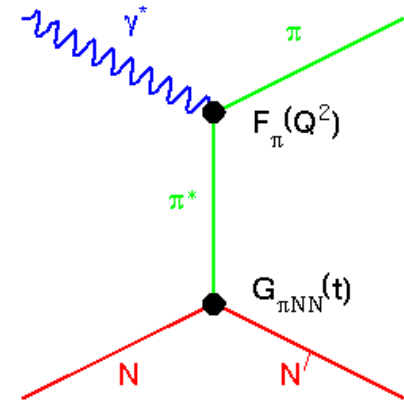
# Two Motivations for Studying DEMP

## 1) Determine the Pion Form Factor at $Q^2 > 0.3 \text{ GeV}^2$ :

- Indirectly measure  $F_\pi$  using the “pion cloud” of the proton via  $p(e, e' \pi^+) n$

$$|p\rangle = |p\rangle_0 + |n\pi^+\rangle + \dots$$

- Pion pole process dominates  $\sigma_L$  in forward kinematics.
- Can a similar method be used to determine the kaon form factor?



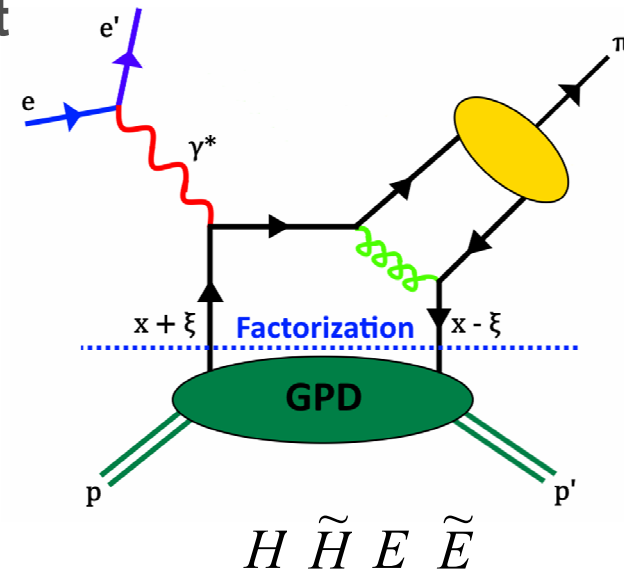
## 2) Study the Hard-Soft Factorization Regime:

Implications for GPD studies, as they can only be extracted from hard exclusive data where hard-soft factorization applies.

- Investigate if  $p(e, e' \pi^+) n$  and  $p(e, e' K^+) \Lambda$  cross sections at fixed  $x$  behave according to the  $Q^{-n}$  scaling expectations of hard QCD.

$$\frac{\sigma_T[n(e, e' \pi^-) p]}{\sigma_T[p(e, e' \pi^+) n]}$$

- Form ratios where soft contributions may cancel, yielding insight to factorization at modest  $Q^2$ .

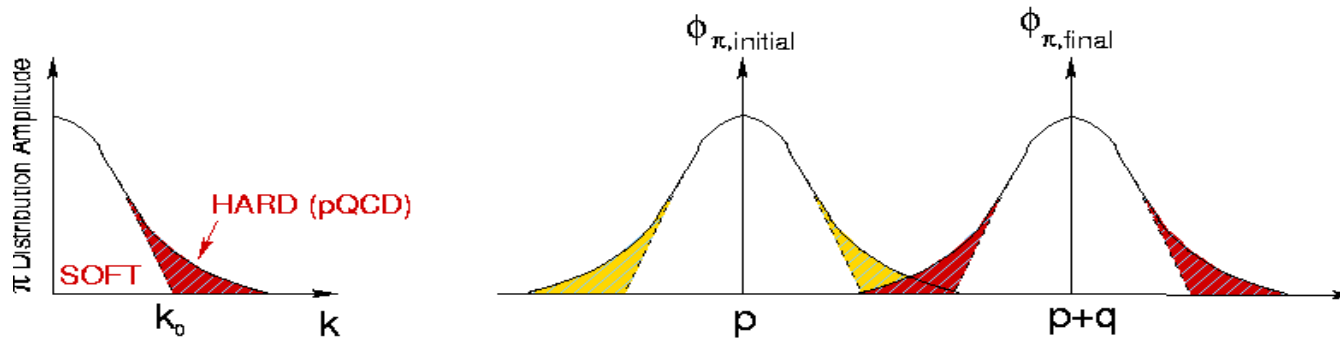


# Charged Meson Form Factors

Simple  $q\bar{q}$  valence structure of mesons presents the ideal testing ground for our understanding of bound quark systems.

In quantum field theory, the form factor is the overlap integral:

$$F_\pi(Q^2) = \int \phi_\pi^*(p) \phi_\pi(p+q) dp$$



The meson wave function can be separated into  $\phi_\pi^{\text{soft}}$  with only low momentum contributions ( $k < k_0$ ) and a hard tail  $\phi_\pi^{\text{hard}}$ .

While  $\phi_\pi^{\text{hard}}$  can be treated in pQCD,  $\phi_\pi^{\text{soft}}$  cannot.

**From a theoretical standpoint, the study of the  $Q^2$ -dependence of the form factor focuses on finding a description for the hard and soft contributions of the meson wave-function.**

# The Pion in perturbative QCD

At very large  $Q^2$ , pion form factor ( $F_\pi$ ) can be calculated using pQCD

$$F_\pi(Q^2) = \frac{4\pi C_F \alpha_s(Q^2)}{Q^2} \left| \sum_{n=0}^{\infty} a_n \left( \log \left( \frac{Q^2}{\Lambda^2} \right) \right)^{-\gamma_n} \right|^2 \left[ 1 + O \left( \alpha_s(Q^2), \frac{m}{Q} \right) \right]$$

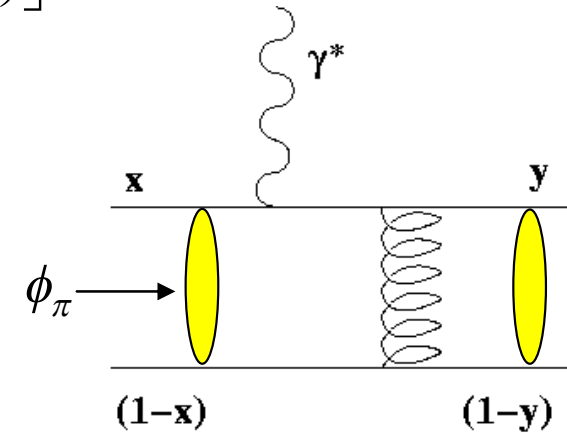
At asymptotically high  $Q^2$ , only hardest portion of pion distribution amplitude contributes

$$\phi_\pi(x) \xrightarrow{Q^2 \rightarrow \infty} \frac{3f_\pi}{\sqrt{n_c}} x(1-x)$$

and  $F_\pi$  takes the very simple form

$$Q^2 F_\pi(Q^2) \xrightarrow{Q^2 \rightarrow \infty} 16\pi \alpha_s(Q^2) f_\pi^2$$

$f_\pi = 93$  MeV is the  $\pi^+ \rightarrow \mu^+ \nu$  decay constant



G.P. Lepage, S.J. Brodsky, Phys.Lett. **87B**(1979)359.

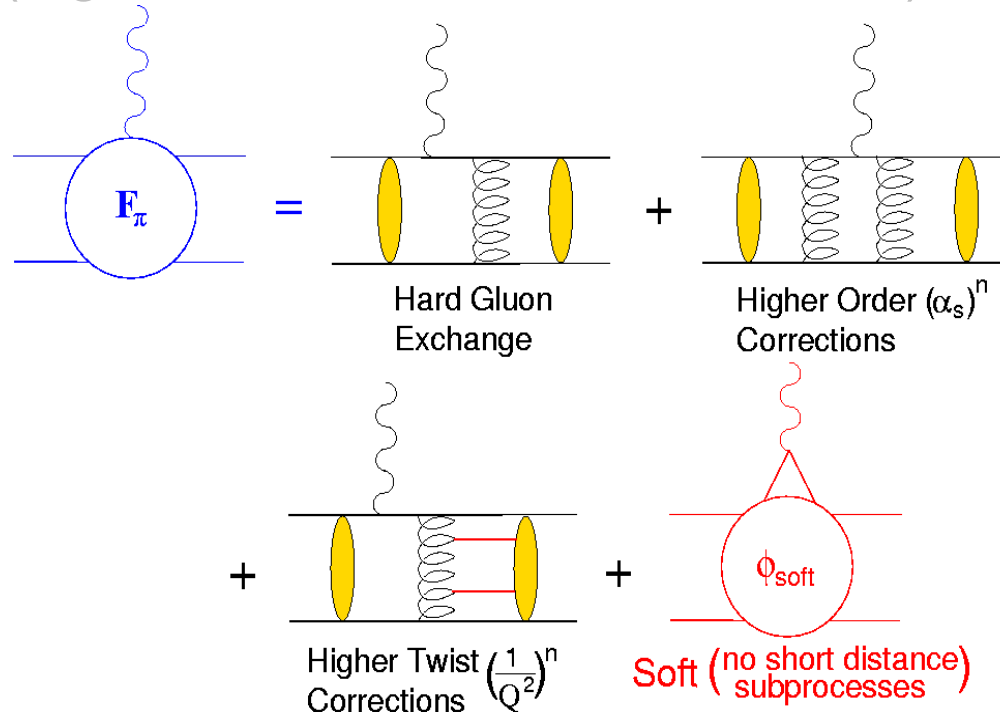
This only relies on asymptotic freedom in QCD, *i.e.*  $(\partial\alpha_s/\partial\mu) < 0$  as  $\mu \rightarrow \infty$ .

**$Q^2 F_\pi$  should behave like  $\alpha_s(Q^2)$  even for moderately large  $Q^2$ .**

→ Pion form factor seems to be best tool for experimental study of nature of the quark-gluon coupling constant renormalization.

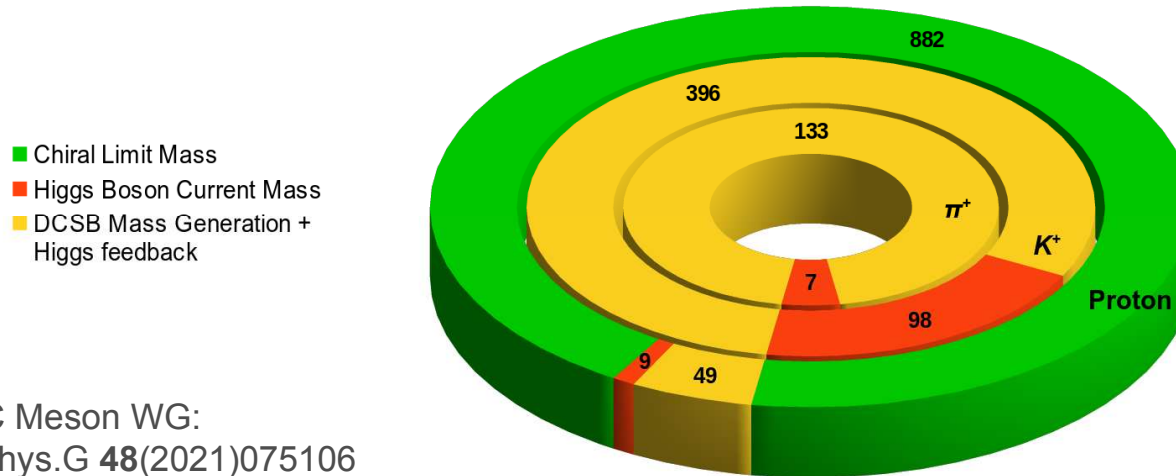
[A.V. Radyushkin, JINR 1977, arXiv:hep-ph/0410276]

At experimentally-accessible  $Q^2$ , both the “hard” and “soft” components (e.g. transverse momentum effects) contribute.



- **The interplay of hard and soft contributions is poorly understood.**
  - Different theoretical viewpoints on whether higher-twist mechanisms dominate until very large momentum transfer or not.
- **The pion elastic and transition form factors experimentally accessible over a wide kinematic range.**
  - A laboratory to study the **transition** from the soft to hard regime.

Hadron Mass Budget



EIC Meson WG:  
J.Phys.G 48(2021)075106

## Stark Differences between proton, $K^+$ , $\pi^+$ mass budgets

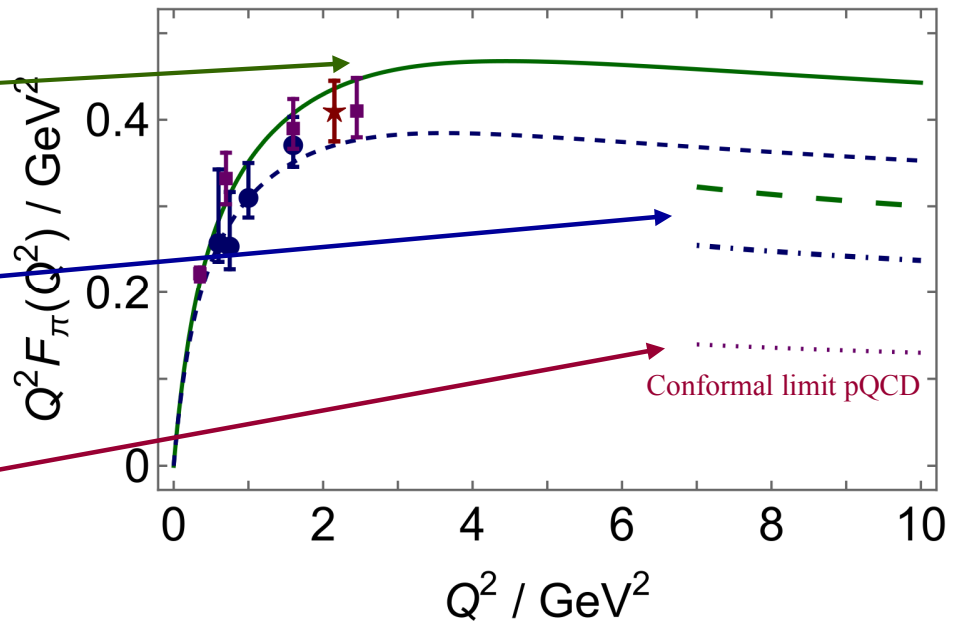
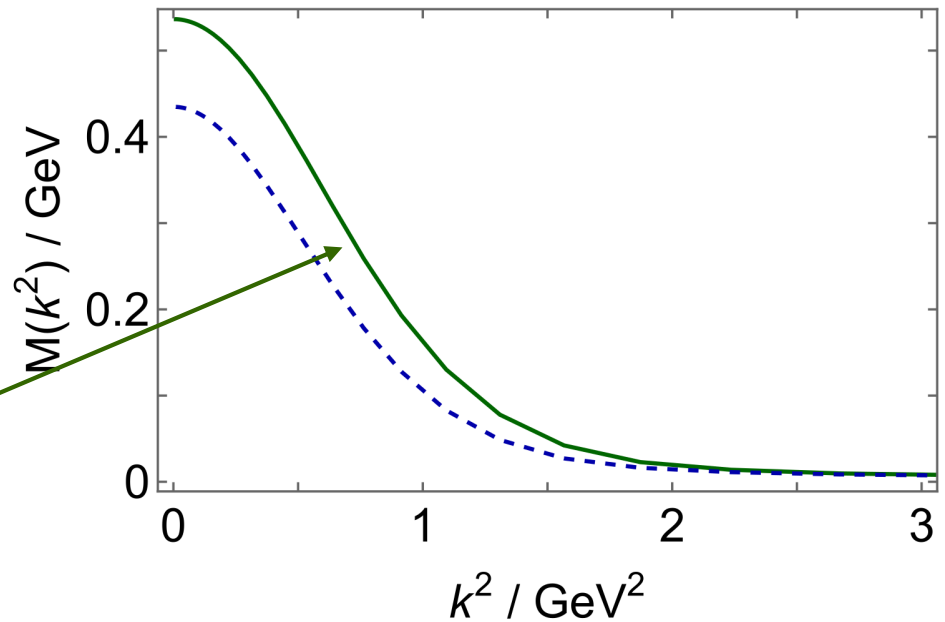
- Due to Emergent Hadronic Mass (EHM), Proton mass large in absence of quark couplings to Higgs boson (chiral limit).
- Conversely, and yet still due to EHM and DCSB,  $K$  and  $\pi$  are massless in chiral limit (i.e. they are Goldstone bosons).
- The mass budgets of these crucially important particles demand interpretation.
- Equations of QCD stress that any explanation of the proton's mass is incomplete, unless it simultaneously explains the light masses of QCD's Goldstone bosons, the  $\pi$  and  $K$ .

# Synergy: Emergent Mass and $\pi^+$ Form Factor

At empirically accessible energy scales,  $\pi^+$  form factor is sensitive to emergent mass scale in QCD

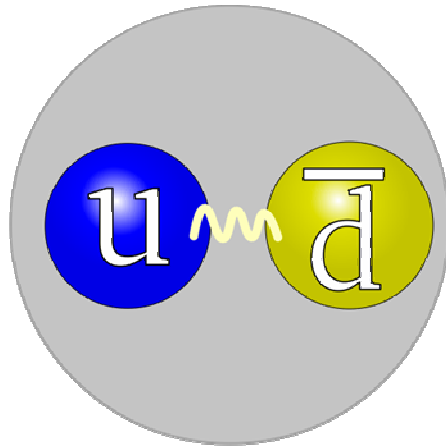
- Two dressed-quark mass functions distinguished by amount of DCSB
  - DCSB emergent mass generation is 20% stronger in system characterized by solid green curve, which is more realistic case
- $F_\pi(Q^2)$  obtained with these mass functions
  - $r_\pi=0.66$  fm with solid green curve
  - $r_\pi=0.73$  fm with solid dashed blue curve
- $F_\pi(Q^2)$  predictions from QCD hard scattering formula (slide #3), obtained with related, computed pion PDAs
- QCD hard scattering formula, using conformal limit of pion's twist-2 PDA

$$\phi_\pi^{cl}(x) = 6x(1-x)$$

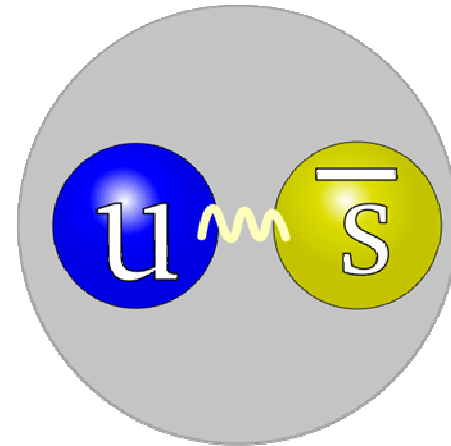


Chen, et al., PRD 98(2018)091505(R); Aguilar et al, EPJA 55(2019)190

# The Charged Kaon – a 2<sup>nd</sup> QCD test case



$\pi^+$



$K^+$

- In the hard scattering limit, pQCD predicts that the  $\pi^+$  and  $K^+$  form factors will behave similarly

$$\frac{F_K(Q^2)}{F_\pi(Q^2)} \xrightarrow{Q^2 \rightarrow \infty} \frac{f_K^2}{f_\pi^2}$$

- It is important to compare the magnitudes and  $Q^2$ -dependences of both form factors.



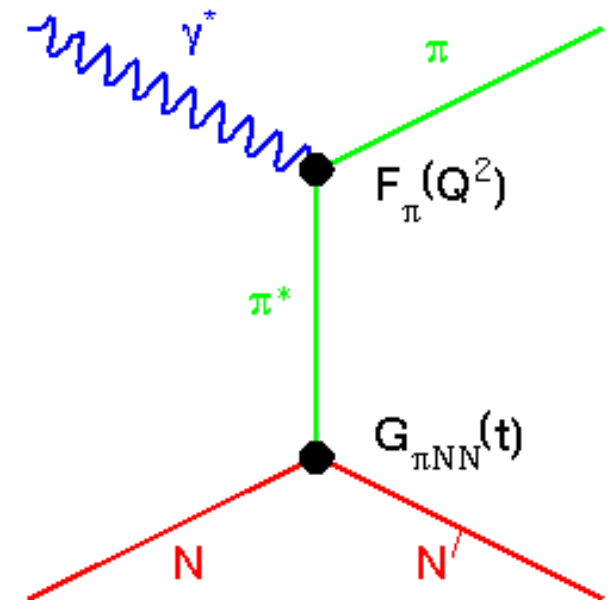
# Measurement of $\pi^+$ Form Factor – Larger $Q^2$

At larger  $Q^2$ ,  $F_\pi$  must be measured indirectly using the “pion cloud” of the proton via pion electroproduction  $p(e, e'\pi^+)n$

$$|p\rangle = |p\rangle_0 + |n\pi^+\rangle + \dots$$

- At small  $-t$ , the pion pole process dominates the longitudinal cross section,  $\sigma_L$
- In Born term model,  $F_\pi^2$  appears as,

$$\frac{d\sigma_L}{dt} \propto \frac{-tQ^2}{(t - m_\pi^2)} g_{\pi NN}^2(t) F_\pi^2(Q^2, t)$$



Drawbacks of this technique

1. Isolating  $\sigma_L$  experimentally challenging
2. Theoretical uncertainty in form factor extraction.

$K^+$  pole is further in the unphysical region, uncertainties will be larger

# Experimental Issues

- Deep Exclusive Meson Production (DEMP) cross section is small, can exclusive  $p(e, e' \pi^+) n$  and  $p(e, e' K^+) \Lambda$  channels be cleanly identified?
  - High momentum, forward angle ( $5.5^\circ$ ) meson detection is required, with good Particle ID to separate  $\pi^+$ ,  $K^+$ ,  $p$
  - Good momentum resolution required to reconstruct crucial kinematics, such as  $M_{miss}$ ,  $Q^2$ ,  $W$ ,  $t$
- Need to measure the longitudinal cross section  $d\sigma_L/dt$  needed for form factor extraction



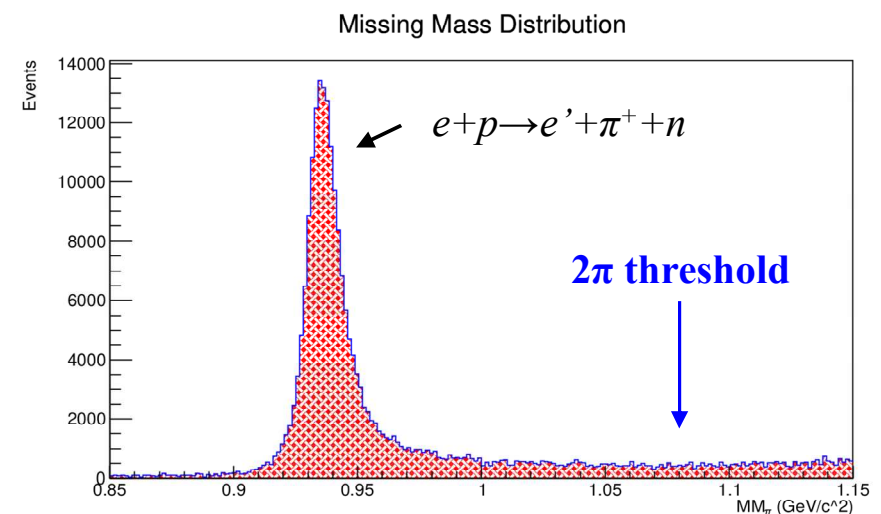
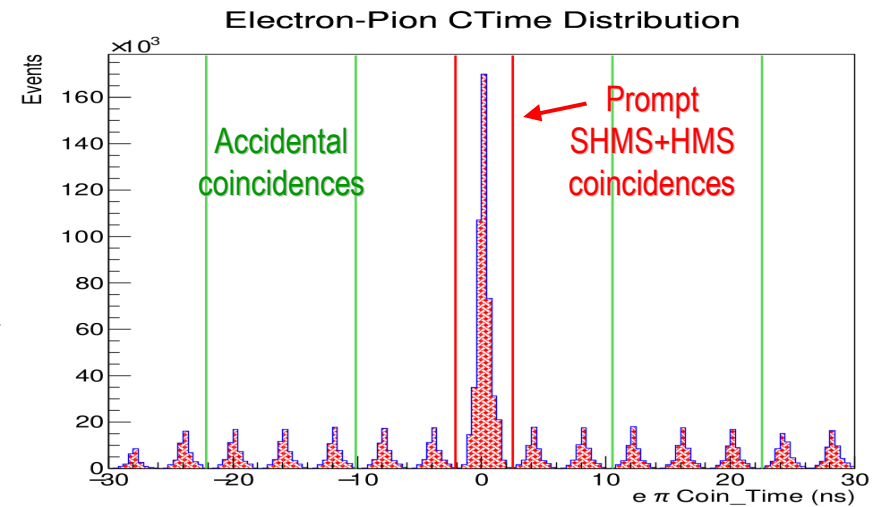
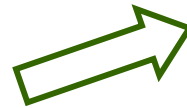
Hall C of  
Jefferson Lab  
has been  
optimized for  
specifically  
such studies

# $p(e, e' \pi^+) n$ Event Selection

## Coincidence measurement between charged pions in SHMS and electrons in HMS

Easy to isolate  
exclusive channel

- Excellent particle identification
- CW beam minimizes “accidental” coincidences
- Missing mass resolution easily excludes 2-pion contributions



PionLT experiment E12-19-006 Data

$Q^2=1.60$ ,  $W=3.08$ ,  $x=0.157$ ,  $\varepsilon=0.685$

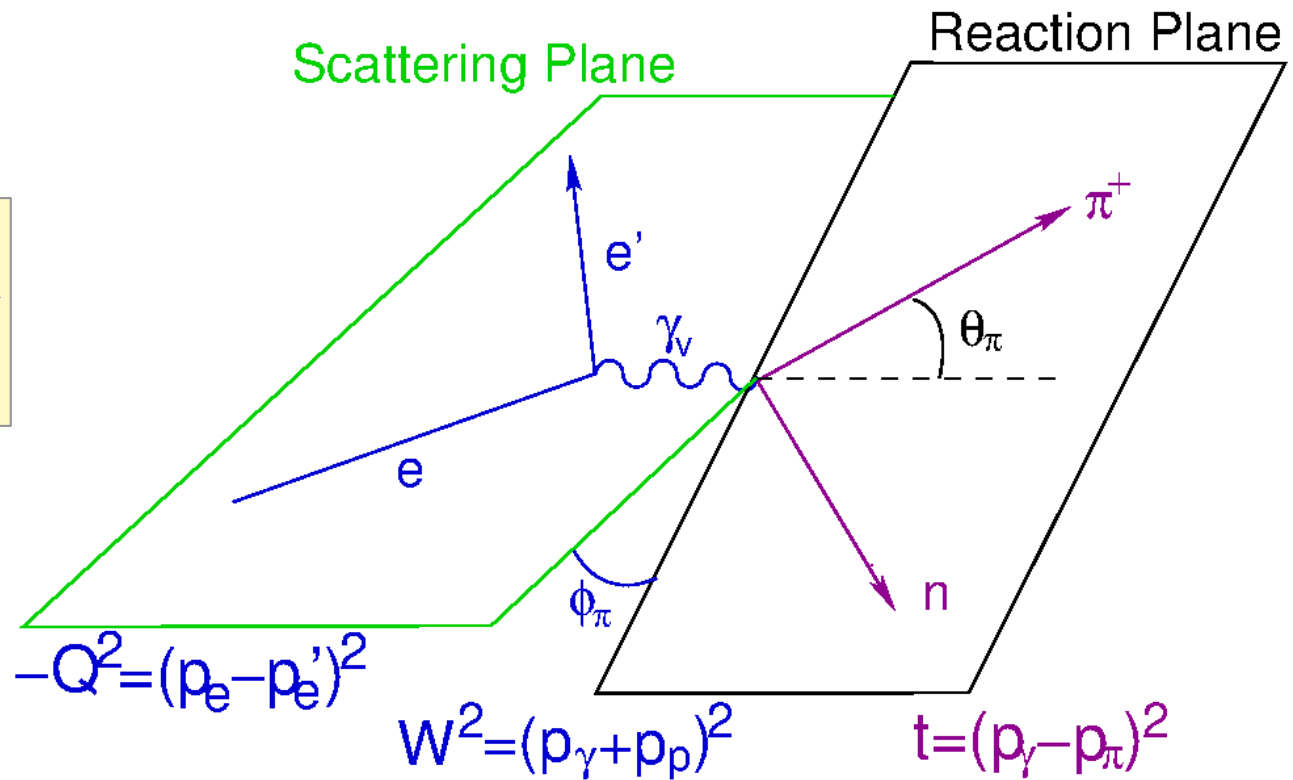
$E_{\text{beam}}=9.177$  GeV,  $P_{\text{SHMS}}=+5.422$  GeV/c,  $\theta_{\text{SHMS}}=10.26^\circ$  (left)

Plots by Muhammad Junaid

$$2\pi \frac{d^2\sigma}{dt d\phi} = \varepsilon \frac{d\sigma_L}{dt} + \frac{d\sigma_T}{dt} + \sqrt{2\varepsilon(\varepsilon+1)} \frac{d\sigma_{LT}}{dt} \cos\phi + \varepsilon \frac{d\sigma_{TT}}{dt} \cos 2\phi$$

Virtual-photon polarization:

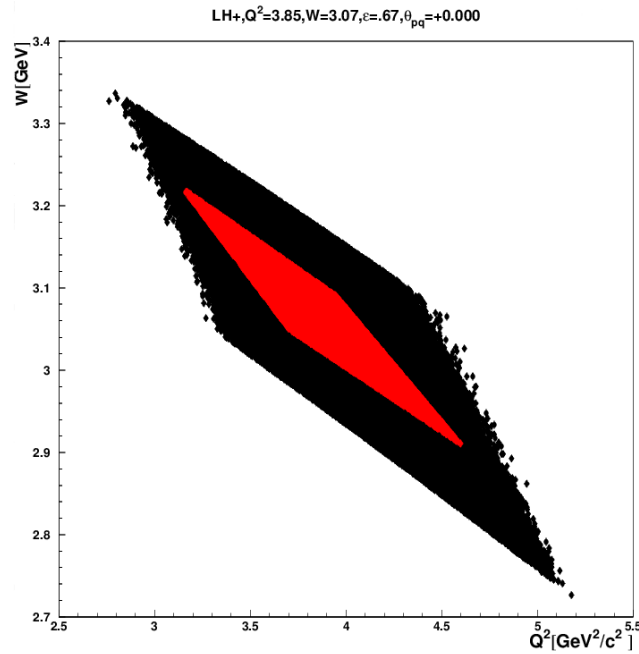
$$\varepsilon = \left( 1 + 2 \frac{(E_e - E_{e'})^2 + Q^2 \tan^2 \frac{\theta_{e'}}{2}}{Q^2} \right)^{-1}$$



- L-T separation required to separate  $\sigma_L$  from  $\sigma_T$
- Need to take data at smallest available  $-t$ , so  $\sigma_L$  has maximum contribution from the  $\pi^+$  pole
- Need to measure  $t$ -dependence of  $\sigma_L$  at fixed  $Q^2, W$

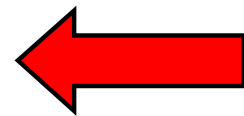
# The different pion arm (SHMS) settings are combined to yield $\phi$ -distributions for each $t$ -bin

$$2\pi \frac{d^2\sigma}{dt d\phi} = \varepsilon \frac{d\sigma_L}{dt} + \frac{d\sigma_T}{dt} + \sqrt{2\varepsilon(\varepsilon+1)} \frac{d\sigma_{LT}}{dt} \cos\phi + \varepsilon \frac{d\sigma_{TT}}{dt} \cos 2\phi$$

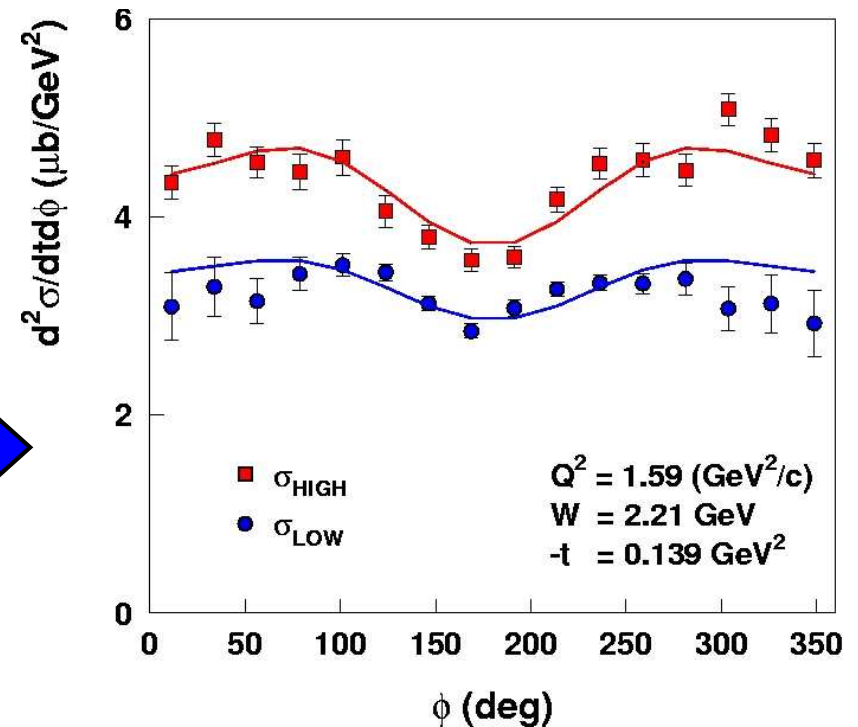


Diamond cuts define common ( $W, Q^2$ ) coverage at both  $\varepsilon$

Simulated SHMS+HMS acceptance at  $Q^2=3.85, W=3.07$   
 ■ High  $\varepsilon=0.67$  ■ Low  $\varepsilon=0.30$



■ Extract  $\sigma_L$  by simultaneous fit of L, T, LT, TT using measured azimuthal angle ( $\phi_\pi$ ) and knowledge of photon polarization ( $\varepsilon$ )



# Extract $F_\pi(Q^2)$ from JLab $\sigma_L$ data

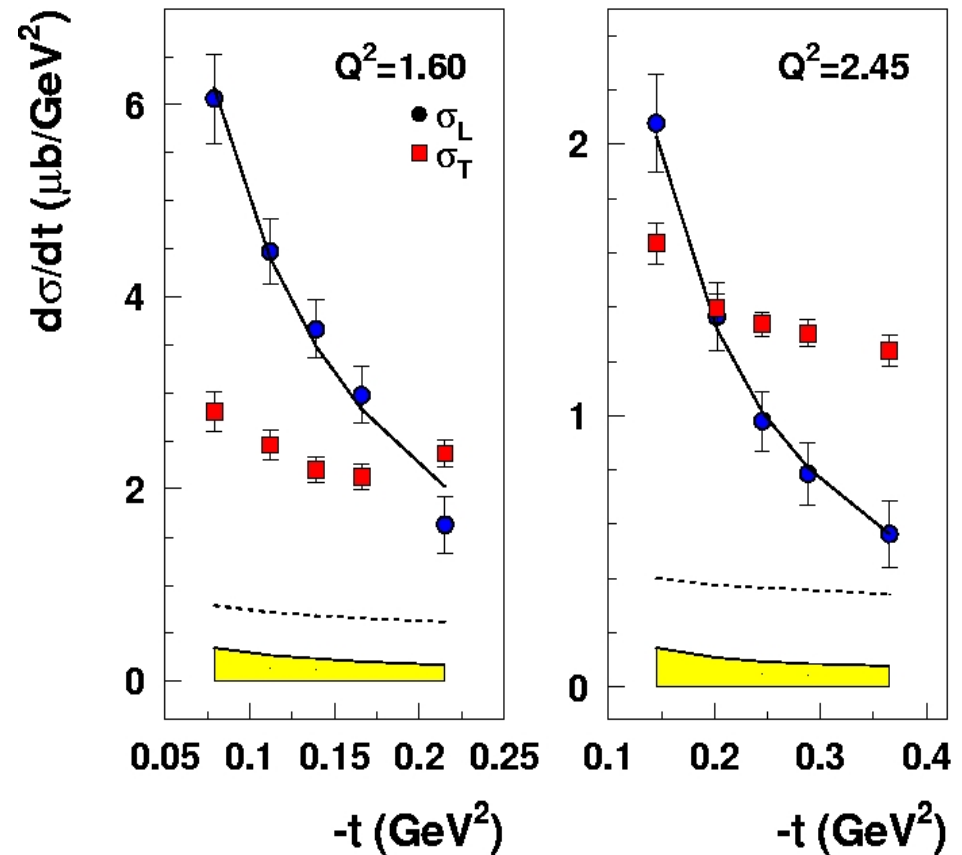
Model incorporates  $\pi^+$  production mechanism and spectator neutron effects:

## VGL Regge Model:

- Feynman propagator  $\left( \frac{1}{t - m_\pi^2} \right)$   
replaced by  $\pi$  and  $\rho$  Regge propagators.
  - Represents the exchange of a series of particles, compared to a single particle.
- Free parameters:  $\Lambda_\pi, \Lambda_\rho$  (trajectory cutoff)  
[Vanderhaeghen, Guidal, Laget, PRC 57(1998)1454]
- At small  $-t$ ,  $\sigma_L$  only sensitive to  $F_\pi$

$$F_\pi = \frac{1}{1 + Q^2 / \Lambda_\pi^2}$$

Fit to  $\sigma_L$  to model  
gives  $F_\pi$  at each  $Q^2$



Error bars indicate statistical and random (pt-pt) systematic uncertainties in quadrature.

Yellow band indicates the correlated (scale) and partly correlated (t-corr) systematic uncertainties.

$$\Lambda_\pi^2 = 0.513, 0.491 \text{ GeV}^2, \Lambda_\rho^2 = 1.7 \text{ GeV}^2.$$

# Current and Projected $F_\pi$ Data

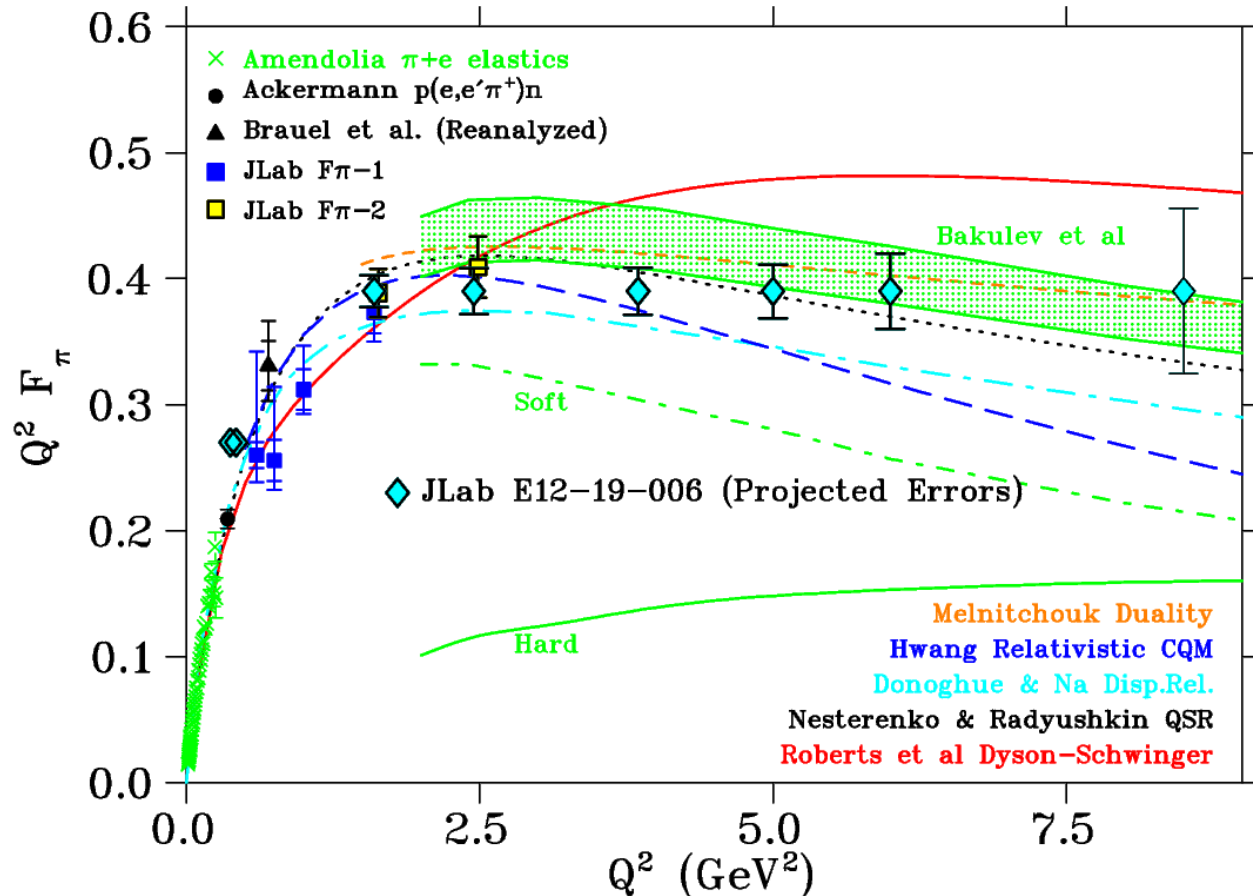
SHMS+HMS will allow measurement of  $F_\pi$  to much higher  $Q^2$

No other facility worldwide can perform this measurement

Data taking completed September 2022 (E12-19-006: G. Huber, D. Gaskell and T. Horn, spokespersons)

y-positions of projected points are arbitrary

Error bars are calculated from obtained statistics and projected systematic uncertainties



The  $\sim 10\%$  measurement of  $F_\pi$  at  $Q^2=8.5 \text{ GeV}^2$  is at higher  $-t_{min}=0.45 \text{ GeV}^2$

The pion form factor is the clearest test case for studies of QCD's transition from non-perturbative to perturbative regions

## Isolate Exclusive Final States via Missing Mass

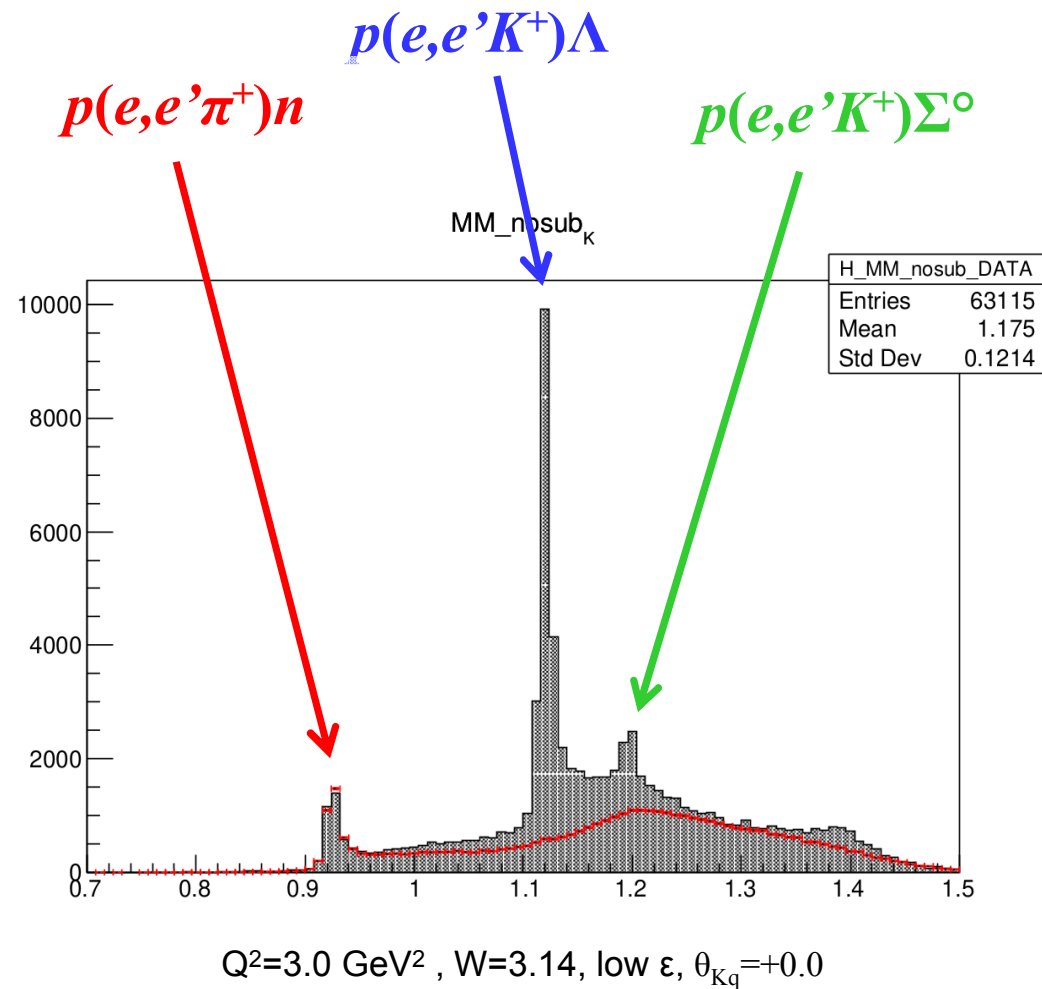
$$M_X = \sqrt{(E_{\text{det}} - E_{\text{init}})^2 - (p_{\text{det}} - p_{\text{init}})^2}$$

- Spectrometer coincidence acceptance allows for simultaneous studies of  $\Lambda$  and  $\Sigma^0$  channels.

- Kaon-pole dominance test through

$$\frac{\sigma_L(\gamma^* p \rightarrow K^+ \Sigma^0)}{\sigma_L(\gamma^* p \rightarrow K^+ \Lambda^0)}$$

- Should be similar to ratio of  $g^2_{pK\Lambda}/g^2_{pK\Sigma}$  coupling constants if t-channel exchange dominates.



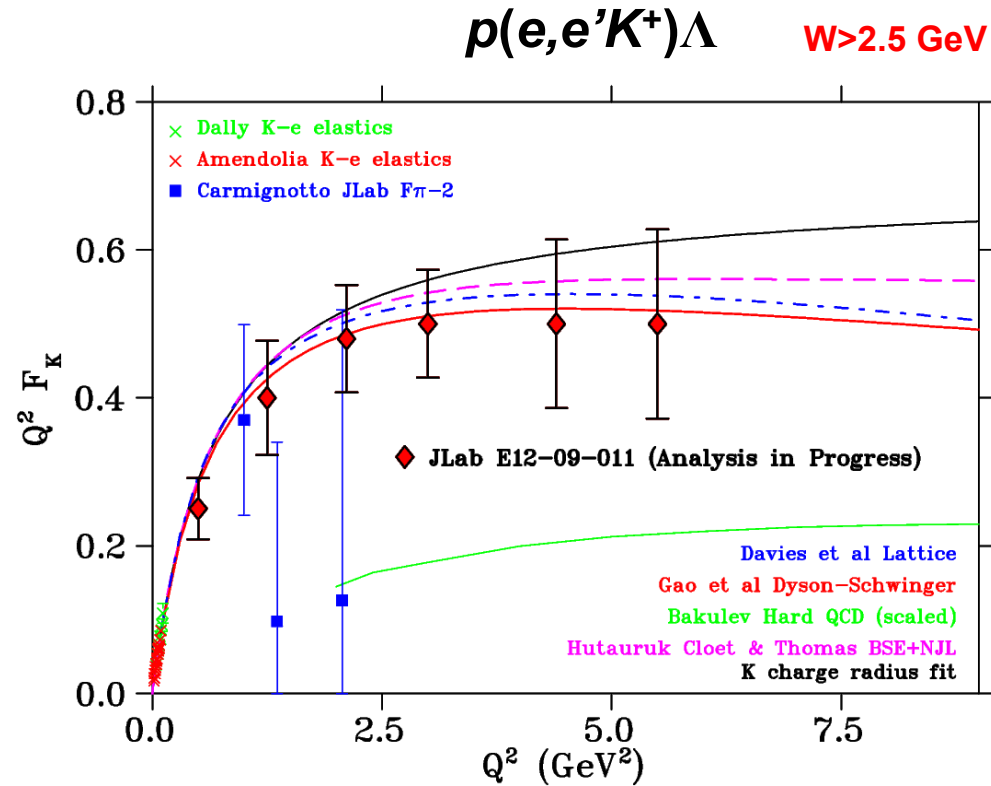
Plot by Richard Trotta (CUA/Virginia)



- Measure the  $-t$  dependence of the  $p(e, e'K^+)\Lambda, \Sigma^0$  cross section at fixed  $Q^2$  and  $W > 2.5$  GeV to search for evidence of  $K^+$  pole dominance in  $\sigma_L$ 
  - Separate the cross section components: L, T, LT, TT
  - First L/T measurement above the resonance region in  $K^+$  production
- If warranted by the data, extract the  $Q^2$  dependence of the kaon form factor to shed new light on QCD's transition to quark-gluon degrees of freedom.
- Even if we cannot extract the kaon form factor, the measurements are important.
  - $K^+\Lambda$  and  $K^+\Sigma^0$  reaction mechanisms provide valuable information in our study of hadron structure
  - Flavor degrees of freedom provide important information for QCD model building and understanding of basic coupling constants

# Projected Uncertainties for $K^+$ Form Factor

- First measurement of  $F_K$  well above the resonance region.
- Measure form factor to  $Q^2=3 \text{ GeV}^2$  with good overlap with elastic scattering data.
  - Limited by  $-t < 0.2 \text{ GeV}^2$  requirement to minimize non-pole contributions.
- Data will provide an important second  $q\bar{q}$  system for theoretical models, this time involving a strange quark.

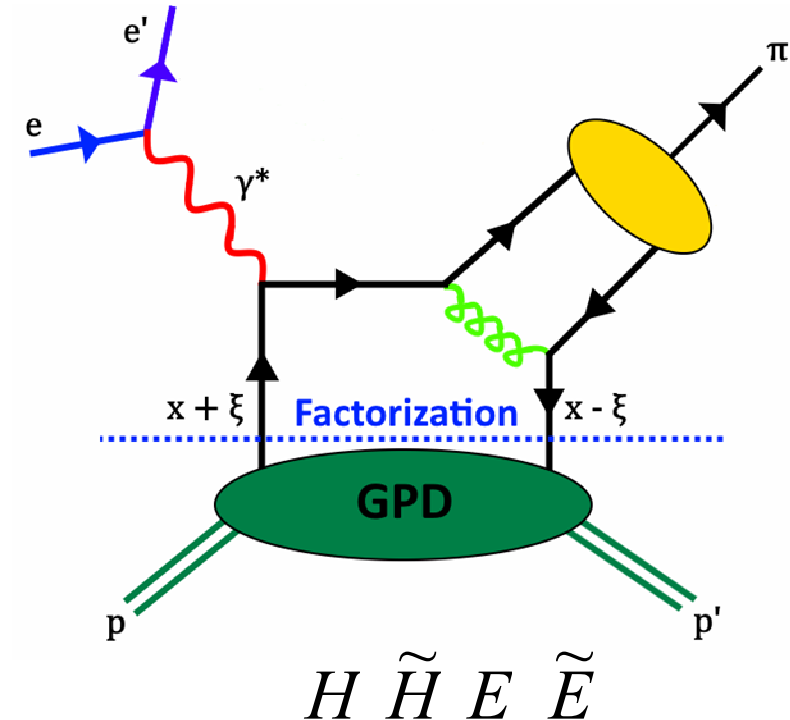


Extraction of  $F_K$  from  $Q^2 > 4 \text{ GeV}^2$  data is more uncertain, due to higher  $-t_{\min}$

- **Partially completed as an early SHMS commissioning experiment: LT-separation**  
(E12-09-011: T. Horn, G. Huber and P. Markowitz, spokespersons)
- **Data under analysis, expecting final results next year**  
— R. Trotta (CUA/Virginia)

- At sufficiently high  $Q^2$ , the Hard–Soft Factorization Theorem separates the reaction amplitude into two parts:

- Hard scattering process, where perturbative QCD can be used
- A non–perturbative (soft) part, where the response of the target nucleon to the virtual photon probe is encoded in GPDs

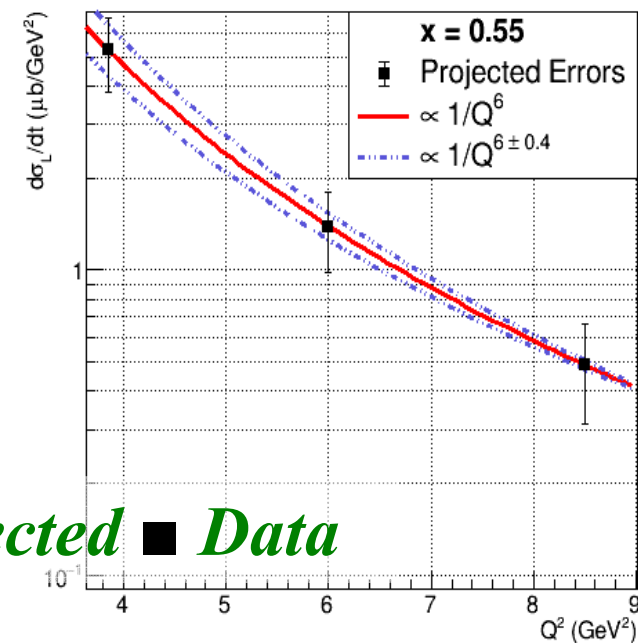
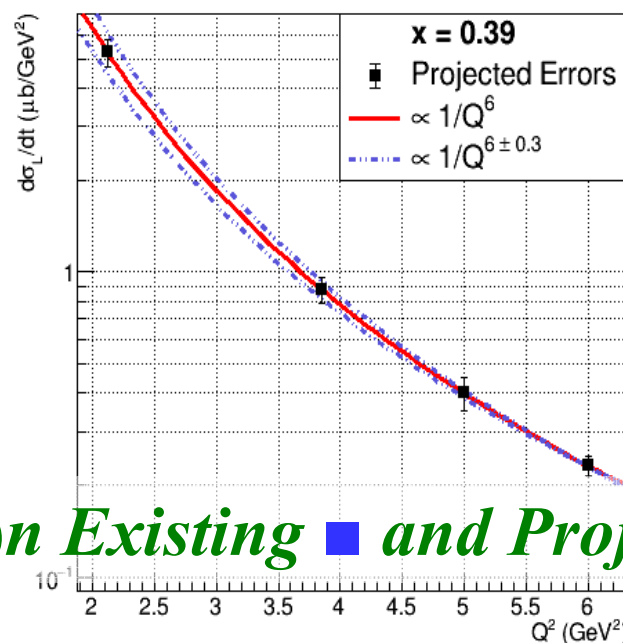
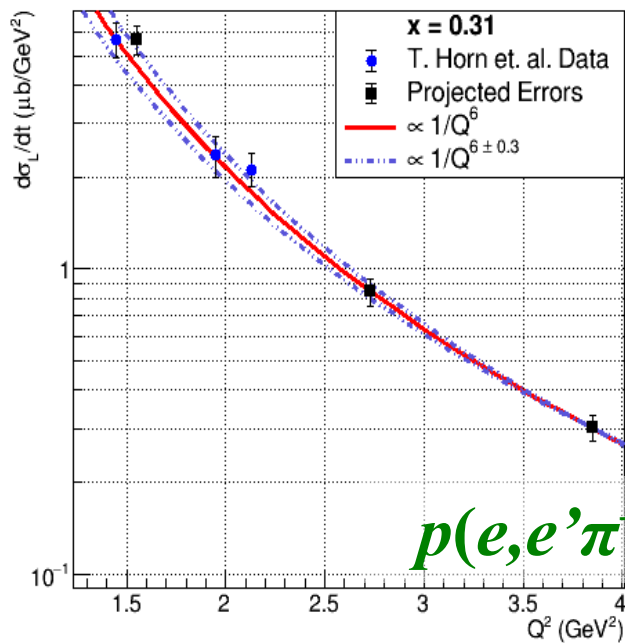


Collins, Frankfurt, Strikman PRD 56(1997)2982

- To access physics contained in GPDs, one is limited to the kinematic regime where hard–soft factorization applies
- No single criterion for applicability, but tests of necessary conditions can provide evidence that  $Q^2$  scaling regime reached

# Testing Factorization: $p(e, e' \pi^+) n$

- One of most stringent tests of factorization is  $Q^2$  dependence of  $\pi/K$  electroproduction cross sections
  - $\sigma_L$  scales to leading order as  $Q^{-6}$
  - As  $Q^2$  becomes large:  $\sigma_L \gg \sigma_T$
- If we show factorization regime is not reached, it will have major implications for meson production GPD experiments in this  $Q^2$  regime (Some of these experiments are already taking data!)

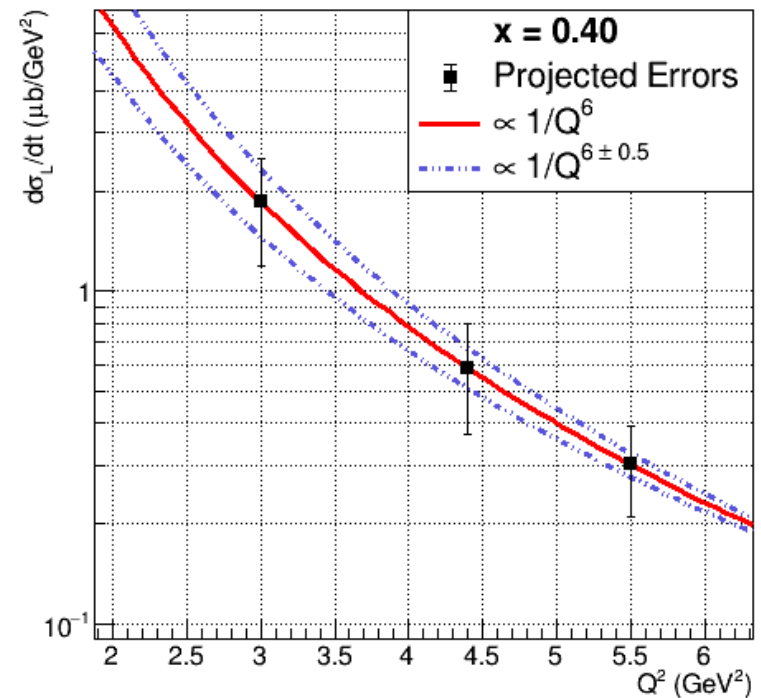
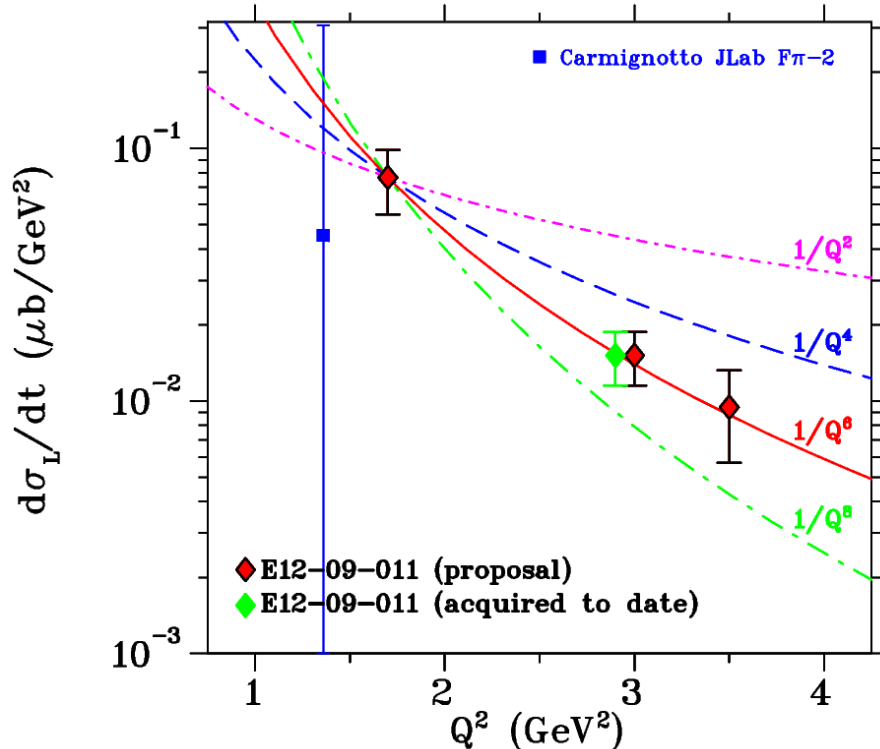


$p(e, e' \pi^+) n$  Existing ■ and Projected ■ Data

- E12-19-006 data taking completed 2022
- PhD students: N. Heinrich, M. Junaid Spokespersons: D. Gaskell, T. Horn, GMH

# Important 2<sup>nd</sup> Test: $p(e, e'K^+)\Lambda$

- Experimental validation of onset of hard scattering regime is essential for reliable interpretation of JLab GPD program results
- Is onset of scaling different for kaons than pions?
- $K^+$  and  $\pi^+$  together provide quasi model-independent study



$p(e, e'K^+)\Lambda$  Existing ■ and Projected ◆♦■ Data

- E12-09-011 data taking partially completed in 2019
- Data for  $x_B=0.40$  scan in hand. Data for  $x_B=0.25$  scan only partly acquired.
- Spokespersons: T. Horn, P. Markowitz, GMH

- Higher  $Q^2$  data on  $\pi^+$  and  $K^+$  form factors are vital to our better understanding of hadronic physics
  - PionLT (E12–19–006) has for the first time, since the pioneering measurements at Cornell in 1970's, acquired the high quality data needed to test these theoretical developments with authority
  - KaonLT (E12–09–011) partially completed. First results hopefully out next year
- Factorization studies are crucial if the field is to fully utilize the information encoded in GPDs, as GPDs are only accessible experimentally in the hard–soft factorization regime
  - PionLT (E12–19–006) has acquired data for LT–separated  $p(e, e'\pi^+)n$   $Q^{-n}$  scans at  $x_B=0.31, 0.39, 0.55$
  - KaonLT (E12–09–011) has acquired  $p(e, e'K^+)\Lambda$  data for  $Q^{-n}$  scan at  $x_B=0.40$ , eventual extension to  $x_B=0.25$